

SASAKI-EINSTEIN METRICS ON A CLASS OF 7-MANIFOLDS

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ABSTRACT. In this note we give an explicit construction of Sasaki-Einstein metrics on a class of simply connected 7-manifolds with the rational cohomology of the 2-fold connected sum of $S^2 \times S^5$. The homotopy types are distinguished by torsion in H^4 .

INTRODUCTION

Sasaki-Einstein metrics on 7-manifolds continue to play an important role in M-theory as well as black hole physics [Spa11, GLPS17, MT18]. An important reason for this is that Sasaki-Einstein manifolds admit supersymmetry and are used in the AdS-CFT correspondence. For such reasons it seems important to have a large list of explicit examples of Sasaki-Einstein 7-manifolds that can be used as possible models. Since these SE metrics are toric, general existence of such metrics is well known [FOW09, Leg16]. What is new here is an explicit construction of such metrics, their relation with Bott manifolds (orbifolds), and the topological description of the 7-manifolds. We give an explicit construction of toric Sasaki-Einstein (SE) 7-manifolds which can be represented as S^1 orbundles over 2-twist stage 3 Bott orbifolds. All of these are obtained by adding orbifold structures to certain stage 3 Bott manifolds which were studied in [BCTF18]. The 7-manifolds all have the rational cohomology of the 2-fold connected sum $2(S^2 \times S^5)$ and are generalizations of the SE 7-manifolds given by Theorem 1.2 of [BTF15].

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1. STAGE THREE BOTT TOWERS AND ORBIFOLDS

Following [GK94] and [BCTF18] we consider *Bott towers* which in arbitrary dimension is represented by a lower triangular unipotent matrix A over \mathbb{Z} . Here we deal only with stage 3 Bott towers, so the matrix A in [GK94, BCTF18] takes the form

$$(1) \quad A = \begin{pmatrix} 1 & 0 & 0 \\ a & 1 & 0 \\ b & c & 1 \end{pmatrix}$$

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with $a, b, c \in \mathbb{Z}$. The Bott manifold can be realized as the quotient of $S^3 \times S^3 \times S^3$ by the \mathbb{T}^3 action

$$(2) \quad (z_j^0, z_j^\infty)_{j=1}^3 \mapsto (t_j z_j^0, \prod_{i=1}^3 t_i^{A_j^i} z_j^\infty).$$

The quotient M_3 which is called a *Bott manifold* can be represented as a sequence, called a *Bott tower*

$$(3) \quad M_3 \xrightarrow{\pi_3} M_2 \xrightarrow{\pi_2} M_1 \xrightarrow{\pi_1} M_0 = pt,$$

where the j th S^3 is written as $|z_j^0|^2 + |z_j^\infty|^2 = 1$. Then M_k is the compact complex manifold arising as the total space of the $\mathbb{C}\mathbb{P}^1$ bundle $\pi_k: \mathbb{P}(\mathbb{1} \oplus \mathcal{L}_k) \rightarrow M_{k-1}$. At each stage we have *zero* and *infinity sections* $\sigma_k^0: M_{k-1} \rightarrow M_k$ and $\sigma_k^\infty: M_{k-1} \rightarrow M_k$ which respectively identify M_{k-1} with $\mathbb{P}(\mathbb{1} \oplus 0)$ and $\mathbb{P}(0 \oplus \mathcal{L}_k)$. We consider these to be part of the structure of the Bott tower $(M_k, \pi_k, \sigma_k^0, \sigma_k^\infty)_{k=1}^n$. Here in our case $M_3 = M_3(a, b, c)$ is a stage 3 Bott manifold, $M_2(a)$ is a Hirzebruch surface $\mathcal{H}_a = \mathbb{P}(\mathbb{1} \oplus \mathcal{O}(a)) \rightarrow \mathbb{C}\mathbb{P}^1$, and $M_1 = \mathbb{C}\mathbb{P}^1$. $M_3(a, b, c)$ can be viewed as the total space of $\mathbb{C}\mathbb{P}^1$ bundle over the Hirzebruch surface \mathcal{H}_a , and also a bundle of Hirzebruch surfaces over $\mathbb{C}\mathbb{P}^1$ with fiber \mathcal{H}_c . Bott towers form the object set \mathcal{BT}_0 of a groupoid whose morphisms \mathcal{BT}_1 are biholomorphisms [BCTF18], and elements of the quotient space $\mathcal{BT}_0/\mathcal{T}_1$ are identified with biholomorphism classes of Bott manifolds. Since Bott manifolds are toric, they are described by a fan, and it follows from the Bott tower description (3) that the fan of the Bott tower $M_3(a, b, c)$ is described by the primitive collections (cf. [CLS11])

$$(4) \quad \{v_1, u_1\}, \quad \{v_2, u_2\}, \quad \{v_3, u_3\}$$

with normal vectors $u_1 = -v_1 - av_2 - bv_3, u_2 = -v_2 - cv_3, u_3 = -v_3$, and thus has the combinatorial type of a cube. The symmetry group of a cube is the Coxeter group $BC_3 \cong \text{Sym}_3 \times \mathbb{Z}_2^3$ where Sym_3 is the symmetric group on 3 letters. However, not all elements of BC_3 are induced by equivalences. We refer to [BCTF18] and references therein for details.

The structure of Bott towers implies the existence of 3 pairs of \mathbb{T}^3 invariant divisors $\{D_{v_j}, D_{u_j}\}$ which are the zero and infinity sections of $M_j \xrightarrow{\pi_j} M_{j-1}$ with $j = 1, 2, 3$. Thus, elements of the subgroup \mathbb{Z}_2^3 are induced by the *fiber inversion* maps that interchange the zero and infinity sections, so these equivalences always exist. However, elements of Sym_3 are induced by equivalences only in special cases (see Lemma 1.11 and Example 1.6 of [BCTF18] for details).

Recall (Definition 1.6 of [BCTF18]) that the *holomorphic twist* of an n dimensional Bott tower is the number $t \in \{0, \dots, n-1\}$ of holomorphically nontrivial $\mathbb{C}\mathbb{P}^1$ bundles in the tower. So for $n = 3$ there are only 3 possibilities $t = 0, 1, 2$. Of course, $t = 0$ is the well understood product $\mathbb{C}\mathbb{P}^1 \times \mathbb{C}\mathbb{P}^1 \times \mathbb{C}\mathbb{P}^1$, whereas $t = 1$ leads to the Koiso-Sakane case, so here we restrict our attention to $t = 2$. For the $t = 2$ case we have a nontrivial holomorphic bundle over a Hirzebruch surface \mathcal{H}_a with $a \neq 0$ whose fiber is $\mathbb{C}\mathbb{P}^1$. The case of interest to us can be obtained as an $S_{\mathbf{w}}^3$ -join with the $Y^{p,q}$ structures of [GMSW04] on $S^2 \times S^3$, that is, $(S^2 \times S^3) \star_{l_1, l_2} S_{\mathbf{w}}^3$. So the stage 3 Bott manifold has an additional orbifold structure which we now describe.

1.1. Invariant Divisors. Here we consider the \mathbb{T}^3 invariant divisors D_{v_i}, D_{u_i} defined by the normal vectors v_i, u_i with equivalences

$$D_{v_1} \sim D_{u_1}, \quad D_{v_2} \sim aD_{u_1} + D_{u_2}, \quad D_{v_3} \sim bD_{u_1} + cD_{u_2} + D_{u_3}.$$

We have 4 sets of distinguished invariant bases of the Chow group $A_2(M_3)$ of invariant divisor classes

$$\begin{aligned} (5) \quad & \{[D_{u_1}], [D_{u_2}], [D_{u_3}]\}, \\ (6) \quad & \{[D_{u_1}], [D_{u_2}], [D_{v_3}]\}, \\ (7) \quad & \{[D_{u_1}], [D_{v_2}], [D_{u_3}]\}, \\ (8) \quad & \{[D_{u_1}], [D_{v_2}], [D_{v_3}]\}. \end{aligned}$$

This gives rise to the 4 sets of dual bases of cohomology classes in $H^2(M_3, \mathbb{Z})$, viz

$$\begin{aligned} (9) \quad & \{x_1, x_2, x_3\}, \\ (10) \quad & \{x_1, x_2, y_3\}, \\ (11) \quad & \{x_1, y_2, x_3\}, \\ (12) \quad & \{x_1, y_2, y_3\}, \end{aligned}$$

where x_i is dual to $[D_{u_i}]$ and y_i is dual to $[D_{v_i}]$. Both the ample and Kähler cones can be easily worked out, see Example 3.3 of [BCTF18].

