

SPECTRAL CLUSTER ESTIMATES FOR METRICS OF SOBOLEV REGULARITY

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ABSTRACT. We investigate spectral cluster estimates for compact manifolds equipped with a Riemannian metric whose regularity is determined by its inclusion in a Sobolev space of sufficiently high order. The problem is reduced to obtaining L^p estimates for the wave equation which are shown by employing wave packet techniques.

1. INTRODUCTION

Spectral cluster estimates refer to a family of inequalities on eigenfunctions of the Laplacian on a compact Riemannian manifold (M, g) . It is well known that the Laplace-Beltrami operator associated to (M, g) , denoted by Δ_g , possesses a sequence of eigenfunctions $\{\phi_j\}_{j=0}^\infty$ that form an orthonormal basis for $L^2(M, dV_g)$, dV_g denoting the Riemannian measure induced by g . We take the convention that the eigenfunctions are ordered so that

$$\Delta_g \phi_j = -\lambda_j^2 \phi_j \text{ with } 0 = \lambda_0 \leq \lambda_1 \leq \dots \leq \lambda_j \leq \lambda_{j+1} \leq \dots$$

We define the spectral projection operator Π_λ for any $\lambda \geq 0$ as

$$\Pi_\lambda f = \sum_{\lambda_j \in [\lambda, \lambda+1]} \langle f, \phi_j \rangle \phi_j.$$

Spectral cluster estimates then state that

$$(1.1) \quad \|\Pi_\lambda f\|_{L^q(M)} \leq C \lambda^{\delta(q)} \|f\|_{L^2(M)} \quad \delta(q) = \begin{cases} \frac{n-1}{2}(\frac{1}{2} - \frac{1}{q}) & 2 \leq q \leq q_n \\ n(\frac{1}{2} - \frac{1}{q}) - \frac{1}{2} & q_n \leq q \leq \infty \end{cases}$$

with $q_n = \frac{2(n+1)}{n-1}$ and C independent of λ in both cases. These inequalities were originally developed by C.D. Sogge in [12] for compact Riemannian manifolds with C^∞ metrics.

It is also of interest to investigate whether or not such estimates hold for manifolds equipped with rougher metrics of limited differentiability. The first work in this direction came in [10], where H. Smith and C.D. Sogge provided examples of $C^{1,\alpha}$ metrics for each $\alpha \in [0, 1)$ where the spectral cluster estimates fail to hold in the range $2 \leq q \leq q_n$. Other examples are in [11]. The natural question then was whether or not such estimates should hold for $C^{1,1}$ metrics. Indeed, the geodesic flow is well-behaved for such metrics, suggesting that such estimates should hold.

Recently, in [8], H. Smith proved these estimates for $C^{1,1}$ metrics by employing wave equation methods. The main idea here is that by examining the effect of the

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wave group on a spectral cluster $\Pi_\lambda f$ one can reduce the estimates (1.1) to certain mixed L^q estimates on solutions to the wave equation. These inequalities, often called “squarefunction estimates” state that if $u(x, t) : [-t_0, t_0] \times M \rightarrow \mathbb{C}$ denotes a solution to the wave equation and $q \geq q_n$, then

$$(1.2) \quad \|u\|_{L_x^q(M; L_t^2[-t_0, t_0])} \leq C(\|(u(0, \cdot), \partial_t u(0, \cdot))\|_{H^{\delta(q)} \times H^{\delta(q)-1}} + \|Pu\|_{L_t^1([-t_0, t_0]; H^{\delta(q)-1}}).$$

Here $P = \partial_t^2 - \Delta_g$ denotes the wave operator and H^s denotes the L^2 Sobolev space of order s over M with $\delta(q)$ as in (1.1). Squarefunction estimates were first developed for C^∞ metrics by Mockenhaupt, Seeger, and Sogge in [6]. They were then shown in the case of $C^{1,1}$ metrics in the aforementioned work of Smith [8].

To handle metrics rougher than $C^{1,1}$, H. Smith then developed a family of weaker spectral cluster estimates for $C^{1,\alpha}$ metrics (see [9]). These estimates bound the operator norm of Π_λ as a map from $L^2(M)$ to $L^q(M)$ by a power of λ higher than that appearing in (1.1). Specifically, if g is a $C^{1,\alpha}$ metric with $0 \leq \alpha \leq 1$ and $\sigma = \frac{1-\alpha}{3+\alpha}$, then his results show that

$$(1.3) \quad \|\Pi_\lambda\|_{L^2(M) \rightarrow L^q(M)} \leq C\lambda^{\delta(q)(1+\sigma)} \quad 2 \leq q \leq \frac{2(n+1)}{n-1}$$

$$(1.4) \quad \|\Pi_\lambda\|_{L^2(M) \rightarrow L^q(M)} \leq C\lambda^{\delta(q) + \frac{\sigma}{q}} \quad \frac{2(n+1)}{n-1} \leq q \leq \infty.$$

The examples of [10], [11] show that such estimates are sharp for q in the range $2 \leq q \leq q_n$. We also remark that a recent work of Koch-Smith-Tataru [5] has developed a version of these estimates for Hölder C^s metrics below Lipschitz regularity which are sharp for $2 \leq q \leq q_n$ and $q = \infty$.

These estimates are equivalent to a weaker version of the squarefunction estimates for wave equations in $C^{1,\alpha}$ metrics which state that

$$(1.5) \quad \|u\|_{L_x^q(M; L_t^2[-t_0, t_0])} \leq C(\|(u(0, \cdot), \partial_t u(0, \cdot))\|_{H^{\delta(q) + \frac{\sigma}{q}} \times H^{\delta(q) + \frac{\sigma}{q} - 1}} + \|Pu\|_{L_t^1([-t_0, t_0]; H^{\delta(q) + \frac{\sigma}{q} - 1}}).$$

Hence there is a gain in the number of derivatives required on the right hand side. These estimates follow from (1.4), but can also be shown directly by adapting the “truncation/rescaling” methods found in the work of D. Tataru (see [13], [14]) and in Bahouri-Chemin (see [1], [2]). Their results demonstrated that a weaker version of the related Strichartz estimates hold, which control the $L_t^p L_x^q$ -norm of solutions to wave equations with C^s coefficients when $0 < s \leq 2$.

The purpose of this paper is to consider spectral cluster estimates in the case where the metric tensor g lies in an L^r Sobolev space of sufficiently high order with $1 < r < \infty$. Specifically, we seek to establish the estimates (1.1) for metrics in $L^{r,\kappa}$ (the L^r Sobolev space of order κ) for $\kappa > \frac{n-1}{r} + 2$. In addition, we use truncation/rescaling methods show that the estimates (1.3), (1.4) hold when the metric lies in $L^{r,\kappa} \cap Lip$ with $\kappa = \frac{n-1}{r} + 1 + \alpha$ for $\alpha \in (0, 1)$ and $\sigma = \frac{1-\alpha}{3+\alpha}$ determined by the α appearing in the definition of κ .

Clearly, the property $g \in L^{r,\kappa}(M)$, $\kappa > \frac{n-1}{r} + 2$ does not imply that the metric tensor is $C^{1,1}$. However, it is enough differentiability to ensure that the geodesic flow exists and is well-behaved. These properties were used in a recent work by the

author (see [3]) to show that in dimensions $n = 2, 3$, the Strichartz estimates hold for solutions to wave equations whose coefficients possess this degree of Sobolev regularity. Strichartz estimates are similar to the squarefunction inequalities in (1.2) in that they control mixed L^p norms of the solution. However, they are distinct in that they depend largely on the dispersive effect of the solution operator, whereas squarefunction inequalities also rely on the decay of the operator away from the light cone. Nevertheless, a high degree of control of the geodesic flow is important in establishing both sets of inequalities. At the moment, we refer the reader to section 1 and section 4 of [3] for a detailed discussion of these ideas in the context of Strichartz estimates, though they will also appear below.