1.2. Bott Orbifolds and log pairs. We are interested in these Bott manifolds, but with an additional special orbifold structure along the \mathbb{T}^3 -invariant divisors D_{v_i}, D_{u_i} . The orbifold structure on $M_3(a, b, c)$ that we are interested in is given by the log pair $(M_3(a, b, c), \Delta_{\mathbf{m}})$ where $\Delta_{\mathbf{m}}$ is the branch divisor

$$(13) \quad \left(1 - \frac{1}{m_1^0}\right)D_{v_1} + \left(1 - \frac{1}{m_1^\infty}\right)D_{u_1} + \left(1 - \frac{1}{m_2^0}\right)D_{v_2} + \left(1 - \frac{1}{m_2^\infty}\right)D_{u_2} + \left(1 - \frac{1}{m_3^0}\right)D_{v_3} + \left(1 - \frac{1}{m_3^\infty}\right)D_{u_3}$$

where $m_j^0, m_j^\infty \in \mathbb{Z}^+$ are the ramification indices. We define $\Delta_{\mathbf{m}_3} = \Delta_{\mathbf{m}}$ and

$$\begin{aligned} \Delta_{\mathbf{m}_2} &= \left(1 - \frac{1}{m_1^0}\right)D_{v_1} + \left(1 - \frac{1}{m_1^\infty}\right)D_{u_1} + \left(1 - \frac{1}{m_2^0}\right)D_{v_2} + \left(1 - \frac{1}{m_2^\infty}\right)D_{u_2} \\ \Delta_{\mathbf{m}_1} &= \left(1 - \frac{1}{m_1^0}\right)D_{v_1} + \left(1 - \frac{1}{m_1^\infty}\right)D_{u_1}. \end{aligned}$$

Clearly we have

$$\Delta_{\mathbf{m}_1} \subset \Delta_{\mathbf{m}_2} \subset \Delta_{\mathbf{m}_3}.$$

From [GK94] or [BCTF18] one easily sees that

Lemma 1.1. *We have the sequence of Bott towers of log pairs*

$$(14) \quad (M_3(a, b, c), \Delta_{\mathbf{m}_3}) \xrightarrow{\pi_3} (M_2(a), \Delta_{\mathbf{m}_2}) \xrightarrow{\pi_2} (M_1, \Delta_{\mathbf{m}_1}) \xrightarrow{\pi_1} (\{pt\}, \emptyset),$$

where $M_2(a) = \mathcal{H}_a$ is a Hirzebruch surface, $M_1 = \mathbb{CP}^1$, and π_i is the natural projection.

The invariant branch divisors are related to the section maps by

$$D_{v_3} = \sigma_2^0(\mathcal{H}_a), \quad D_{u_3} = \sigma_2^\infty(\mathcal{H}_a), \quad D_{v_2} = \pi_3^{-1}(\sigma_1^0(\mathbb{CP}^1)), \quad D_{u_2} = \pi_3^{-1}(\sigma_1^\infty(\mathbb{CP}^1))$$

and

$$D_{v_1} = (\pi_2 \circ \pi_3)^{-1}(\sigma_1^0(\{pt\})), \quad D_{u_1} = (\pi_2 \circ \pi_3)^{-1}(\sigma_1^\infty(\{pt\})).$$

We denote by \mathfrak{BD}_A the set of all such Bott orbifolds.

1.3. The Orbifold First Chern Class. We can now compute the orbifold canonical divisor K^{orb} and dually the orbifold first Chern class. We compute the orbifold first Chern class c_1^{orb} in the $\{x_j\}$ basis for any n -dimensional Bott manifold $M_n(A)$;

$$\begin{aligned} c_1^{orb}(M_n(A), \Delta_{\mathbf{m}}) &= c_1(M_n(A)) - \sum_{i=1}^n \left(\left(1 - \frac{1}{m_i^0}\right)y_i + \left(1 - \frac{1}{m_i^\infty}\right)x_i \right) \\ &= \sum_{i=1}^n (x_i + y_i) - \sum_{i=1}^n \left(\left(1 - \frac{1}{m_i^0}\right)y_i + \left(1 - \frac{1}{m_i^\infty}\right)x_i \right) \\ &= \sum_{j=1}^n \left(\frac{1}{m_j^0}y_j + \frac{1}{m_j^\infty}x_j \right) = \sum_{i=1}^n \left(\left(\frac{1}{m_i^0} + \frac{1}{m_i^\infty} \right)x_i + \sum_{j=1}^{i-1} \frac{A_i^j}{m_i^0}x_j \right) \\ (15) \quad &= \sum_{j=1}^{n-1} \left(\frac{1}{m_j^0} + \frac{1}{m_j^\infty} + \sum_{i=j+1}^n \frac{A_i^j}{m_i^0} \right)x_j + \left(\frac{1}{m_n^0} + \frac{1}{m_n^\infty} \right)x_n. \end{aligned}$$

In dimension n there are 2^{n-1} invariant bases in which to compute c_1^{orb} . This becomes much more manageable for $n = 3$. We have $c_1^{orb}(M_3(a, b, c), \Delta_{\mathbf{m}})$ in the four bases (9)-(12), respectively

$$(16) \quad \left(\frac{1}{m_1^0} + \frac{1}{m_1^\infty} + \frac{a}{m_2^0} + \frac{b}{m_3^0} \right)x_1 + \left(\frac{1}{m_2^0} + \frac{1}{m_2^\infty} + \frac{c}{m_3^0} \right)x_2 + \left(\frac{1}{m_3^0} + \frac{1}{m_3^\infty} \right)x_3,$$

$$(17) \quad \left(\frac{1}{m_1^0} + \frac{1}{m_1^\infty} + \frac{a}{m_2^0} - \frac{b}{m_3^\infty} \right)x_1 + \left(\frac{1}{m_2^0} + \frac{1}{m_2^\infty} - \frac{c}{m_3^\infty} \right)x_2 + \left(\frac{1}{m_3^0} + \frac{1}{m_3^\infty} \right)y_3,$$

$$(18) \quad \left(\frac{1}{m_1^0} + \frac{1}{m_1^\infty} - \frac{a}{m_2^\infty} + \frac{b-ac}{m_3^0} \right)x_1 + \left(\frac{1}{m_2^0} + \frac{1}{m_2^\infty} + \frac{c}{m_3^0} \right)y_2 + \left(\frac{1}{m_3^0} + \frac{1}{m_3^\infty} \right)x_3,$$

$$(19) \quad \left(\frac{1}{m_1^0} + \frac{1}{m_1^\infty} - \frac{a}{m_2^\infty} - \frac{b-ac}{m_3^\infty} \right)x_1 + \left(\frac{1}{m_2^0} + \frac{1}{m_2^\infty} - \frac{c}{m_3^\infty} \right)y_2 + \left(\frac{1}{m_3^0} + \frac{1}{m_3^\infty} \right)y_3.$$

Note also that as a cohomology class $c_1^{orb}(M_3(a, b, c), \Delta_{\mathbf{m}})$ as an element of $H^{1,1}(M_n(a, b, c), \mathbb{R})$ makes perfect sense for all $m_j^0, m_j^\infty \in \mathbb{R}^+$. We shall make use of this fact shortly.

Equations (16)-(19) implies

Lemma 1.2. *Let $(M_3(a, b, c), \Delta_{\mathbf{m}})$ be a Bott orbifold. Then the following are equivalent:*

- (1) $(M_3(a, b, c), \Delta_{\mathbf{m}})$ is log Fano,
- (2) $c_1^{orb}(M_3(a, b, c), \Delta_{\mathbf{m}})$ lies in the Kähler cone $\mathcal{K}(M_3(a, b, c))$,

(3) *the inequalities*

$$\begin{aligned}
\frac{1}{m_1^0} + \frac{1}{m_1^\infty} + \frac{a}{m_2^0} + \frac{b}{m_3^0} &> 0, & \frac{1}{m_2^0} + \frac{1}{m_2^\infty} + \frac{c}{m_3^0} &> 0 \\
\frac{1}{m_1^0} + \frac{1}{m_1^\infty} + \frac{a}{m_2^0} - \frac{b}{m_3^\infty} &> 0, & \frac{1}{m_2^0} + \frac{1}{m_2^\infty} - \frac{c}{m_3^\infty} &> 0 \\
\frac{1}{m_1^0} + \frac{1}{m_1^\infty} - \frac{a}{m_2^\infty} + \frac{(b-ac)}{m_3^0} &> 0, & \frac{1}{m_2^0} + \frac{1}{m_2^\infty} + \frac{c}{m_3^0} &> 0 \\
\frac{1}{m_1^0} + \frac{1}{m_1^\infty} - \frac{a}{m_2^\infty} - \frac{(b-ac)}{m_3^\infty} &> 0, & \frac{1}{m_2^0} + \frac{1}{m_2^\infty} - \frac{c}{m_3^\infty} &> 0
\end{aligned}$$

hold.

2. M^7 AS THE JOIN $Y^{p,q} \star_{l_1, l_2} S_{\mathbf{w}}^3$

Not every S^1 orbibundle over a Bott orbifold can be realized as a join; however, the 2 twist stage 3 Bott orbifolds that we study here can be realized as Kähler quotients of the join $Y^{p,q} \star_{l_1, l_2} S_{\mathbf{w}}^3$ where $Y^{p,q}$ are the well known SE structures on $S^2 \times S^3$ discovered by the physicists [GMSW04]. In fact, since $Y^{p,q}$ is itself a join of two S^3 's, it is an iterated join of three S^3 's which in the terminology of [BHLTF18] is completely cone decomposable. Now from Example 6.8 of [BTF16] we have

$$(20) \quad Y^{p,q} = S^3 \star_{l,p} S_{\frac{p+q}{l}, \frac{p-q}{l}}^3$$

where $l = \gcd(p+q, p-q)$ which equals 2 if p, q are both odd, and equals 1 if p, q have opposite parity. Note that we choose the standard SE structure on the lefthand S^3 factor, whereas it is the weighted Sasakian structure with weights $(\frac{p+q}{l}, \frac{p-q}{l})$ on the righthand S^3 .

If we assume that $p > q \geq 1$ are such that $\sqrt{4p^2 - 3q^2} \in \mathbb{N}$, then the Sasaki-Einstein structure on $Y^{p,q}$ is quasi-regular [GMSW04]. Indeed, the η -Einstein structure corresponds to a ray in the so-called \mathbf{w} subcone is determined by co-prime solutions (v_2^0, v_2^∞) of

$$(21) \quad \int_{-1}^1 ((v_2^0 - v_2^\infty) - (v_2^0 + v_2^\infty)\mathfrak{z})(((p+q)v_2^\infty + (p-q)v_2^0) + ((p+q)v_2^\infty - (p-q)v_2^0)\mathfrak{z}) d\mathfrak{z} = 0$$

(from e.g. (68) in [BTF16]). Following Theorem 3.8 in [BTF16], for any choice of quasi-regular ray determined by the co-prime pair (v_2^0, v_2^∞) , the quotient Hirzebruch orbifold is $(\mathcal{H}_a, \Delta_{\mathbf{m}_2})$ where $\Delta_{\mathbf{m}_2} = (1 - \frac{1}{m_2^0})D_{v_2} + (1 - \frac{1}{m_2^\infty})D_{u_2}$ is the branch divisor with

$$\begin{aligned}
\mathbf{m}_2 &= (m_2^0, m_2^\infty) = m_2(v_2^0, v_2^\infty), \\
m_2 &= \frac{p}{\gcd(p, |\frac{p+q}{l}v_2^\infty - \frac{p-q}{l}v_2^0|)}, \\
a &= \frac{(p+q)m_2^\infty - (p-q)m_2^0}{p}.
\end{aligned}$$

The join $M_{l_1, l_2, \mathbf{w}} = Y^{p, q} \star_{l_1, l_2} S_{\mathbf{w}}^3$, where $Y^{p, q}$ has a quasi-regular Sasaki structure as above, can then be obtained as the quotient of the following \mathbb{T}^2 action on $S^3 \times S^3 \times S^3$:

$$(23) \quad (x, u; u_1, u_2; z_1, z_2) \mapsto (x, e^{ipl_2\theta} u; e^{i(l_2 m_2 v_2^0 \phi - (p+q)\theta)} u_1, e^{i(l_2 m_2 v_2^\infty \phi - (p-q)\theta)} u_2; e^{-il_1 w_1 \phi} z_1, e^{-il_1 w_2 \phi} z_2).$$

First we notice that without loss of generality we can assume that $\gcd(l_1, m_2) = 1$, for otherwise we can redefine ϕ . So our gcd conditions are $\gcd(l_1, l_2 m_2) = 1$, $\gcd(w_1, w_2) = 1 = \gcd(v_2^0, v_2^\infty)$, and $\gcd(p, q) = 1$. Note that $p+q$ and $p-q$ can have a common factor if and only if both p and q are odd, in which case the common factor is 2. However, $M_{l_1, l_2, \mathbf{w}}$ may not be a smooth manifold. we have

Lemma 2.1. *The join $M_{l_1, l_2, \mathbf{w}} = Y^{p, q} \star_{l_1, l_2} S_{\mathbf{w}}^3$, with Sasakian structure on $Y^{p, q}$ given by the Reeb field $\xi_{\mathbf{v}_2}$ with $\mathbf{v}_2 = (v_2^0, v_2^\infty)$, is a smooth manifold if and only if $\gcd(l_2 m_2 v_2^i, l_1 w_j) = 1$ for $i = 0, \infty$ and $j = 1, 2$, where $m_2 = \frac{p}{\gcd(p, |\frac{p+q}{l_1} v_2^\infty - \frac{p-q}{l_1} v_2^0|)}$.*

Proof. From Proposition 7.6.7 of [BG08] $M_{l_1, l_2, \mathbf{w}}$ is smooth if and only if $\gcd(l_1 \Upsilon_2, l_2 \Upsilon_1) = 1$ where Υ_1 is the order of $Y^{p, q}$ and Υ_2 is the order of $S_{\mathbf{w}}^3$. The latter is $\Upsilon_2 = w_1 w_2$ with w_1, w_2 coprime. The order Υ_1 with quasi-regular Reeb field $\xi_{\mathbf{v}_2}$ is $\Upsilon_1 = m_2 v_2^0 v_2^\infty$. \square

The analysis in Section 3 of [BTF16] holds equally well when the manifold M in the join $M \star_{l_1, l_2} S_{\mathbf{w}}^3$ has any quasi-regular Sasakian structure. The major difference is having more complicated computations. For example we need the Fano index of $Y^{p, q}$. As described above, the quotient of any quasi-regular Sasakian structure in the \mathbf{w} subcone of $Y^{p, q}$ is a Hirzebruch orbifold of the form $(\mathcal{H}_a, \Delta_{\mathbf{m}_2})$. We have

Lemma 2.2. *Let $\xi_{\mathbf{v}_2}$ be a quasiregular Reeb vector field with quotient Hirzebruch orbifold $(\mathcal{H}_a, \Delta_{\mathbf{m}_2})$ with $a > 0$. Then its Fano index $J_{\mathbf{v}_2}$ is given by*

$$J_{\mathbf{v}_2} = \gcd(2m_2 v_2^0 + a)v_2^\infty, v_2^0 + v_2^\infty)$$

where $\mathbf{m}_2 = m_2(v_2^0, v_2^\infty)$ and v_2^0, v_2^∞ are coprime.

Proof. Recall (Definition 4.4.24 of [BG08]) that the Fano index J of an orbifold \mathcal{Z} is the largest positive integer k such that $p^* c_1^{orb}/k$ is an element of $H_{orb}^2(\mathcal{Z}, \mathbb{Z}) = H^2(B\mathcal{Z}, \mathbb{Z})$. Now the classifying map $p : B\mathcal{Z} \rightarrow \mathcal{Z}$ is an $m_2 v_2^0 v_2^\infty$ -fold cover, and c_1^{orb} is

$$(24) \quad c_1^{orb} = (2 + \frac{a}{m_2 v_2^0})x_1 + \frac{v_2^0 + v_2^\infty}{m_2 v_2^0 v_2^\infty} x_2 = \frac{1}{m_2 v_2^0 v_2^\infty} ((2m_2 v_2^0 + a)v_2^\infty x_1 + (v_2^0 + v_2^\infty)x_2).$$

So $p^* c_1^{orb} = (2m_2 v_2^0 + a)v_2^\infty p^* x_1 + (v_2^0 + v_2^\infty)p^* x_2$ from which the result follows. \square

From now on we assume that $p > q \geq 1$ are such that $\sqrt{4p^2 - 3q^2} \in \mathbb{N}$ and $Y^{p, q}$ has the quasi-regular Sasaki-Einstein structure. Thus co-prime (v_2^0, v_2^∞) are chosen such that (21) is satisfied. We then present a version of Theorem 3.8 of [BTF16] that allows the join $M_{l_1, l_2, \mathbf{w}} = Y^{p, q} \star_{l_1, l_2} S_{\mathbf{w}}^3$

Theorem 2.3. *Consider the join $M_{l_1, l_2, \mathbf{w}} = Y^{p, q} \star_{l_1, l_2} S_{\mathbf{w}}^3$ where $Y^{p, q}$ has a quasi-regular Sasaki-Einstein structure determined by the co-prime pair (v_2^0, v_2^∞) satisfying Equation (21). Let $\mathcal{S}_{\mathbf{v}_3} = (\xi_{\mathbf{v}_3}, \eta_{\mathbf{v}_3}, \Phi_{\mathbf{v}_3}, g_{\mathbf{v}_3})$ be a quasi-regular Sasakian structure that lies in the \mathbf{w} subcone of the Sasaki cone with $\mathbf{v}_3 = (v_3^0, v_3^\infty)$ where v_3^0, v_3^∞ are coprime. Then the*

quotient of $M_{l_1, l_2, \mathbf{w}}$ by the S^1 action generated by $\xi_{\mathbf{v}_3}$ is the Bott orbifold given by the log pair $(M_3(a, b, c), \Delta_{\mathbf{m}})$, where a is determined by (22),

$$\begin{aligned} b &= n\hat{b} = n \frac{(2m_2v_2^0 + a)v_2^\infty}{\mathcal{J}_{\mathbf{v}_2}} \\ c &= n\hat{c} = n \frac{(v_2^0 + v_2^\infty)}{\mathcal{J}_{\mathbf{v}_2}}, \end{aligned}$$

$n = l_1 \frac{w_1v_3^\infty - w_2v_3^0}{\gcd(|w_1v_3^\infty - w_2v_3^0|, l_2)}$, $\mathcal{J}_{\mathbf{v}_2}$ is given by Lemma 2.2, and $\Delta_{\mathbf{m}}$ is the branch divisor

$$\left(1 - \frac{1}{m_1^0}\right)D_{v_1} + \left(1 - \frac{1}{m_1^\infty}\right)D_{u_1} + \left(1 - \frac{1}{m_2^0}\right)D_{v_2} + \left(1 - \frac{1}{m_2^\infty}\right)D_{u_2} + \left(1 - \frac{1}{m_3^0}\right)D_{v_3} + \left(1 - \frac{1}{m_3^\infty}\right)D_{u_3}$$

with $(m_1^0, m_1^\infty) = (1, 1)$, (m_2^0, m_2^∞) given by (22), and

$$(m_3^0, m_3^\infty) = m_3(v_3^0, v_3^\infty) = \frac{l_2}{\gcd(|w_1v_3^\infty - w_2v_3^0|, l_2)}(v_3^0, v_3^\infty).$$