Metrics possessing this degree of Sobolev regularity arise naturally if one looks only at L^r bounds on the Ricci tensor. If one assumes a mild degree of regularity on the metric g , it can be shown that the manifold can be covered by harmonic coordinate charts, that is, coordinates (y_1, \dots, y_n) such that $\Delta_g y_l = 0$ for each coordinate function y_l . See Proposition 9.1 and the surrounding discussion in Chapter 3 of [17] for an approach that requires only that $g \in C^s(M)$. In such coordinates, it can be shown that the metric satisfies a quasilinear elliptic system of equations

$$\partial_j g^{jk}(y) \partial_k g_{lm}(y) + Q_{lm}(g, \nabla g) = Ric_{lm}$$

where $Q_{lm}(g, \nabla g)$ is a quadratic form in ∇g with coefficients determined by g . By applying elliptic regularity results for such systems (see Corollary 12B.5 in Chapter 14 of [15] and Proposition 10.1 of [17]), we have that if $p \geq \frac{n}{2}$, $1 < r < \infty$, and $\kappa \geq 2$, then

$$g_{jk} \in H^1(B_1) \cap C^0(B_1), Ric_{lm} \in L^{r, \kappa-2}(B_1) \cap L^p(B_1) \Rightarrow g_{jk} \in L^{r, \kappa}(B_{\frac{1}{2}})$$

with B_r denoting a ball of radius r in the coordinate chart. See Proposition 4.10 in Chapter 14 of [15] for another result of this type which requires less regularity from the Ricci tensor. Hence our results will provide a weaker form of the spectral cluster estimates (1.5) for “ L^r -pinched” metrics whose Ricci tensor (defined in the sense of distributions) satisfies $Ric \in L^r(M)$, $r \geq \frac{n}{2}$ and α satisfying $\frac{n-1}{r} + 1 + \alpha = 2$.

Finite speed of propagation allows one to work in coordinate charts and easily reduce the estimates (1.2), (1.5) to establishing squarefunction inequalities over \mathbb{R}^n for solutions to the Cauchy problem with $\rho(x) = \sqrt{\det(g_{ij})}$, $a_{ij}(x) = g^{ij}(x)\rho(x)$

$$(1.6) \quad P(x, D)u(x, t) := \rho(x) \partial_t^2 u(x, t) - \sum_{1 \leq i, j \leq n} \partial_{x_i} (a_{ij}(x) \partial_{x_j} u(x, t)) = F(x, t)$$

$$u(x, 0) = f, \quad \partial_t u(x, 0) = g.$$

One of the main challenges in proving estimates for wave equations in this context lies in constructing a suitable parametrix for the equation. When the coefficients of P are smooth, there are asymptotic methods available for this purpose, but are inapplicable in the non-smooth setting. An effective substitute is to use wave packets in constructing a parametrix. In our context, wave packets are approximate solutions to the wave equation that are highly localized in space and in frequency. They can be used to construct an effective parametrix for the wave equation by writing general solutions as a superposition of these packets. In [8], H. Smith employed a wave packet transform, similar to the Fourier-Bros-Iaglonitzer (FBI) transform used by Tataru in [13], [14] and the wave packets of Cordoba-Fefferman [4]. This

provides an effective way of representing solutions to equations with time-dependent coefficients satisfying

$$(1.7) \quad \partial_{x,t}^\beta \rho(x,t), \partial_{x,t}^\beta a_{ij}(x,t) \in L_t^1([-t_0, t_0], L_x^\infty(\mathbb{R}^n)) \quad |\beta| \leq 2.$$

These techniques are not immediately applicable when coefficients possess the Sobolev regularity outlined above. Indeed, the ability to control the L^∞ -norm of the second order derivatives of the coefficients over all spatial variables plays a crucial role in Smith's construction. Therefore, a different perspective on the problem is needed. Suppose for now that $\|\rho - 1\|_{L^\infty}, \|a_{ij} - \delta_{ij}\|_{L^\infty} \leq \frac{1}{4}$ so that P is in some sense a perturbation of the constant coefficient wave operator. Write the principal symbol of $-P(x, \xi, \tau)$ as

$$a_{11}(x)\xi_1^2 + 2 \sum_{j=2}^n a_{1j}(x)\xi_1\xi_j - \left(\rho(x)\tau^2 - \sum_{2 \leq i, j \leq n} a_{ij}(x)\xi_i\xi_j \right).$$

We now see that P can be viewed as a operator that is hyperbolic in x_1 when restricted to space-time frequencies lying in a set of the form

$$(1.8) \quad \{(\xi, \tau) : |\tau| \gg |(\xi_2, \dots, \xi_n)|\}$$

since this means that the principal symbol of P is a quadratic in ξ_1 with 2 distinct real roots. But if the coefficients lie in $L^{r,\kappa}$ with $\kappa > \frac{n-1}{r} + 2$, Sobolev embedding implies that for any T in the orthogonal group $O(n)$ and any multi-index $|\beta| \leq 2$

$$(1.9) \quad \partial_x^\beta ((a_{ij} \circ T)(x) - \delta_{ij}) \in L_{x_1}^r(\mathbb{R}; L_{x'}^\infty(\mathbb{R}^{n-1})) \subset L_{x_1,loc}^1(\mathbb{R}; L_{x'}^\infty(\mathbb{R}^{n-1})).$$

where $x = (x_1, x')$. The latter set here is the space of locally integrable functions in the x_1 variable with values in the Banach space of L^∞ functions over the x' variable. When T is the identity, this is akin to (1.7) for operators hyperbolic in x_1 . This suggests that squarefunction estimates should hold for solutions to the wave equation that are localized to frequencies in the set (1.8). It should then be possible to prove squarefunction estimates for arbitrary solutions to the wave equation by representing them as a sum of solutions localized in frequency to regions of this type. Indeed, it was this perspective that allowed for a proof of the Strichartz estimates in [3] and will now form the backbone of our arguments here.

Our main approach will be to combine the frequency localization arguments of [3] with the wave packet transform techniques of [8]. We first reduce the problem to showing estimates on components of our solution that are localized in frequency to dyadic regions contained in (1.8). Once this is done, we show that these components are also solutions to a first-order pseudodifferential equation that is hyperbolic in x_1 . With these reductions in place, we are then able to use a wave packet transform as in Smith [8] to represent these localized functions as a continuous superposition of wave packets parameterized by x_1 . It is then a matter of showing that this parametrix yields the desired $L_x^q L_t^2$ estimates.

We will also prove squarefunction estimates for coefficients in the space $L^{r,\kappa}$, $\kappa = \frac{n-1}{r} + 1 + \alpha$. Such coefficients satisfy the property that for any T in the orthogonal group $O(n)$ and any multi-index $|\beta| \leq 1$

$$(1.10) \quad \partial_x^\beta ((a_{ij} \circ T)(x) - \delta_{ij}) \in L_{x_1,loc}^1(\mathbb{R}; C_{x'}^{1-|\beta|,\alpha}(\mathbb{R}^{n-1})).$$

This property will allow us to use elements of the previously mentioned “truncation/rescaling” approach to prove inequalities of the form (1.5).

The main results of this paper are the following:

Theorem 1.1. *Assume that $\rho - 1, a_{ij} - \delta_{ij} \in L^{r,\kappa}(\mathbb{R}^n)$, with $\kappa > \frac{n-1}{r} + 2$ and that $0 \leq \delta(q) \leq \frac{3}{2}$. Then there exists C depending only on the dimension n and on the Sobolev norms of the coefficients such that*

$$\|u\|_{L^q_x(\mathbb{R}^n; L^2_t[-t_0, t_0])} \leq C(\|f\|_{H^{\delta(q)}} + \|g\|_{H^{\delta(q)-1}} + \|F\|_{L^1_t([-t_0, t_0]; H^{\delta(q)-1}(\mathbb{R}^n)))$$

for $q \geq q_n$ and any solution u to the Cauchy problem (1.6).