Proof. We can follow the proof of Theorem 3.8 of [BTF16] with the caveat that N is a Hirzebruch orbifold $(\mathcal{H}_a, \Delta_{\mathbf{m}_2})$. As in Equation (3) of [BTF16] we have the commutative diagram

$$(25) \quad \begin{array}{ccc} S^3 \times S^3 \times S^3 & & \\ \downarrow & & \\ Y^{p,q} \times S_{\mathbf{w}}^3 & & \\ \downarrow \pi_2 & \searrow \pi_L & \\ & & Y^{p,q} \star_{l_1, l_2, \mathbf{w}} S_{\mathbf{w}}^3 \\ & \swarrow \pi_1 & \\ (\mathcal{H}_a, \Delta_{\mathbf{m}_2}) \times \mathbb{C}P^1[\mathbf{w}] & & \end{array}$$

where the π s are the obvious projections, and the orbifold $(\mathcal{H}_a, \Delta_{\mathbf{m}_2})$ is the quotient by the locally free S^1 action on $Y^{p,q}$ generated by the quasi-regular Reeb vector field $\xi_{\mathbf{m}_2}$ where $\mathbf{m}_2 = m_2(v_2^0, v_2^\infty)$ and a is given in Equations (22). The holomorphic line bundle $L_n = L^n$ now becomes the holomorphic line orbibundle with L determined by the “primitive” Kähler class in $H^2((\mathcal{H}_a, \Delta_{\mathbf{m}}), \mathbb{Q})$, namely

$$(26) \quad \frac{c_1^{orb}(\mathcal{H}_a, \Delta_{\mathbf{m}_2})}{\mathcal{J}_{\mathbf{v}_2}} = \frac{1}{m_2v_2^0v_2^\infty} \frac{((2m_2v_2^0 + a)v_2^\infty x_1 + (v_2^0 + v_2^\infty)x_2)}{\mathcal{J}_{\mathbf{v}_2}}.$$

We make note of the fact that $Y^{p,q} \rightarrow (\mathcal{H}_a, \Delta_{\mathbf{m}_2})$ is an $m_2v_2^0v_2^\infty$ -fold covering map.

Now as in Section 3.5 of [BTF16] we want to describe the base orbifold $B_{1, \mathbf{v}, \mathbf{w}}$ of the S^1 orbibundle generated by the locally free action of the quasi-regular Reeb vector field $\xi_{\mathbf{m}_3}$ of the SE structure on the join $Y^{p,q} \star_{\ell} S_{\mathbf{w}}^3$. We see that the analysis of Section 3.5 of [BTF16] goes through verbatim through Remark 3.7 with $M = Y^{p,q}$ and $N = (\mathcal{H}_a, \Delta_{\mathbf{m}_2})$. In particular, from Lemma 3.6 of [BTF16] we obtain the base orbifold

$B_{\ell, \mathbf{v}, \mathbf{w}} \approx (B_{\ell, 1, \mathbf{w}'}, \Delta)$ with

$$\Delta = \left(1 - \frac{1}{m_3^0}\right) D_{v_3} + \left(1 - \frac{1}{m_3^\infty}\right) D_{u_3}, \quad \mathbf{w}' = (v_3^\infty w_1, v_3^0 w_2),$$

and from Theorem 3.8 of [BTF16] we have

$$(27) \quad \mathbf{m}_3 = m_3(v_3^0, v_3^\infty), \quad s = \gcd(|w_1 v_3^\infty - w_2 v_3^0|, l_2), \quad l_2 = s m_3, \quad n = \frac{l_1}{s} (w_1 v_3^\infty - w_2 v_3^0).$$

Then from the proof of Theorem 3.8 we see that the quotient is the total space of the projective orbibundle $\mathbb{P}(\mathbb{1} \oplus L^n)$ over $(\mathcal{H}_a, \Delta_{\mathbf{m}_2})$ whose invariant divisors are generally branch divisors of an orbifold. But then using Lemma 1.1 this is precisely a stage 3 Bott orbifold $(M_3(a, b, c), \Delta_{\mathbf{m}})$ for some b, c and where a is given in Equations (22). The fact that $Y^{p,q}$ as a join has the form of Equation (20) with the standard regular Sasakian structure on the first S^3 implies that the ramification indices $\mathbf{m}_1 = (1, 1)$.

It remains to check the equations for b and c . For this we make use of an orbifold version of Proposition 1.5 in [BCTF18]. Explicitly, we have

Lemma 2.4. *The pullback of $c_1(L^n)$ of the orbifold line bundle L^n to $M_3(a, b, c)$ is $b x_1 + c x_2$ where n is given in Equations (27).*

We know that L^n is the n th tensor product of the line orbibundle L which is determined by the Kähler class

$$(28) \quad \frac{c_1^{orb}(\mathcal{H}_a, \Delta_{\mathbf{m}_2})}{\mathcal{J}_{\mathbf{v}_2}} = \frac{((2m_2 v_2^0 + a)v_2^\infty x_1 + (v_2^0 + v_2^\infty)x_2)}{m_2 v_2^0 v_2^\infty \mathcal{J}_{\mathbf{v}_2}}.$$

Now the projection $p : (M_3(a, b, c), \Delta) \rightarrow (\mathcal{H}_a, \Delta_{\mathbf{m}_2})$ is a $m_2 v_2^0 v_2^\infty$ -fold covering map. So pulling back we have $c_1(L^n) = n c_1(L) = n p^* \left(\frac{c_1^{orb}(\mathcal{H}_a, \Delta_{\mathbf{m}_2})}{\mathcal{J}_{\mathbf{v}_2}} \right)$. The equations for b and c then follow by equating coefficients in this and in Lemma 2.4. \square

Remark 2.5. The Poincaré dual to $c_1^{orb}(M_3(a, b, c), \Delta_{\mathbf{m}})$ is a \mathbb{Q} -divisor on $M_3(a, b, c)$ which is ample when $c_1^{orb}(M_3(a, b, c), \Delta_{\mathbf{m}})$ is positive. Such a class gives a polarization to the orbifold $(M_3(a, b, c), \Delta_{\mathbf{m}})$, and a \mathbb{T}^3 invariant orbifold 2-form representing $c_1^{orb}(M_3(a, b, c), \Delta_{\mathbf{m}})$ gives an orbifold Kähler metric $g_{a,b,c,\mathbf{m}}$ on $M_3(a, b, c)$.

Remark 2.6. Note that the real cohomology class $c_1^{orb}(M_3(a, b, c), \Delta_{\mathbf{m}})$ makes perfect sense for $m_j^0, m_j^\infty \in \mathbb{R}^+$, and we denote the set of all such classes by \mathfrak{BL}_A . In this case $(M_3(a, b, c), \Delta_{\mathbf{m}})$ can be understood as having Kähler metrics with conical singularities along the corresponding \mathbb{R} -divisors D_{v_i} and D_{u_i} with cone angle $\frac{2\pi}{m_i^0}$ and $\frac{2\pi}{m_i^\infty}$, respectively [Don12, CDS15]. By this we mean that there is a Kähler metric ω which is smooth on $M_n(a, b, c) \setminus \Delta_{\mathbf{m}}$ which extends to $M_n(a, b, c)$ as a closed positive $(1, 1)$ current satisfying certain uniformity requirements. See Definition 1.3 of [CDS15] for the precise statement. For arbitrary \mathbf{m} we say that $c_1^{orb}(M_3(a, b, c), \Delta_{\mathbf{m}})$ represents a cone singularity along the divisor $\Delta_{\mathbf{m}}$. The rational entries in the interval $(0, 1)$ are related to the so-called ‘ramifolds’ [RT11]. As in this reference we shall also assume hereafter that $m_j^0, m_j^\infty \in (0, \infty)$. We denote by \mathfrak{LF}_A the subset of \mathfrak{BL}_A whose Kähler metrics are log Fano and have cone singularities along the divisor $\Delta_{\mathbf{m}}$. Then the natural map from \mathfrak{LF}_A to the

Kähler cone $\mathcal{K}(M_3(a, b, c))$ is surjective, and the conclusion of Lemma 1.2 holds for all $(M_3(a, b, c), \Delta_{\mathbf{m}}) \in \mathfrak{L}\mathfrak{F}_A$.

Remark 2.7. We consider the action of the affine monoid $\mathfrak{M}(\mathbb{R})$ on $(\mathbb{R}^+)^6$ defined by the affine linear map

$$(29) \quad (m_j^0, m_j^\infty) \mapsto (\lambda_j^0 m_j^0 + a_j^0, \lambda_j^\infty m_j^\infty + a_j^\infty) = (\tilde{m}_j^0, \tilde{m}_j^\infty)$$

with $1 \leq \lambda_j < \infty$ and $0 \leq a_j < \infty$. Restricting to the positive integers gives an action of the submonoid $\mathfrak{M}(\mathbb{Z})$ on $(\mathbb{Z}^+)^6$. One easily checks that this induces an action of $\mathfrak{M}(\mathbb{R})$ on $\mathfrak{B}\mathfrak{L}$ that leaves the subset $\mathfrak{L}\mathfrak{F}_A$ invariant for all $\lambda_j^0, \lambda_j^\infty \in [1, \infty)$ and $a_i^0, a_j^\infty \in [0, \infty)$ sending $c_1^{orb}(M_3(a, b, c), \Delta_{\mathbf{m}})$ to $c_1^{orb}(M_3(a, b, c), \Delta_{\tilde{\mathbf{m}}})$.

3. THE TOPOLOGY OF $M^7 = Y^{p,q} \star_{l_1, l_2} S_{\mathbf{w}}^3$

It is important to remember that generally the topology of a join depends on the choice of Sasakian structure (through its Reeb vector field) of each factor. We assume that (p, q) are relatively prime with $1 \leq q < p$ and that l_1, l_2, w_1, w_2 are chosen such that M^7 is smooth. We show first that our Sasaki 7-manifolds M^7 have the rational cohomology of the 2-fold connected sum $(S^2 \times S^5) \# (S^2 \times S^5)$. The integer cohomology groups are only distinguished by torsion in H^4 . Moreover, the torsion depends on the choice of quasi-regular Sasakian structures on $Y^{p,q}$ and S^3 . For the most generality we choose arbitrary quasi-regular Sasakian structure in the so-called \mathbf{w} subcone of the Sasaki cones for both $Y^{p,q}$ and S^3 (of course, the \mathbf{w} cone of S^3 is its entire Sasaki cone).

First we note that any quotient of a quasi-regular Reeb vector field $\xi_{\mathbf{m}}$ in the \mathbf{w} cone of $Y^{p,q}$ has the form of a Hirzebruch orbifold $(\mathcal{H}_a, \Delta_{\mathbf{m}})$. Moreover, $Y^{p,q}$ is itself the join $S^3 \star_{l_1, p} S_{\mathbf{w}}^3$ with $\mathbf{w} = (\frac{p+q}{l_1}, \frac{p-q}{l_1})$. Here $l_1 = 2$ if p, q are both odd, and $l_1 = 1$ if p, q have opposite parities. In any case the relation with the ramification indices is $\mathbf{m} = m(v^0, v^\infty)$ with v^0, v^∞ coprime and $m = p$.

The purpose of this section is to prove

Theorem 3.1. *Let $Y^{p,q}$ have a quasi-regular Sasakian structure with Reeb vector field $\xi_{\mathbf{m}}$. Then the 7-manifolds $M^7 = Y^{p,q} \star_{l_1, l_2} S_{\mathbf{w}}^3$ have the rational cohomology of the connected sum $(S^2 \times S^5) \# (S^2 \times S^5)$. Furthermore, the only torsion that occurs is $H^4(M^7, \mathbb{Z}) \approx \mathbb{Z}_{v^0 v^\infty m^2 l_2^2} \oplus \mathbb{Z}_{w_1 w_2 l_1^2}$.*

We begin with some lemmas.

Lemma 3.2. *The 7-manifolds $M^7 = Y^{p,q} \star_{l_1, l_2} S_{\mathbf{w}}^3$ satisfy the following conditions:*

- (1) $H_1(M^7, \mathbb{Z}) = \pi_1(M^7) = 0$,
- (2) $\pi_2(M^7) = \mathbb{Z}^2$,
- (3) $H^2(M^7, \mathbb{Z}) = H_2(M^7, \mathbb{Z}) = \mathbb{Z}^2$,
- (4) $H^3(M^7, \mathbb{Z})$ is torsion free.
- (5) $b_3(M^7) = b_4(M^7)$ is even.

Proof. From the long exact homotopy sequence for the fibration

$$(30) \quad \mathbb{T}^2 \longrightarrow S^3 \times S^3 \times S^3 \longrightarrow M^7$$

we conclude that M^7 is simply connected and that $\pi_2(M^7) = \mathbb{Z}^2$. Thus, by Hurewicz $H_2(M^7, \mathbb{Z}) = \mathbb{Z}^2$ which implies (4) by universal coefficients, and then by item (1)

$H^2(M^7, \mathbb{Z}) = \mathbb{Z}^2$. Item (5) follows from Poincaré duality and the fact that M^7 admits a Sasakian structure. □

Actually we have

Lemma 3.3. $H^3(M^7, \mathbb{Z}) = 0$.

Proof. First by (4) of Lemma 3.2 $H^3(M^7, \mathbb{Z})$ is torsion free, so it suffices to work with \mathbb{Q} coefficients. Since M^7 is simply connected and is an S^1 orbibundle over a stage 3 Bott orbifold $(M_3(a, b, c), \Delta_{\mathbf{m}})$ we can apply the Leray-Serre Theorem with \mathbb{Q} coefficients. The differential $d_2 : E_2^{0,1} \rightarrow E_2^{2,0} = H^2(M^7, \mathbb{Q})$ sends the class α of the fiber S^1 to the Kähler class $c_1x_1 + c_2x_2 + c_3x_3$ where $c_i \in \mathbb{Q}^+$. So by naturality we have $d_2(\alpha \otimes x_i) = (c_1x_1 + c_2x_2 + c_3x_3)x_i$. Suppose there would exist a class $w = w_1x_1 + w_2x_2 + w_3x_3 \in E_2^{2,0}$ such that $d_2(\alpha \otimes w) = 0$. Then the 3-class $\alpha \otimes w$ would survive to the limit giving a nonzero element in $H^3(M^7, \mathbb{Q})$ by the Leray-Serre Theorem. Now the cohomology ring of $M_3(a, b, c)$ is [CMS10]

$$\mathbb{Z}[x_1, x_2, x_3]/(x_1^2, x_2(ax_1 + x_2), x_3(bx_1 + cx_2 + x_3)).$$

So computing the d_2 differential we have

$$\begin{aligned} 0 &= d_2(\alpha \otimes w) = d_2(\alpha) \otimes w = (c_1x_1 + c_2x_2 + c_3x_3)(w_1x_1 + w_2x_2 + w_3x_3) \\ &= (c_1w_2 + c_2w_1 - c_2w_2a)x_1x_2 + (c_1w_3 + c_3w_1 - c_3w_3b)x_1x_3 + (c_2w_3 + c_3w_2 - c_3w_3c)x_2x_3 \end{aligned}$$

which gives

$$(31) \quad \begin{pmatrix} c_2 & c_1 - c_2a & 0 \\ c_3 & 0 & c_1 - c_3b \\ 0 & c_3 & c_2 - c_3c \end{pmatrix} \begin{pmatrix} w_1 \\ w_2 \\ w_3 \end{pmatrix} = 0.$$

Since the coefficients c_1, c_2, c_3 are all positive, the rank of the matrix

$$C = \begin{pmatrix} c_2 & c_1 - c_2a & 0 \\ c_3 & 0 & c_1 - c_3b \\ 0 & c_3 & c_2 - c_3c \end{pmatrix}$$

is either 3 or 2 which can be seen by putting C in Jordan canonical form. If the rank of C were 2 there would be exactly one solution to Equation (31) which is an element of $E_2^{1,2}$ and since $E_2^{3,0} = H^3(M_3(a, b, c), \mathbb{Q}) = 0$, there would be precisely one generator in $H^3(M^7, \mathbb{Q})$ which contradicts the fact that $b_3(M^7)$ is even. □

Proof of Theorem 3.1: It follows from Lemmas 3.2 and 3.3 and the fact that the homology (cohomology) of a connected sum is the direct sum of the homology (cohomology) of the two factors that M^7 has the rational homology (cohomology) of the connected sum $(S^2 \times S^5) \# (S^2 \times S^5)$. By Poincaré duality and Lemma 3.2 it follows that $H^5(M^7, \mathbb{Z})$ and $H^6(M^7, \mathbb{Z})$ have no torsion. So it suffices to compute the torsion in $H^4(M^7, \mathbb{Z})$. First

Lemma 3.4. *We have*

$$H_{orb}^r((\mathcal{H}_a, \Delta_{\mathbf{m}}), \mathbb{Z}) = \begin{cases} \mathbb{Z} & \text{if } r = 0 \\ \mathbb{Z}^2 & \text{if } r = 2 \\ \mathbb{Z} \oplus \mathbb{Z}_{m^0} \oplus \mathbb{Z}_{m^\infty} & \text{if } r = 4 \\ \mathbb{Z}_{m^0} \oplus \mathbb{Z}_{m^\infty} & \text{if } r = 6, 8, \dots \\ 0 & \text{if } r \text{ is odd.} \end{cases}$$