Theorem 1.2. *Assume that $\rho - 1, a_{ij} - \delta_{ij} \in L^{r,\kappa}(\mathbb{R}^n) \cap Lip(\mathbb{R}^n)$, with $\kappa = \frac{n-1}{r} + 1 + \alpha$. Suppose also that $\sigma = \frac{1-\alpha}{3+\alpha}$, $q \geq q_n$ and $\delta(q) + \frac{\sigma}{q} \leq 1$. Then there exists C depending only on the dimension n and on the Sobolev/Lipschitz norms of the coefficients such that*

$$\|u\|_{L^q_x(\mathbb{R}^n; L^2_t[-t_0, t_0])} \leq C(\|f\|_{H^{\delta(q)+\frac{\sigma}{q}}} + \|g\|_{H^{\delta(q)+\frac{\sigma}{q}-1}} + \|F\|_{L^1_t([-t_0, t_0]; H^{\delta(q)+\frac{\sigma}{q}-1}(\mathbb{R}^n)))$$

for any solution u to the Cauchy problem (1.6).

Note that $\delta(q_n) + \frac{\sigma}{q_n} \leq 1$, yielding squarefunction estimates (1.2), (1.5) over manifolds for a range of indices that includes the critical index q_n . By the techniques of [8], this in turn proves spectral cluster estimates at the critical index for manifolds with metrics of the corresponding degree of Sobolev regularity. In [9], Smith used heat kernel methods to show the spectral cluster estimates (1.1) for $q = \infty$ with no gain in the power of λ when $g \in Lip(M)$. As a consequence we can interpolate Smith’s results with our spectral cluster estimates at the critical index to obtain the following:

Theorem 1.3. *Let M be C^∞ manifold equipped with a $L^{r,\kappa}(M) \cap Lip(M)$ metric g with $\kappa = \frac{n-1}{r} + 1 + \alpha$ where either $\alpha > 1$ or $0 < \alpha < 1$. The spectral projector Π_λ induced by Δ_g then satisfies, for $\sigma = \max(0, \frac{1-\alpha}{3+\alpha})$,*

$$(1.11) \quad \|\Pi_\lambda\|_{L^2(M) \rightarrow L^q(M)} \leq C\lambda^{\frac{n-1}{2}(\frac{1}{2}-\frac{1}{q})(1+\sigma)} \quad 2 \leq q \leq \frac{2(n+1)}{n-1}$$

$$(1.12) \quad \|\Pi_\lambda\|_{L^2(M) \rightarrow L^q(M)} \leq C\lambda^{n(\frac{1}{2}-\frac{1}{q})-\frac{1}{2}+\frac{\sigma}{q}} \quad \frac{2(n+1)}{n-1} \leq q \leq \infty.$$

We remark that this theorem yields an improvement over what would be obtained by embedding the metric into a space C^s and applying the current results in [9], [5]. This is because it allows us to control the operator norm of Π_λ by a smaller power of λ than would otherwise be needed.

In light of recent examples of Koch-Smith-Tataru [5], the assumption that g is a Lipschitz metric is crucial for large values of q in Theorem 1.3. In their work, for each $0 < s < 1$, they produce functions $\rho_\lambda, \psi_\lambda$ over \mathbb{R}^n such that

$$\Delta\psi_\lambda + \lambda^2\rho_\lambda\psi_\lambda = 0, \quad \frac{\|\psi_\lambda\|_{L^\infty}}{\|\psi_\lambda\|_{L^2}} \approx \lambda^{\frac{n-s}{2}}, \quad \rho_\lambda = 1 + \lambda^{-s}q_1(\lambda|x|) + \lambda^{-2s}q_2(\lambda|x|)$$

where Δ is the standard Laplacian on Euclidean space and q_1, q_2 are smooth functions independent of λ . It is also seen that $\|\rho_\lambda - 1\|_{C^s(M)}$ is uniformly bounded in λ . The ψ_λ are exponentially concentrated to balls of radius λ^{s-1} allowing one to lift a sequence of these examples to a compact manifold and produce a compact

(M, g) with $g \in C^s(M)$ a Hölder metric for which the best possible L^∞ spectral cluster estimate is

$$\|\Pi_\lambda\|_{L^2(M) \rightarrow L^\infty(M)} \leq C\lambda^{\frac{n-s}{2}}.$$

Computing the $L^{r,\kappa}$ Sobolev norm of $\rho_\lambda - 1$ yields

$$\|\rho_\lambda - 1\|_{L^{r,\kappa}} = \|\lambda^{-s}q_1(\lambda|\cdot|) + \lambda^{-2s}q_2(\lambda|\cdot|)\|_{L^{r,\kappa}} \leq C\lambda^{\kappa - \frac{n}{r} - s}.$$

Thus when $0 < \kappa - \frac{n}{r} < 1$ and $L^{r,\kappa}$ does not embed into a space of Lipschitz functions, setting $s = \kappa - \frac{n}{r}$ above also provides an example of an $L^{r,\kappa}$ metric for which the L^∞ spectral cluster estimate $\|\Pi_\lambda\|_{L^2(M) \rightarrow L^\infty(M)} \leq C\lambda^{\frac{n-1}{2}}$ fails.

Using an argument of H. Smith [9], we can also see that we must have

$$\|\Pi_\lambda\|_{L^2(M) \rightarrow L^q(M)} \gtrsim \lambda^{\frac{n-s}{2} - \frac{n}{q}}.$$

This uses the heat kernel methods alluded to above that were employed by Smith to show the $L^2 \rightarrow L^\infty$ spectral cluster estimate for Lipschitz metrics. Consider $H_\lambda = \exp(-\lambda^{-2}\Delta)$, the heat evolution operator at time λ^{-2} . By Theorem 6.3 of Saloff-Coste [7] we have a pointwise bound over the integral kernel h_λ of H_λ

$$|h_\lambda(x, y)| \leq C\lambda^n \exp(-c\lambda^2 d_g(x, y)^2)$$

with $d_g(x, y)$ denoting Riemannian distance. By Young's inequality we now have

$$\lambda^{\frac{n-s}{2}} \leq C \frac{\|\psi_\lambda\|_{L^\infty}}{\|\psi_\lambda\|_{L^2}} \leq C' \frac{\lambda^{\frac{n}{q}} \|H_\lambda^{-1}\psi_\lambda\|_{L^q}}{\|\psi_\lambda\|_{L^2}} \leq C'' \lambda^{\frac{n}{q}} \exp\left(\frac{\lambda^2}{\lambda^2}\right) \frac{\|\psi_\lambda\|_{L^q}}{\|\psi_\lambda\|_{L^2}}$$

which shows the claim. This now shows that (1.12) fails for a range of q in an interval of the form $(c, \infty]$. For small r , values of α can be chosen so that $c < \frac{2(n+1)}{n-1}$, and hence (1.11) fails for certain q as well. Hence the assumption that $g \in Lip(M)$ is indeed crucial.

Unfortunately, our methods do not handle metrics in $L^{r,\kappa}$ when $\kappa = \frac{n-1}{r} + 1 + \alpha$ and $\alpha = 0$ or 1 . In both cases, Sobolev embedding fails to provide the necessary control over the derivatives of the metric needed to prove the estimates. Indeed, if $\alpha = 1$, (1.9) fails to hold for multi-indices $|\beta| = 2$, precluding us from showing estimates of the form (1.2). When $\alpha = 0$, (1.10) fails to hold, meaning that the truncation/rescaling methods here do not yield (1.5) with $\sigma = \frac{1}{3}$. In other words when $\alpha = 0$ we are unable to provide any results, and when $\alpha = 1$, the best we can provide is squarefunction estimates with an arbitrarily small loss of derivatives.