Proof. The Leray sheaf of the map

$$p : \mathbf{B}(\mathcal{H}_a, \Delta_{\mathbf{m}}) \longrightarrow \mathcal{H}_a$$

is the derived functor sheaf $R^s p\mathbb{Z}$, that is, the sheaf associated to the presheaf $U \mapsto H^s(p^{-1}(U), \mathbb{Z})$. For $s > 0$ the stalks of $R^s p\mathbb{Z}$ at points of U vanish if U lies in the regular locus of $(\mathcal{H}_a, \Delta_{\mathbf{m}})$ which is the complement of the union of the zero e_0 and infinity e_∞ sections of the natural projection $\mathcal{H}_a \rightarrow \mathbb{C}\mathbb{P}^1$. However, at points of e_0 and e_∞ the fibers of p are (up to homotopy) the Eilenberg-MacLane spaces $K(\mathbb{Z}_{m^0}, 1)$ and $K(\mathbb{Z}_{m^\infty}, 1)$, respectively. So at points of $e_0(e_\infty)$ the stalks are the group cohomology $H^s(\mathbb{Z}_{m^0}, \mathbb{Z})(H^s(\mathbb{Z}_{m^\infty}, \mathbb{Z}))$. This is \mathbb{Z} for $s = 0$ and $\mathbb{Z}_{m^0}(\mathbb{Z}_{m^\infty})$ at points of $e_0(e_\infty)$ when $s > 0$ is even; it vanishes when s is odd. The E_2 term of the Leray spectral sequence of the map p is

$$E_2^{r,s} = H^r(\mathcal{H}_a, R^s p\mathbb{Z})$$

and by Leray's theorem this converges to the orbifold cohomology $H_{orb}^{r+s}((\mathcal{H}_a, \Delta_{\mathbf{m}}), \mathbb{Z})$. Now $E_2^{r,0} = H^r(\mathcal{H}_a, \mathbb{Z})$ and $E_2^{r,s} = 0$ for r or s odd. For $r = 0$ since $R^s p\mathbb{Z}$ has its support in the orbifold singular locus $e_0 \cup e_\infty$, the only continuous section of $R^s p\mathbb{Z}$ is the 0 section which implies that $E_2^{0,s} = 0$ for all s . Now we have $E_2^{2r,2s} = 0$ for $r > 1$ and

$$E_2^{2,2s} = H^2(\mathcal{H}_a, R^{2s} p) = H^2(e_0, \mathbb{Z}_{m^0}) \oplus H^2(e_\infty, \mathbb{Z}_{m^\infty}) = \mathbb{Z}_{m^0} \oplus \mathbb{Z}_{m^\infty}.$$

One easily sees this spectral sequence collapses whose limit is the orbifold cohomology $H_{orb}^r((\mathcal{H}_a, \Delta_{\mathbf{m}}), \mathbb{Z})$ which implies the result. \square

To continue the proof of Theorem 3.1, as in [BTF16, BTF15], we use the commutative diagram of fibrations

$$(32) \quad \begin{array}{ccccc} Y^{p,q} \times S_{\mathbf{w}}^3 & \longrightarrow & M_{l_1, l_2, \mathbf{w}} & \longrightarrow & \mathbf{B}S^1 \\ \downarrow = & & \downarrow & & \downarrow \psi \\ Y^{p,q} \times S_{\mathbf{w}}^3 & \longrightarrow & \mathbf{B}(\mathcal{H}_a, \Delta_{\mathbf{m}}) \times \mathbf{BCP}^1[\mathbf{w}] & \longrightarrow & \mathbf{B}S^1 \times \mathbf{B}S^1. \end{array}$$

Here $\mathbf{B}G$ is the classifying space of a group G or Haefliger's classifying space [Hae84] of an orbifold if G is an orbifold. The lower exact fibration is a product of fibrations. We denote the orientation classes of $Y^{p,q} \times S^3 = S^2 \times S^3 \times S^3$ by α, β, γ , respectively. As in 3.2.2 of [BTF15] we have $d_4(\gamma) = w_1 w_2 s_2^2$. For the fibration

$$Y^{p,q} \longrightarrow \mathbf{B}(\mathcal{H}_a, \Delta_{\mathbf{m}}) \longrightarrow \mathbf{B}S^1$$

we have $d_4(\beta) = m^2 v^0 v^\infty s_1^2$. The map ψ in Diagram (32) is induced by the map $e^{i\theta} \mapsto (e^{il_2\theta}, e^{-il_1\theta})$, so $\psi^* s_1 = l_2 s$ and $\psi^* s_2 = -l_1 s$ the result follows by the commutativity of Diagram (32). \square

Remark 3.5. Although the case $(p, q) = (1, 0)$ does not fit directly into the $Y^{p,q}$ scheme of [GMSW04], it can, nevertheless, be identified with the homogeneous Sasaki-Einstein structure on $S^2 \times S^3$. Then if we take $(v^0, v^\infty) = (w_1, w_2) = (l_1, l_2) = (1, 1)$ and $m = 1$ we obtain the homogeneous Sasaki-Einstein structure on an S^1 bundle over $\mathbb{C}P^1 \times \mathbb{C}P^1 \times \mathbb{C}P^1$ which is a 7-manifold with the integral cohomology of the 2-fold connected sum $2(S^2 \times S^5)$. See Sections 11.1.1 and 11.4.2 of [BG08] for details. The general S_w^3 join with $Y^{1,0}$ was treated in Section 3.2.2 of [BTF15].

4. A GENERALIZED ORBIFOLD CALABI CONSTRUCTION

We now discuss a family of explicit examples or orbifold Kähler-Einstein metrics that we may view as arising from a special case of the generalized Calabi construction as presented in [ACGTF04, Section 2.5] and further discussed in [ACGTF11, Section 2.3] and [BCTF18, Section 5.1] - here generalized to allowing certain mild orbifold singularities.

The base of the construction will in this case be a Kähler-Einstein Hirzebruch orbifold $(\mathcal{H}_a, \Delta_{\mathbf{m}_2})$. This is the Kähler quotient of the quasi-regular Sasaki-Einstein $Y^{p,q}$ examples produced in [GMSW04]). Now $Y^{p,q}$ may be viewed as the total space of a S^1 principal orbi-bundle, P , over $(\mathcal{H}_a, \Delta_{\mathbf{m}_2})$ defined by the class $\frac{c_1^{orb}(\mathcal{H}_a, \Delta_{\mathbf{m}_2})}{J_{\mathbf{v}_2}} = \frac{(2m_2v_2^0+a)v_2^\infty x_1 + (v_2^0+v_2^\infty)x_2}{m_2v_2^0v_2^\infty J_{\mathbf{v}_2}}$, where the notation is as in Lemma 2.2. In particular, the index of $(\mathcal{H}_a, \Delta_{\mathbf{m}_2})$ is $J_{\mathbf{v}_2} = \gcd(2m_2v_2^0 + a)v_2^\infty, v_2^0 + v_2^\infty)$, where $(m_2^0, m_2^\infty) = m_2(v_2^0, v_2^\infty)$ and v_2^0, v_2^∞ are coprime. Let g_{base} denote the Kähler-Einstein orbifold metric whose Kähler form, ω_{base} , satisfies that

$$\left[\frac{\omega_{base}}{2\pi} \right] = \frac{(2m_2v_2^0 + a)v_2^\infty x_1 + (v_2^0 + v_2^\infty)x_2}{m_2v_2^0v_2^\infty J_{\mathbf{v}_2}}.$$

As we saw in Example 6.8 of [BTF16], the metric g_{base} is explicit and "admissible" in the sense of [ACGTF08]. Note that the Ricci form of g_{base} is given by $\rho_{base} = J_{\mathbf{v}_2}\omega_{base}$.

We consider the generalized Calabi construction of orbifold Kähler metrics on the bundle $P \times_{S^1} \mathbb{C}P^1_{m_3^0, m_3^\infty} \rightarrow (\mathcal{H}_a, \Delta_{\mathbf{m}_2})$. This may also be viewed as an admissible construction - extended to mild orbifold cases.

Definition 4.1. Generalized orbifold Calabi data for our purposes.

- (1) A log pair $(\mathcal{H}_a, \Delta_{\mathbf{m}_2})$ with Kähler-Einstein structure $(\omega_{base}, g_{base})$ such that $\left[\frac{\omega_{base}}{2\pi} \right] = \frac{(2m_2v_2^0+a)v_2^\infty x_1 + (v_2^0+v_2^\infty)x_2}{m_2v_2^0v_2^\infty J_{\mathbf{v}_2}}$.
- (2) The weighted projective line $(\mathbb{C}P^1_{m_3^0, m_3^\infty} = \mathbb{C}P^1_{v_3^0, v_3^\infty} / \mathbb{Z}_{m_3}, g_{\mathbf{m}_3}, \omega_{\mathbf{m}_3})$ with rational Delzant polytope $[-1, 1] \subseteq \mathbb{R}^*$ and momentum map $\mathfrak{z}: \mathbb{C}P^1_{m_3^0, m_3^\infty} \rightarrow [-1, 1]$. Here $(m_3^0, m_3^\infty) = m_3(v_3^0, v_3^\infty)$ and v_3^0, v_3^∞ are coprime.
- (3) A principal S^1 orbi-bundle, $P_n \rightarrow (\mathcal{H}_a, \Delta_{\mathbf{m}_2})$, with a principal connection of curvature $n\omega_{base} \in \Omega^{1,1}((\mathcal{H}_a, \Delta_{\mathbf{m}_2}), \mathbb{R})$, where S^1 acts on $\mathbb{C}P^1_{m_3^0, m_3^\infty}$, $n \in \mathbb{Z} \setminus \{0\}$, and $\gcd(n, m_3) = 1$. Note that $n \in \text{span}_{\mathbb{Z}}\{v_3^0, v_3^\infty\}$ (since v_3^0, v_3^∞ are coprime), so $m_3n \in \text{span}_{\mathbb{Z}}\{m_3^0, m_3^\infty\}$.
- (4) A constant $0 < |r_3| < 1$ of same sign as n [ensuring that the $(1, 1)$ -form $(1/r_3 + \mathfrak{z})n\omega_{base}$ is positive for $\mathfrak{z} \in [-1, 1]$].

From this data we may define the orbifold

$$M_3 = P_n \times_{S^1} \mathbb{C}P_{m_3^0, m_3^\infty}^1 = \mathring{M}_3 \times_{\mathbb{C}^*} \mathbb{C}P_{m_3^0, m_3^\infty}^1 \rightarrow (\mathcal{H}_a, \Delta_{\mathbf{m}_2}),$$

where $\mathring{M}_3 = P_n \times_{S^1} (\mathfrak{z}^{-1}(-1, 1))$. Since the curvature 2-form of P_n has type $(1, 1)$, \mathring{M}_3 is a holomorphic principal \mathbb{C}^* bundle with connection $\theta \in \Omega^1(\mathring{M}_3, \mathbb{R})$ and M_3 is a complex orbifold.

On \mathring{M}_3 we define Kähler structures of the form

$$(33) \quad \begin{aligned} g &= (1/r_3 + \mathfrak{z})n g_{base} + \frac{1}{\Theta(\mathfrak{z})} d\mathfrak{z}^2 + \Theta(\mathfrak{z})\theta^2 \\ \omega &= (1/r_3 + \mathfrak{z})n \omega_{base} + d\mathfrak{z} \wedge \theta \\ d\theta &= n \omega_{base}, \end{aligned}$$

where $\frac{1}{\Theta(\mathfrak{z})} = \frac{d^2 U}{d\mathfrak{z}^2}$ and U is the symplectic potential [Gui94] of the chosen toric Kähler structure $g_{\mathbf{m}_3}$ on $\mathbb{C}P_{m_3^0, m_3^\infty}^1$.

The *generalized Calabi construction* arises from seeing (33) as a blueprint for the construction of various orbifold Kähler metrics on M_3 by choosing various smooth functions $\Theta(\mathfrak{z})$ on $(-1, 1)$ satisfying that

- [boundary values] the following endpoint conditions are satisfied
- $$(34) \quad \Theta(\pm 1) = 0 \quad \text{and} \quad \Theta'(-1) = 2/m_3^\infty \quad \text{and} \quad \Theta'(1) = -2/m_3^0;$$
- [positivity] the function $\Theta(\mathfrak{z})$ is positive for $\mathfrak{z} \in (-1, 1)$.

Then (33) extends to an orbifold Kähler metric on $(M(a, b, c), \Delta_{\mathbf{m}})$, where $b = n \frac{(2m_2 v_2^0 + a)v_2^\infty}{\mathcal{J}_{v_2}}$, $c = n \frac{v_2^0 + v_2^\infty}{\mathcal{J}_{v_2}}$, and $\mathbf{m} = (1, 1, m_2^0, m_2^\infty, m_3^0, m_3^\infty)$. Metrics constructed this way are called *compatible Kähler metrics* with *compatible Kähler classes* parametrized by r_3 . From [ACG06] (and specifically to our notation from [BTF16, Proposition 5.4]) we have that the compatible metric defined by $\Theta(\mathfrak{z})$ is Kähler-Einstein exactly when

$$(35) \quad 2r_3 \mathcal{J}_{v_2} / n = (1 + r_3) / m_3^\infty + (1 - r_3) / m_3^0$$

and

$$(36) \quad \int_{-1}^1 ((1 - \mathfrak{z}) / m_3^\infty - (1 + \mathfrak{z}) / m_3^0) (1 + r_3 \mathfrak{z})^2 d\mathfrak{z} = 0.$$

5. EXPLICIT SASAKI-EINSTEIN METRICS

We now look more closely for explicit Sasaki-Einstein examples arising from the join from Theorem 2.3. The arguments in Sections 6.1 and 6.2 of [BTF16] (see specifically page 1053) carry through so that we have an adapted version of Theorem 1.4 of [BTF16]:

Theorem 5.1. *Consider the join $M^7 = Y^{p,q} \star_{l_1, l_2} S_{\mathbf{w}}^3$ where $Y^{p,q}$ has a quasi-regular Sasaki-Einstein structure with quotient Hirzebruch orbifold $(\mathcal{H}_a, \Delta_{v_2})$ with $a > 0$ and l_1, l_2 are given by*

$$l_1 = \frac{\mathcal{J}_{v_2}}{\gcd(|\mathbf{w}|, \mathcal{J}_{v_2})}, \quad l_2 = \frac{|\mathbf{w}|}{\gcd(|\mathbf{w}|, \mathcal{J}_{v_2})}.$$

Then for each vector $\mathbf{w} = (w_1, w_2) \in \mathbb{Z}^+ \times \mathbb{Z}^+$ with relatively prime components satisfying $w_1 > w_2$ there exist a Reeb vector field $\xi_{\mathbf{v}_3}$ in the 2-dimensional \mathbf{w} -Sasaki cone on M^7 such that the corresponding Sasakian structure is Sasaki-Einstein.

Specifically, using equation (59) from [BTF16], we know that if the ray defined by co-prime (v_3^0, v_3^∞) is quasi-regular, then we ought to look at the Kähler class determined by $r_3 = \frac{w_1 v_3^\infty - w_2 v_3^0}{w_1 v_3^\infty + w_2 v_3^0}$. Now with this r_3 , and the above choice of (l_1, l_2) , (35) is automatically solved and (36) becomes (similarly to (68) in [BTF16])

$$(37) \quad \int_{-1}^1 ((v_3^0 - v_3^\infty) - (v_3^0 + v_3^\infty)\mathfrak{z}) ((w_1 v_3^\infty + w_2 v_3^0) + (w_1 v_3^\infty - w_2 v_3^0)\mathfrak{z})^2 d\mathfrak{z} = 0.$$

This equation defines a priori a quasi-regular Sasaki η -Einstein ray (and thus, up to transverse homothety, a Sasaki-Einstein structure), but, by the same arguments as in Section 6.1 of [BTF16], any solution $(v_3^0, v_3^\infty) \in \mathbb{R}^+ \times \mathbb{R}^+$ of (37) defines a Sasaki η -Einstein ray in the \mathbf{w} -cone, which is irregular unless $v_3^\infty/v_3^0 \in \mathbb{Q}$. This gives

Corollary 5.2. *For any pair of relatively prime positive integers (w_1, w_2) satisfying $w_1 > w_2$, we obtain an SE metric by solving the cubic equation*

$$(38) \quad 3w_2 k^3 + (2w_2 - w_1)k^2 - (2w_1 - w_2)k - 3w_1 = 0.$$

The ray of the Reeb vector field of the SE metric is then given by

$$(39) \quad \frac{v_3^\infty}{v_3^0} = \frac{3 + 2k + k^2}{(1 + 2k + 3k^2)} = k \frac{w_2}{w_1}.$$

The SE metric is quasi-regular when k is a rational root of (38) and irregular when it is an irrational root.

It follows from the analysis in Section 6.2 of [BTF16] that the positive real root k lies in the open interval $(1, \infty)$.

5.1. Quasi-regular examples. Changing our point of view we see that for any $k \in (1, +\infty) \cap \mathbb{Q}$ we can choose (w_1, w_2) such that

$$(40) \quad w_2/w_1 = \frac{3 + 2k + k^2}{k(1 + 2k + 3k^2)}$$

and then the Sasaki-Einstein metric from Theorem 5.1 is quasi-regular with ray \mathbf{v}_3 defined by Equation (39). We now give some examples.

Example 5.3. Here we give an example that builds on a bouquet of Sasaki cones from Example¹ 6.8 of [BTF16]. Corollary 5.5 of [BP14] describes the well known $Y^{p,q}$ structures [GMSW04] on $S^2 \times S^3$ as a $|\phi(p)|$ -bouquet of Sasaki cones where $|\phi(p)|$ denotes the order of the Euler phi function $\phi(p)$. Let us consider the example when $p = 13$. Since 13 is prime the bouquet consists of $p - 1 = 12$ Sasaki cones labeled by the 12 positive integers $1 \leq q < 13$ and as such contains 12 Sasaki-Einstein metrics. However, in order to construct SE metrics on our 7-manifolds, we need the SE structure on $Y^{p,q}$ to be quasi-regular. It is easy to check that for the bouquet $\bigcup_{q=1}^{12} \{Y^{13,q}\}$ the only values of q where the SE metric is quasi-regular is for $q = 7, 8$, all the other SE metrics in the bouquet are irregular. Let us look at these two cases a bit closer.

¹Example 6.8 in this reference has a small typo, namely v_2^0 and v_2^∞ got switched.