We also comment that the results of this paper yield Strichartz estimates in dimensions $n \geq 4$. Indeed, the reductions in section 2 below are similar to those in [3] and the estimates (4.11), (5.8), (5.10) can be combined to establish the key dispersive estimates. However, we shall not pursue this issue in detail.

Before proceeding in the proofs of Theorem 1.1, 1.2, we will make a few additional assumptions about our coefficients and our solution. We will always assume that $a_{ij} - \delta_{ij}, \rho - 1 \in Lip(\mathbb{R}^n)$ and are compactly supported with

$$(1.13) \quad \|a_{ij} - \delta_{ij}\|_{Lip}, \|\rho - 1\|_{Lip} < c_0$$

for some sufficiently small but fixed $c_0 > 0$. In addition, when $\kappa > \frac{n-1}{r} + 2$ we assume that

$$\sup_{T \in O(n)} \sup_{|\beta| \leq 2} \|\partial_x^\beta (a_{ij}(T \cdot) - \delta_{ij})\|_{L_{x_1}^r L_{x'}^\infty} < c_0.$$

Lastly, we assume that $t_0 = 1$ and that $u(\cdot, t)$ is supported in the cube $(-\frac{3}{2}, \frac{3}{2})^n$ for all $t \in (-\frac{5}{4}, \frac{5}{4})$. The extra hypotheses are superfluous as the general case can be reduced to these assumptions. This involves showing that we can restrict our attention to solutions supported in small sets and then rescaling the problem. See [13], pp. 394-5, for a suitable approach that also works in the context of squarefunction estimates. We will also assume that u solves the homogeneous Cauchy problem with $F \equiv 0$ in (1.6). By Duhamel's principle it is not difficult to see that the general case reduces to such an assumption.

The paper is organized as follows, section 2 parallels the preliminaries in [3], reducing the problem to showing $L_x^q L_t^2$ estimates on solutions to the wave equation that are dyadically localized in frequency to regions contained in a set of the form (1.8). In section 3, we then reduce the problem further, showing that it will suffice to find estimates on solutions to a first order equation that is hyperbolic in x_1 with coefficients satisfying (1.9). section 4 then uses a continuous wave packet transform to construct a parametrix for the first order equation. It is then shown in section 5 that this parametrix yields an effective method for proving the desired estimates, culminating in a proof of Theorem 1.1.

Notation. In what follows d will denote the gradient operator which maps scalar functions to vector fields and vector fields to matrix functions in the natural way. d_x will denote that spatial derivatives are taken and $d_{x,t}$ denotes the full space-time gradient. Also, given a Lebesgue measurable rectangle $U \times V \subset \mathbb{R}^n \times \mathbb{R}$, $L_x^q L_t^p(U \times V)$ and $L_x^q L_t^p$ will abbreviate the Banach space $L^q(U; L^p(V))$, the latter being used when $U \times V = \mathbb{R}^{n+1}$. Similarly, $L_t^p H^z$ abbreviates $L^p(I; H^z)$ when it is I is an interval understood from the context of the argument.

The expression $X \lesssim Y$ means that $X \leq CY$ for some C depending only on n and on the Lipschitz/ $L^{r,\kappa}$ norms of the coefficients. Additionally, $X \ll Y$ means that $X \leq C(c_0)Y$ where $C(c_0)$ denotes a constant depending on the choice of c_0 that can be made arbitrarily small by choosing c_0 small.

Lastly, $b_j(x)$ will always denote the function $b_j(x) = \sum_{i=1}^n (\partial_i a_{ij})(x)$ so that we can write our wave operator in non-divergence form with a first order term

$$P(x, D) = \rho(x) \partial_t^2 - \sum_{1 \leq i, j \leq n} a_{ij}(x) \partial_i \partial_j - \sum_{j=1}^n b_j(x) \partial_j.$$

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2. PRELIMINARY REDUCTIONS

We begin by laying out the first set of reductions in showing the squarefunction estimates. This is very similar the approach used to show Strichartz estimates

in [3]. The main idea here is to reduce the problem to showing estimates on the components of u frequency localized to dyadic regions contained in cones of the form $|\tau| \gg |(\xi_2, \dots, \xi_n)|$.

First, we comment on the energy estimates that can be established for our wave operator. We have that

$$\|v(\cdot, t)\|_{H^z} + \|\partial_t v(\cdot, t)\|_{H^{z-1}} \lesssim \|v(\cdot, 0)\|_{H^z} + \|\partial_t v(\cdot, 0)\|_{H^{z-1}} + \int_0^t \|(Pv)(\cdot, s)\|_{H^{z-1}} ds$$

provided that $0 \leq z \leq 1$ when $a_{ij} \in Lip(\mathbb{R}^n)$ or that $0 \leq z \leq \frac{3}{2}$ when $a_{ij} - \delta_{ij} \in L^{r, \kappa}(\mathbb{R}^n)$, $\kappa > \frac{n-1}{r} + 2$. To see this, first observe that the case $z = 1$ follows from a standard computation. When $\rho \equiv 1$, the $z = 0$ case follows by applying the $z = 1$ inequality to $\langle D \rangle^{-1} v$ ($\langle D \rangle$ denoting the Fourier multiplier with symbol $\langle \xi \rangle = (1 + |\xi|^2)^{\frac{1}{2}}$) along with an application of the commutator estimate (2.4) below. The $z = 0$ case for variable ρ follows by applying this estimate to the equation satisfied by the function $w(x, t) = \rho(x)v(x, t)$ and using the fact that

$$\|\partial_{x_j}(a_{ij}(\partial_{x_i} \rho^{-1}))w(\cdot, t)\|_{H^{-1}} \lesssim \|w(\cdot, t)\|_{L^2}.$$

The estimates when $0 \leq z \leq 1$ now follow by interpolation. To establish inequalities when $1 \leq z \leq \frac{3}{2}$ we use energy estimates for operators in non-divergence form. This provides the desired estimates as the additional regularity of the coefficients implies that multiplication by b_j is a bounded operation on H^{z-1} .

Before we begin localizing our solution in frequency, we make a small alteration of our solution so that it is compactly supported in time. Replace $u(x, t)$ by $\phi(t)u(x, t)$ where $\phi(t)$ is a smooth bump function equal to 1 on $[-1, 1]$ and supported in $(-\frac{5}{4}, \frac{5}{4})$. Since this defines $u(x, t)$ as a function over all of \mathbb{R}^{n+1} , it now makes sense to look at u as an element of a Sobolev space in $n + 1$ variables. Hence $H_{x,t}^w$ will denote the L^2 Sobolev space of order w on \mathbb{R}^{n+1} in the (x, t) variables and H^w will denote the Sobolev space on \mathbb{R}^n in only the x variables. This modification of u means that it is no longer a solution to the homogeneous problem. However, energy estimates developed above allow us to control $[P, \phi]$ so that we have the following inequality for all $t \in \mathbb{R}$

$$(2.1) \quad \|u(\cdot, t)\|_{H^z} + \|\partial_t u(\cdot, t)\|_{H^{z-1}} + \|(Pu)(\cdot, t)\|_{H^{z-1}} \lesssim \|f\|_{H^z} + \|g\|_{H^{z-1}}$$

provided that $0 \leq z \leq 1$ when $a_{ij} \in Lip(\mathbb{R}^n)$ or that $0 \leq z \leq \frac{3}{2}$ when $a_{ij} - \delta_{ij} \in L^{r, \kappa}(\mathbb{R}^n)$, $\kappa > \frac{n-1}{r} + 2$.

To localize in frequency we first take a sequence of Littlewood-Paley cutoffs $\{\beta_k\}_{k \geq 0}$, so that $\beta_k(\xi)$ is a smooth function for $\xi \in \mathbb{R}^n$ with $\sum_k \beta_k \equiv 1$, $\text{supp}(\beta_k) \subset \{\xi \in \mathbb{R}^n : 2^{k-\frac{1}{2}} \leq |\xi| \leq 2^{k+\frac{3}{2}}\}$ when $k \neq 0$, and $\text{supp}(\beta_0) \subset B_1(0)$. We can also take the sequence so that $\beta_k(\cdot) = \beta(2^{-k}\cdot)$ for $k > 1$ for some function $\beta \in C_0^\infty(\mathbb{R}^n)$.

By multiplying β_0 by another smooth cutoff in τ , we can extend β_0 to be smooth cutoff supported in the unit ball in \mathbb{R}^{n+1} and one in a neighborhood of the origin in \mathbb{R}^{n+1} , that is, there exists $\bar{\beta}_0(\xi, \tau)$ with the above properties such that $\bar{\beta}_0(\xi, 0) = \beta_0(\xi)$. Treating $\bar{\beta}_0(D)$ as a Fourier multiplier on u in the (x, t) variables, $\bar{\beta}_0(D)$ is a smoothing operator and hence by Sobolev embedding and energy estimates we

obtain

$$\|\bar{\beta}_0(D)u\|_{L_x^q L_t^2([-1,1] \times \mathbb{R}^n)} \lesssim \|u\|_{L_t^2 H^{\delta(q)+\frac{\sigma}{q}}} \lesssim \|f\|_{H^{\delta(q)+\frac{\sigma}{q}}} + \|g\|_{H^{\delta(q)+\frac{\sigma}{q}-1}}.$$

Now take a cutoff $\Lambda^+ \in C^\infty(\mathbb{R}^{n+1} \setminus \{0\})$ homogeneous of degree 0 such that $\text{supp}(\Lambda^+) \subset \{(\xi, \tau) : 9^{-1}|\xi| \leq \tau \leq 9|\xi|\}$ and is identically 1 on the slightly smaller cone $\{(\xi, \tau) : 8^{-1}|\xi| \leq \tau \leq 8|\xi|\}$. Since we already have estimates on $\bar{\beta}_0 u$ we will replace Λ^+ by $\Lambda^+(I - \bar{\beta}_0)$ so that Λ^+ is smooth on all of \mathbb{R}^{n+1} and vanishes on in a neighborhood of the origin. Also, set $\Lambda^-(\xi, \tau) = \Lambda^+(\xi, -\tau)$. We now need to control the $L_x^q L_t^2$ norm of the following:

$$\Lambda^+(D)u, \Lambda^-(D)u, \text{ and } (I - \Lambda^+(D) - \Lambda^-(D) - \bar{\beta}_0(D))u.$$

The latter component can be estimated using elliptic regularity.

Proposition 2.1. $(I - \Lambda^+(D) - \Lambda^-(D) - \bar{\beta}_0(D))u$ satisfies the following estimate

$$\|(I - \Lambda^+(D) - \Lambda^-(D) - \bar{\beta}_0(D))u\|_{L_x^q L_t^2(\mathbb{R}^{n+1})} \lesssim \|f\|_{H^{\delta(q)}} + \|g\|_{H^{\delta(q)-1}}.$$

Proof. For $k \geq 1$, set $\alpha_k(\xi, \tau) = (I - \Lambda^+(\xi, \tau) - \Lambda^-(\xi, \tau))\beta_k(\xi, \tau)$. By Minkowski's inequality, Littlewood-Paley theory and (2.1), it suffices to show that

$$\sum_{k \geq 1} \|\alpha_k(D)u\|_{L_t^2 L_x^q}^2 \lesssim \|u\|_{L_t^2 H^{\delta(q)}}^2 + \|\partial_t u\|_{L_t^2 H^{\delta(q)-1}}^2 + \|Pu\|_{L_t^2 H^{\delta(q)-1}}^2.$$

First observe that by Sobolev embedding we have

$$\|\alpha_k u\|_{L_t^2 L_x^q} \lesssim 2^{k(\delta(q)+\frac{1}{2})} \|\alpha_k u\|_{L_{x,t}^2}$$

Set $a_{ij}^k = \sum_{l \leq \frac{k}{2}} \beta_l(D) a_{ij}$ and let $P^k = \rho^k \partial_t^2 - \sum_{ij} \partial_i(a_{ij}^k \partial_j)$. Observing that since $a_{ij}(x) \in Lip(\mathbb{R}^n)$ we always have the crude estimates $\|a_{ij}^k - a_{ij}\|_{L^\infty} \lesssim 2^{-\frac{k}{2}}$ and $\|d_x a_{ij}^k\| \lesssim 1$. By choosing c_0 sufficiently small in (1.13), P^k is also perturbation of the constant coefficient wave operator. Thus $P^k(x, D) \circ \alpha_k(D)$ is an elliptic operator with symbol estimates in $S_{1, \frac{1}{2}}^2$ uniformly bounded in k . Hence there exist symbols $q^k(x, \xi, \tau) \in S_{1, \frac{1}{2}}^{-2}(\mathbb{R}^{2n+2})$, $r^k(x, \xi, \tau) \in S_{1, \frac{1}{2}}^{-\frac{1}{2}}(\mathbb{R}^{2n+2})$ with

$$I = q^k(x, D) \circ P^k(x, D) \circ \alpha_k(x, D) + r^k(x, D),$$

I denoting the identity operator. Since $\alpha_k(x, D)u$, $P^k(x, D)(\alpha_k(x, D)u)$ are localized to frequencies $|(\xi, \tau)| \gtrsim 2^k$, we then conclude that

$$\|\alpha_k(x, D)u\|_{L_{x,t}^2} \lesssim 2^{-2k} \|P^k(x, D)(\alpha_k(x, D)u)\|_{L_{x,t}^2} + 2^{-\frac{k}{2}} \|\alpha_k u\|_{L_{x,t}^2}.$$

Next observe that since ρ, ρ^k are bounded from below we have

$$\|P^k(\alpha_k u)\|_{L_{x,t}^2} \lesssim \left\| \left(\frac{1}{\rho^k} P^k - \frac{1}{\rho} P \right) \alpha_k u \right\|_{L_{x,t}^2} + \|[P, \alpha_k]u\|_{L_{x,t}^2} + \|\alpha_k Pu\|_{L_{x,t}^2}.$$

Since $\|a_{ij}^k - a_{ij}\|_{L^\infty} \lesssim 2^{-\frac{k}{2}}$ we can conclude that

$$\left\| \left(\frac{1}{\rho^k} P^k - \frac{1}{\rho} P \right) \alpha_k u \right\|_{L_{x,t}^2} \lesssim 2^{-\frac{k}{2}} \|d_x^2 \alpha_k u\|_{L_{x,t}^2} + \|d_x \alpha_k u\|_{L_{x,t}^2} \lesssim 2^{\frac{3k}{2}} \|\alpha_k u\|_{L_{x,t}^2}$$

The inequality

$$\sum_k 2^{2k\delta(q)} (\|\alpha_k u\|_{L_{x,t}^2}^2 + 2^{-2k} \|\alpha_k Pu\|_{L_{x,t}^2}^2) \lesssim \|u\|_{L_t^2 H^{\delta(q)+\frac{\sigma}{q}}}^2 + \|Pu\|_{L_t^2 H^{\delta(q)+\frac{\sigma}{q}-1}}^2$$

is straightforward, so it suffices to control the sum over the terms involving the commutator $[P, \alpha_k]$.