In the case $Y^{13,8}$ we have $a = 70$, $v_2^0 = 7$, $v_2^\infty = 5$, $m_2 = 13$, so $m_2^0 = 91$ and $m_2^\infty = 65$. This gives us that $\mathcal{J}_{\mathbf{v}_2} = 12$. If we put $k = 2$ in the quasi-regular SE prescription above given by Equation (38), we see that $(w_1, w_2) = (34, 11)$, so $l_1 = 4$ and $l_2 = 15$, and $(v_3^0, v_3^\infty) = (17, 11)$. Now we calculate that $s = \gcd(|w_1 v_3^\infty - w_2 v_3^0|, l_2) = 1$ so $n = l_1 \frac{w_1 v_3^\infty - w_2 v_3^0}{s} = 748$, and $(m_3^0, m_3^\infty) = 15(17, 11)$. Then the quotient is the Bott orbifold given by the log pair $(M(a, b, c), \Delta_{\mathbf{m}})$, where

$$\begin{aligned} a &= 70, \\ b &= n\hat{b} = n \frac{(2m_2 v_2^0 + a)v_2^\infty}{\mathcal{J}_{\mathbf{v}_2}} = 78540, \\ c &= n\hat{c} = n \frac{v_2^0 + v_2^\infty}{\mathcal{J}_{\mathbf{v}_2}} = 748, \\ \mathbf{m} &= (1, 1, 91, 65, 255, 165). \end{aligned}$$

In the case where $p = 7$ in $Y^{p,q}$ we have $a = 36$, $v_2^0 = 4$, $v_2^\infty = 3$, $m_2 = 13$, so $m_2^0 = 52$ and $m_2^\infty = 39$. This gives us that $\mathcal{J}_{\mathbf{v}_2} = 7$. Again we put $k = 2$ in Equation (38) which gives $(w_1, w_2) = (34, 11)$ and $(v_3^0, v_3^\infty) = (17, 11)$ respectively. Now $l_1 = 7$ and $l_2 = 45$, and $s = \gcd(|w_1 v_3^\infty - w_2 v_3^0|, l_2) = 1$ so $n = l_1 \frac{w_1 v_3^\infty - w_2 v_3^0}{s} = 1309$, and $(m_3^0, m_3^\infty) = 45(17, 11)$. Then the quotient is the Bott orbifold given by the log pair $(M(a, b, c), \Delta_{\mathbf{m}})$, where

$$\begin{aligned} a &= 36, \\ b &= n\hat{b} = n \frac{(2m_2 v_2^0 + a)v_2^\infty}{\mathcal{J}_{\mathbf{v}_2}} = 78540, \\ c &= n\hat{c} = n \frac{v_2^0 + v_2^\infty}{\mathcal{J}_{\mathbf{v}_2}} = 1309, \\ \mathbf{m} &= (1, 1, 52, 39, 765, 495). \end{aligned}$$

One can easily check from the torsion in Theorem 3.1 that the two SE 7-manifolds $Y^{13,8} \star_{4,15} S_{(34,11)}^3$ and $Y^{13,7} \star_{7,45} S_{(34,11)}^3$ are not homotopy equivalent. For both of these 7-manifolds the Reeb field that gives the SE metric is quasi-regular. Moreover, they are both induced from the same Reeb ray, namely $\{\xi_{a(17,11)}\}_{a>0}$ of the same $S_{(34,11)}^3$, and b is the same for the quotient orbifolds.

For each choice of the rational number $k > 1$ we obtain a pair of quasi-regular SE 7-manifolds induced by the $S_{\mathbf{w}}^3$ join and its Reeb field $\xi_{\mathbf{v}_3}$ where \mathbf{w} and \mathbf{v}_3 are determined by Equations (38) and (39).

Example 5.4. Here we give a 1-parameter family of smooth quasi-regular examples. First, let $k_2 \in \mathbb{Z}^{\geq 0}$ be given (using the subscript "2" to indicate that this is a choice at the second stage). Then from [GMSW04]) we get a quasi-regular Sasaki-Einstein $Y^{p,q}$ example by choosing

$$p = 12k_2^2 + 18k_2 + 7 \quad \text{and} \quad q = 12k_2^2 + 16k_2 + 5.$$

It is not hard to check that $\gcd(p, q) = 1$ and $\gcd(p + q, p - q) = 2$. Accordingly we recognize from (20) that

$$Y^{p,q} = S^3 \star_{l,p} S_{\frac{p+q}{2}, \frac{p-q}{2}}^3 = S^3 \star_{2, (12k_2^2 + 18k_2 + 7)} S_{(2+3k_2)(3+4k_2), (1+k_2)}^3$$

and that the quotient Hirzebruch orbifold of the quasi-regular Sasaki-Einstein metric is $(\mathcal{H}_a, \Delta_{\mathbf{m}_2})$ where, using (21) and (22),

$$\mathbf{m}_2 = (m_2^0, m_2^\infty) = (12k_2^2 + 18k_2 + 7)((3 + 4k_2), 2(1 + k_2)),$$

$$m_2 = 12k_2^2 + 18k_2 + 7,$$

$$a = 6(1 + k_2)(1 + 2k_2)(3 + 4k_2).$$

We can calculate from Lemma 2.2 that $\mathcal{J}_{\mathbf{v}_2} = 5 + 6k_2$.

Now choosing $k = 3$ in (38) and (39), we have, from Theorem 5.1, a quasi-regular Sasaki-Einstein structure on $M^7 = Y^{p,q} \star_{l_1, l_2} S_{\mathbf{w}}^3$ with quotient log pair $(M_3(a, b, c), \Delta_{\mathbf{m}})$ given by Theorem 2.3. Indeed, we have

$$a = 6(1 + k_2)(1 + 2k_2)(3 + 4k_2)$$

$$b = 4(1 + k_2)(2 + 3k_2)(3 + 4k_2)n$$

$$c = n$$

$$\mathbf{m} = (m_1^0, m_1^\infty, m_2^0, m_2^\infty, m_3^0, m_3^\infty)$$

$$(m_1^0, m_1^\infty) = (1, 1)$$

$$(m_2^0, m_2^\infty) = (12k_2^2 + 18k_2 + 7)((3 + 4k_2), 2(1 + k_2))$$

$$(m_3^0, m_3^\infty) = m_3(v_3^0, v_3^\infty),$$

where

$$n = l_1 \frac{102}{\gcd(102, l_2)}$$

$$m_3 = \frac{l_2}{\gcd(102, l_2)}$$

$$l_1 = \frac{\mathcal{J}_{\mathbf{v}_2}}{\gcd(20, 5 + 6k_2)}$$

$$l_2 = \frac{20}{\gcd(20, 5 + 6k_2)}$$

$$(w_1, w_2) = (17, 3)$$

$$(v_3^0, v_3^\infty) = (17, 9).$$

Using Lemma 2.1, we know that the corresponding Sasaki structure is a smooth manifold if and only if

$$\begin{aligned} \gcd(l_2(12k_2^2 + 18k_2 + 7)(3 + 4k_2), 17l_1) &= 1 \\ \gcd(l_2((12k_2^2 + 18k_2 + 7)(3 + 4k_2)), 3l_1) &= 1 \\ \gcd(l_2(12k_2^2 + 18k_2 + 7)2((1 + k_2)), 17l_1) &= 1 \\ \gcd(l_2(12k_2^2 + 18k_2 + 7)2(1 + k_2), 3l_1) &= 1 \end{aligned}$$

In order to get smooth examples, let us now assume $k_2 = 255t + 10$ with $t \in \mathbb{Z}^+$. Then we have

$$J_{\mathbf{v}_2} = 5(306t + 13)$$

$$l_1 = 306t + 13$$

$$l_2 = 4$$

$$n = 51(306t + 13)$$

$$m_3 = 2$$

and then the corresponding Sasaki structure is a smooth manifold if and only if

$$\begin{aligned} \gcd(4(1387 + 65790t + 780300t^2)(1020t + 43), 17(306t + 13)) &= 1 \\ \gcd(4(1387 + 65790t + 780300t^2)(1020t + 43), 3(306t + 13)) &= 1 \\ \gcd(8(1387 + 65790t + 780300t^2)(255t + 11), 17(306t + 13)) &= 1 \\ \gcd(8(1387 + 65790t + 780300t^2)(255t + 11), 3(306t + 13)) &= 1. \end{aligned}$$

if and only if

$$\begin{aligned} \gcd((1387 + 65790t + 780300t^2)(1020t + 43), (306t + 13)) &= 1 \\ \gcd((1387 + 65790t + 780300t^2)(255t + 11), (306t + 13)) &= 1. \end{aligned}$$

Since $\forall t \in \mathbb{Z}^+$,

$$6(255t + 11) - 5(306t + 13) = 1$$

$$10(306t + 13) - 3(1020t + 43) = 1$$

$$3(780300t^2 + 65790t + 1387) - (7650t + 320)(306t + 13) = 1,$$

we have that this is always satisfied. Note that with $k_2 = 255t + 10$ we get

$$a = 6(255t + 11)(510t + 21)(1020t + 43)$$

$$b = 204(255t + 11)(765t + 32)(1020t + 43)(306t + 13)$$

$$c = 51(306t + 13)$$

$$\mathbf{m} = (m_1^0, m_1^\infty, m_2^0, m_2^\infty, m_3^0, m_3^\infty)$$

$$(m_1^0, m_1^\infty) = (1, 1)$$

$$(m_2^0, m_2^\infty) = (1387 + 65790t + 780300t^2)(1020t + 43, 2(255t + 11))$$

$$(m_3^0, m_3^\infty) = 2(17, 9).$$

Finally, the p and q in $Y^{p,q}$ are here given by

$$p = 780300t^2 + 65790t + 1387$$

and

$$q = 15(170t + 7)(306t + 13),$$

so the smooth Sasaki Einstein structures live on

$$Y^{780300t^2+65790t+1387,15(170t+7)(306t+13)} \star_{306t+13,4} S_{17,3}^3.$$

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