The key idea here will be to apply Lemma 2.4 of [3] which implies that if $R_\theta = \sum_k e^{ik\theta} 2^{k(\delta(q) - \frac{3}{2})} \alpha_k$ and

$$(2.2) \quad \|[P, R_\theta]u\|_{L_{x,t}^2} \lesssim \|\partial_t u\|_{L_t^2 H^{\delta(q) + \frac{\sigma}{q} - 1}} + \|u\|_{L_t^2 H^{\delta(q) + \frac{\sigma}{q}}}$$

then

$$\sum_k 2^{2k(\delta(q) - \frac{3}{2})} \|[P, \alpha_k]u\|_{L_{x,t}^2}^2 \lesssim \|\partial_t u\|_{L_t^2 H^{\delta(q) + \frac{\sigma}{q} - 1}}^2 + \|u\|_{L_t^2 H^{\delta(q) + \frac{\sigma}{q}}}^2.$$

In order to obtain (2.2), we use an important corollary of the Coifman-Meyer commutator theorem. Suppose R is a Fourier multiplier of order z and a is a Lipschitz function. Then the commutator of R with the multiplication operation $f \mapsto a \cdot f$ enjoys the following continuous mapping properties:

$$(2.3) \quad [a, R] : H^{z-1}(\mathbb{R}^{n+1}) \rightarrow L^2(\mathbb{R}^{n+1}) \quad 0 \leq z \leq 1$$

$$(2.4) \quad [a, R] : H^z(\mathbb{R}^{n+1}) \rightarrow H^1(\mathbb{R}^{n+1}) \quad -1 \leq z \leq 0$$

The Coifman-Meyer result handles the $z = 1$ case and the $z = 0, -1$ cases then follow easily. Interpolation then provides the result for the intermediate values of z . See Proposition 3.6B and the surrounding discussion in Section 3.6 of [16] for a more detailed discussion of this result.

We can now exploit the fact that P is a divergence form operator to obtain

$$(2.5) \quad \begin{aligned} \|[P, R_\theta]u\|_{L_{x,t}^2} &\lesssim \|[\rho, R_\theta]\partial_t u\|_{H_{x,t}^1} + \sum_{ij} \|[a_{ij}, R_\theta]\partial_j u\|_{H_{x,t}^1} \\ &\lesssim \|d_{x,t}u\|_{H_{x,t}^{\max(-1, \delta(q) - \frac{3}{2})}} \lesssim \|u\|_{L_t^2 H^{\delta(q)}} + \|\partial_t u\|_{L_t^2 H^{\delta(q) - 1}}. \end{aligned}$$

□

The squarefunction estimates for $\Lambda^+ u, \Lambda^- u$ will require further localization. Let $\Phi(\xi) \in C^\infty(\mathbb{R}^n \setminus \{0\})$ be a cutoff such that

$$(2.6) \quad \text{supp}(\Phi) \subset \{\xi : -\xi_1 > M|\xi'|\}$$

and identically 1 on a slightly smaller cone, where M is a sufficiently large constant to be determined. Set $\Gamma^\pm(\xi, \tau) = \Phi(\xi)\Lambda^\pm(\xi, \tau)$. By employing a family of similar cutoffs, we can represent $\Lambda^\pm u$ as a finite sum of functions supported in cones in ξ of comparable aperture, so it will suffice to show estimates for the localized $\Gamma^\pm(D)u$. Also, we will show them for $\Gamma^- u$ as estimates for $\Gamma^+ u$ will follow by an identical argument, so from now on we will suppress the ‘-’ in Γ^- and refer to it simply as Γ .

Set $\Gamma_k(\xi, \tau) = \beta_k(\xi)\Gamma(\xi, \tau)$. Also, let ψ be a smooth cutoff in x_1 that is identically 1 on a neighborhood of $[-\frac{3}{2}, \frac{3}{2}]$ and supported in $(-2, 2)$ so that $\psi(x_1)u(x, t) = u(x, t)$. We pause to mention that for any $x_1 \in \mathbb{R}$, the partial Fourier transform of $\psi\Gamma_k u$ in x', t is supported in the set

$$\text{supp}(\widehat{\psi\Gamma_k u})(x_1, \cdot) \subset \{(\xi', \tau) : -\tau \geq \widetilde{M}|\xi'|\} \cap \{(\xi', \tau) : 2^{k-5} \leq |(\xi', \tau)| \leq 2^{k+6}\}$$

where \widetilde{M} is a constant depending on M in (2.6) that can be made arbitrarily large by choosing M sufficiently large.

The key result of this section will be to show that the inequalities on Γu reduce to estimates on the $\psi\Gamma_k u$:

Proposition 2.2. *Suppose $\psi\Gamma_k u$ satisfies the estimates*

$$(2.7) \quad \begin{aligned} \|\psi\Gamma_k u\|_{L_x^q L_t^2([-1,1] \times \mathbb{R}^n)} &\lesssim 2^{k(\delta(q) + \frac{\sigma}{q})} \left(\|\psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2} + 2^{-k} \|d_{x,t}(\psi\Gamma_k u)\|_{L_{x_1}^\infty L_{x',t}^2} \right. \\ &\quad \left. + \|\Gamma_k u\|_{L_{x,t}^2} + 2^{-k} \|P(\psi\Gamma_k u)\|_{L_{x_1}^1 L_{x',t}^2} \right) \end{aligned}$$

uniformly in k . Then u satisfies the desired squarefunction estimates.

Note that $L_x^q L_t^2([-1,1] \times \mathbb{R}^n)$ abbreviates $L_{x_1,x'}^q([-1,1] \times \mathbb{R}^{n-1}; L_t^2(\mathbb{R}))$.

Proof. The proof here is very similar to that in Theorem 2.1 in [3]. We first observe that the projections of the $\Gamma_k(\xi, \tau)$ onto the τ axis have bounded overlap, meaning that Littlewood-Paley theory implies

$$\|\Gamma_k u\|_{L_x^q L_t^2} \approx \|\Gamma_k u\|_{L_x^q L_t^2 l_k^2} \leq \|\psi\Gamma_k u\|_{l_k^2 L_x^q L_t^2} + \|(1-\psi)\Gamma_k u\|_{l_k^2 L_x^q L_t^2}$$

Since $\Gamma_k u$ is given by convolution of u against a Schwartz function concentrated in a ball of radius 2^{-k} about the origin, we have that for any $N > 0$

$$\|(1-\psi)\Gamma_k u\|_{L_x^q L_t^2} \leq C_N 2^{-kN} \|u\|_{L_{x,t}^2}$$

with C_N independent of k . This implies that

$$\|(1-\psi)\Gamma_k u\|_{l_k^2 L_x^q L_t^2} \lesssim \|u\|_{L_t^2 H^{\delta(q)}}$$

and hence it will suffice to control the square sum over the $\psi\Gamma_k u$.

We now need to control the $L_{x_1}^\infty L_{x',t}^2$ norm of $d_{x,t}(\psi\Gamma_k u)$. For this we apply the flux estimates in Proposition 2.2 of [3] which imply that for any $r_1, y_1 \in [-2, 2]$ we have that

$$\|d_{x,t}(\psi\Gamma_k u)(y_1, \cdot)\|_{L_{x',t}^2} \lesssim \|d_{x,t}(\psi\Gamma_k u)(r_1, \cdot)\|_{L_{x',t}^2} + \int_{r_1}^{y_1} \|P(\psi\Gamma_k u)(z_1, \cdot)\|_{L_{x',t}^2} dz_1.$$

The key idea here is that since the $\psi\Gamma_k u$ are localized to a region where P behaves as an operator that is hyperbolic in x_1 , we can formulate energy estimates in the x_1 variable for such functions. Since $2^k \|\psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2} \lesssim \|\partial_t \psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2}$, we can choose $r_1 = -2$ in this inequality to obtain

$$2^k \|\psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2} + \|d_{x,t}(\psi\Gamma_k u)\|_{L_{x_1}^\infty L_{x',t}^2} \lesssim \|P(\psi\Gamma_k u)\|_{L_{x_1}^1 L_{x',t}^2}.$$

Hence (2.7) implies that

$$\|\psi\Gamma_k u\|_{L_x^q L_t^2([-1,1] \times \mathbb{R}^n)} \lesssim 2^{k(\delta(q) + \frac{\sigma}{q} - 1)} \left(2^k \|\Gamma_k u\|_{L_{x,t}^2} + \|P(\psi\Gamma_k u)\|_{L_{x_1}^1 L_{x',t}^2} \right).$$

We now write

$$P(\psi\Gamma_k u) = [P, \psi]\Gamma_k u + \psi[P, \Gamma_k]u + \psi\Gamma_k P u.$$

Since the inequality

$$\begin{aligned} \sum_k 2^{2k(\delta(q) + \frac{\sigma}{q} - 1)} (\|\psi \Gamma_k P u\|_{L_{x_1}^1 L_{x',t}^2}^2 + \|[P, \psi] \Gamma_k u\|_{L_{x,t}^2}^2 + 2^{2k} \|\Gamma_k u\|_{L_{x,t}^2}^2) \\ \lesssim \|P u\|_{L_t^2 H^{\delta(q) + \frac{\sigma}{q} - 1}}^2 + \|u\|_{L_t^2 H^{\delta(q) + \frac{\sigma}{q}}}^2 \end{aligned}$$

is straightforward, the key step in the proof will lie in showing that

$$(2.8) \quad \sum_k 2^{2k(\delta(q) + \frac{\sigma}{q} - 1)} \|\psi [P, \Gamma_k] u\|_{L_{x_1}^1 L_{x',t}^2}^2 \lesssim \|u\|_{L_t^2 H^{\delta(q)}}^2 + \|\partial_t u\|_{L_t^2 H^{\delta(q) - 1}}^2 + \|P u\|_{L_t^2 H^{\delta(q) - 1}}^2.$$

This is where the proof of our statement varies from the aforementioned Theorem 2.1 in [3].

We now employ an application of Lemma 2.4 in [3] similar to that used in Proposition 2.1, this time setting $R_\theta = \sum_{k=0}^{\infty} e^{ik\theta} 2^{k(\delta(q) + \frac{\sigma}{q} - 1)} \Gamma_k$. In the case where $0 \leq \delta(q) + \frac{\sigma}{q} \leq 1$, the corollary of the Coifman-Meyer theorem stated above together with an argument similar to that in (2.5) imply that

$$\|\psi [P, R_\theta] u\|_{L_{x,t}^2} \lesssim \|u\|_{L_t^2 H^{\delta(q) + \frac{\sigma}{q}}} + \|\partial_t u\|_{L_t^2 H^{\delta(q) + \frac{\sigma}{q} - 1}}$$

and hence (2.8) follows as before.

The case $1 \leq \delta(q) + \frac{\sigma}{q} \leq \frac{3}{2}$ is only considered when $\sigma = 0$ and $\kappa > \frac{n-1}{r} + 2$ as we will need the additional regularity of the coefficients in this instance. Once again, it suffices to establish mapping properties of $[P, R_\theta]$. In this case, we write

$$[P, R_\theta] = [\rho, R_\theta] \partial_t^2 - \sum_{ij} [a_{ij}, R_\theta] \partial_{x_i x_j}^2 - \sum_j [b_j, R_\theta] \partial_{x_j}.$$

We now use (2.3) to get that

$$[a_{ij}, R_\theta] : H^{\delta(q) - 2}(\mathbb{R}^{n+1}) \rightarrow L^2(\mathbb{R}^{n+1})$$

with operator norm independent of θ . This in turn implies that

$$\|\psi [P + \sum_{j=1}^n b_j \partial_{x_j}, R_\theta] u\|_{L_{x,t}^2}^2 \lesssim \sum_{l=0}^2 \|\partial_t^l u\|_{L_t^2 H^{\delta(q) - l}}^2.$$

To control the sum on the right, observe that since $\frac{1}{\rho(x)} \in Lip(\mathbb{R}^n)$, multiplication by this function is a bounded operation on $H^{\delta(q) - 2}(\mathbb{R}^n)$ and hence the equation gives us that

$$\|\partial_t^2 u\|_{L_t^2 H^{\delta(q) - 2}} \lesssim \|\rho \partial_t^2 u\|_{L_t^2 H^{\delta(q) - 2}} \lesssim \|P u\|_{L_t^2 H^{\delta(q) - 2}} + \|A(\cdot, D) u\|_{L_t^2 H^{\delta(q) - 2}}.$$

Since $A(x, D)$ is a divergence form operator with Lipschitz coefficients we have that $A(x, D) : H^{\delta(q)} \rightarrow H^{\delta(q) - 2}$ and hence

$$\sum_{l=0}^2 \|\partial_t^l u\|_{L_t^2 H^{\delta(q) - l}} \lesssim \|u\|_{L_t^2 H^{\delta(q)}} + \|\partial_t u\|_{L_t^2 H^{\delta(q) - 1}} + \|P u\|_{L_t^2 H^{\delta(q) - 1}}.$$

Turning our attention to the first order terms, observe that the higher degree of regularity in the coefficients implies that multiplication by the function b_j is a bounded operation on $H^{\delta(q) - 1}(\mathbb{R}^n)$. Indeed, when $1 < r \leq 2$, Sobolev embedding

tells us that $\partial_i a_{ij} \in H^{\frac{n}{2}+1-\frac{1}{r}}$ and when $2 < r < \infty$, $\partial_i a_{ij} \in C^{1-\frac{1}{r}}$. In either case, this implies that

$$\|[b_j, R_\theta] \partial_{x_j} u\|_{L^2} \lesssim \|\partial_{x_j} u\|_{L^2 H^{\delta(q)-1}}$$

and hence $[P, R_\theta]$ possesses the desired mapping properties. \square

3. THE HALF-WAVE OPERATOR

In this section, we show that the $\psi \Gamma_k u$ are also solutions to a first-order pseudodifferential equation that is hyperbolic in x_1 . In future sections, we will construct a parametrix for this first-order operator that yields the desired squarefunction estimates. Throughout this discussion, we assume that $a_{ij} - \delta_{ij} \in L^{r, \kappa}$ with $\kappa = \frac{n-1}{r} + 1 + \alpha$, $1 < r < \infty$, and either $\alpha > 1$ or $0 < \alpha < 1$.

To construct the first order pseudodifferential operator, we let $\Psi \in C^\infty(\mathbb{R}^n)$ be a cutoff function in (ξ', τ) such that $\text{supp}(\Psi) \subset \{(\xi', \tau) : -\tau \geq 5|\xi'|\}$ and is identically 1 on the smaller region $\{(\xi', \tau) : -\tau \geq 6|\xi'|, |(\xi', \tau)| \geq \frac{1}{2}\}$. We can also take the function so that

$$\Psi(\xi', \tau) = \Psi(|(\xi', \tau)|/|(\xi', \tau)|) \quad \text{for } |(\xi', \tau)| \geq \frac{3}{4}.$$

Let $\tilde{a}_{ij} = a_{11}^{-1} a_{ij}$, $\tilde{b}_j = \rho^{-1} \sum_i \partial_i a_{ij}$, and set

$$\tilde{a}_{ij}^k = \sum_{l \leq \varsigma k} \beta_l(D) \tilde{a}_{ij}$$

where $\varsigma = \max(\frac{1}{2}, \frac{2}{3+\alpha})$. Define Q as the pseudodifferential operator with rough symbol

$$\begin{aligned} Q(x, \xi, \tau) &= \xi_1^2 + \left(2 \sum_{j=2}^n \tilde{a}_{1j}(x) \xi_1 \xi_j - (\tilde{\rho}(x) \tau^2 - \sum_{2 \leq i, j \leq n} \tilde{a}_{ij}(x) \xi_i \xi_j) \right) \Psi(\xi', \tau) \\ &\quad - (\tau^2 + \sum_{i=2}^n \xi_i^2) (1 - \Psi(\xi', \tau)) \end{aligned}$$

and let Q_k be the analogous operator defined by replacing the \tilde{a}_{ij} in the symbol defining Q with their regularized counterparts \tilde{a}_{ij}^k . Observe that by construction, if $w(x, t)$ is a function such that $\text{supp}(\hat{w}) \subset \Psi^{-1}(1)$, then

$$Qw = \frac{1}{\rho} Pw + \sum_{ij} \frac{\partial_i a_{ij}}{\rho} \partial_j w.$$

In particular, by choosing $M \geq 200$ in (2.6), this holds for the $\psi \Gamma_k u$. On the other hand, since $Q(x, \xi, \tau)$ and $Q_k(x, \xi, \tau)$ are quadratics in ξ_1 with real roots, these operators are hyperbolic in x_1 rather than in t .

We now let $\varrho(x_1)$ be the Hardy-Littlewood maximal function associated to the $L^r(\mathbb{R})$ function

$$\sum_{|\beta| \leq \min(\alpha, 2)} \sum_{i, j} \|\partial_x^\beta a_{ij}(x_1, \cdot)\|_{L_{x'}^\infty} + \sum_{|\beta| \leq \min(\alpha, 2)} \|\partial_x^\beta \rho(x_1, \cdot)\|_{L_{x'}^\infty}.$$

Note that since $1 < r < \infty$, $\varrho(x_1) \in L^r(\mathbb{R})$ as well. By the arguments in section 3 and section 5 of [3], we have that

$$(3.1) \quad \|(\tilde{a}_{ij} - \tilde{a}_{ij}^k)(x_1, \cdot)\|_{L_{x'}^\infty} \leq C2^{-k \min(1, \varsigma(1+\alpha))} \varrho(x_1)$$

$$(3.2) \quad \|\partial_x^\beta \tilde{a}_{ij}^k(x_1, \cdot)\|_{L_{x'}^\infty} \leq C_\beta 2^{k\varsigma \max(0, |\beta| - 1 - \alpha)} \varrho(x_1)$$

$$(3.3) \quad \|\partial_x^\beta \tilde{a}_{ij}^k\|_{L^\infty} \leq C_\beta 2^{k\varsigma \max(0, |\beta| - 1)}$$

We now decompose

$$Q_k(x, \xi, \tau) = (\xi_1 + q_k^-(x, \xi', \tau))(\xi_1 - q_k^+(x, \xi', \tau))$$

so that $q_k^\pm(x, \xi', \tau)$ are positive functions that yield the roots of $Q_k(x, \xi, \tau)$ as a quadratic in ξ_1 . Next set

$$p_k(x, \xi', \tau) = \sum_{l \leq \varsigma k} \beta_l(D_x) q_k^-(x, \xi', \tau)$$

so that when $|\eta| \geq 2^{k\varsigma}$,

$$\int e^{-i\langle x, \eta \rangle} p_k(x, \xi', \tau) = 0.$$

In addition, if $\hat{w}_k(\xi', \tau)$ is a function supported in a set K , then the Fourier transform of $p_k(x, D_{x', t}) w_k$ will be supported in a set of the form $K + \{|\xi', \tau| \leq C2^{k\varsigma}\}$. Observe also that $p_k(x, \xi', \tau)$ approximates $q_k^-(x, \xi', \tau)$ and satisfies estimates analogous to (3.1), (3.2), and (3.3), uniformly for $|\xi', \tau| = 1$.

For the rest of this section we will take the convention that $\sigma = \max(0, \frac{1-\alpha}{3+\alpha})$. By the arguments in section 5 of [3], there exists a collection of $\mathcal{O}(2^{k\sigma})$ subintervals, $\{I_m^k\}_m$, such that $\cup_m I_m^k = (-2, 2)$, $|I_m^k| \leq 2^{-k\sigma}$ for each m , and

$$(3.4) \quad \int_{I_m^k} \varrho(x_1) dx_1 \ll 2^{-k\sigma}$$

for every m as well.

We now want to rescale the problem. Set $w_k(x, t) = \psi \Gamma_k u(2^{-k\sigma}(x, t))$, $\lambda = 2^{k(1-\sigma)}$, and $\tilde{p}_\lambda(x, \xi', \tau) = p_k(2^{-k\sigma}x, \xi', \tau)$. We also set $\tilde{\varrho}(x_1) = \varrho(2^{-k\sigma}x_1)$ so that $\|\tilde{\varrho}\|_{L^r(I)} \ll 1$ by (3.4). We now have that $\tilde{p}_\lambda(x, \xi', \tau)$ satisfies

$$(3.5) \quad \sup_{x \in \mathbb{R}^n, (\xi', \tau) \in \Xi} \sum_{|\alpha| + |\beta| \leq 2, |\alpha| \leq 1} |\partial_x^\alpha \partial_\xi^\beta (\tilde{p}_\lambda(x, \xi', \tau) - \sqrt{\tau^2 - |\xi'|^2})| \ll 1$$

$$(3.6) \quad \sup_{x' \in \mathbb{R}^{n-1}, (\xi', \tau) \in \Xi} \sum_{|\alpha| + |\beta| \leq 2} |\partial_x^\alpha \partial_\xi^\beta (\tilde{p}_\lambda(x, \xi', \tau) - \sqrt{\tau^2 - |\xi'|^2})| \leq \tilde{\varrho}(x_1)$$

where $\Xi = \{(\xi', \tau) \in \mathbb{S}^{n-1} : -\tau \geq 6|\xi'|\}$. In addition, we have the following estimates on the partial derivatives of \tilde{p}_λ

$$(3.7) \quad \sup_{(x, \xi', \tau) \in \mathbb{R}^n \times \mathbb{S}^{n-1}} |\partial_x^\alpha \partial_\xi^\beta \tilde{p}_\lambda(x, \xi', \tau)| \leq C_{\alpha, \beta} \lambda^{\max(0, |\alpha| - 1)}$$

$$(3.8) \quad \sup_{(x', \xi', \tau) \in \mathbb{R}^{n-1} \times \mathbb{S}^{n-1}} |\partial_x^\alpha \partial_\xi^\beta \tilde{p}_\lambda(x, \xi', \tau)| \leq C_{\alpha, \beta} \tilde{\varrho}(x_1) \lambda^{\max(0, |\alpha| - 2)}$$

with $C_{\alpha,\beta}$ independent of λ . We also have

$$(3.9) \quad \int e^{i\langle x,\eta\rangle + it\zeta} \tilde{p}_\lambda(x, \xi', \tau) dt dx = 0 \text{ for } |(\eta, \zeta)| \geq \lambda^{\frac{1}{2}}$$

In the next two sections, we prove the following theorem:

Theorem 3.1. *Let $I \subset [-2, 2]$ be any interval containing 0 and $p(x, \xi', \tau)$ a symbol satisfying (3.5), (3.6), (3.7), (3.8), and (3.9). Suppose v solves the first order equation*

$$\begin{aligned} (\partial_{x_1} + ip(x, D_{x',t}))v(x, t) &= F(x, t) \in L_{x_1}^1(I; L_{x',t}^2(\mathbb{R}^n)) \\ v(0, x', t) &= v_0(x', t) \in L^2(\mathbb{R}^n) \end{aligned}$$

and that

$$(3.10) \quad \widehat{v}(x_1, \xi', \tau) \subset \{(\xi', \tau) : 2^{-5}\lambda \leq |(\xi', \tau)| \leq 2^6\lambda\} \cap \{(\xi', \tau) : -\tau \geq \widetilde{M}|\xi'|\}$$

for some λ, \widetilde{M} sufficiently large. Then

$$\|v\|_{L_{x_1, x'}^q(I \times \mathbb{R}^{n-1}; L_t^2(\mathbb{R}))} \lesssim \lambda^{\delta(q)} (\|v_0\|_{L^2} + \|F\|_{L_{x_1}^1(I; L_{x',t}^2(\mathbb{R}^n))})$$

This theorem will imply the following estimate on w_k over the set $2^{k\sigma}I_m^k = \{2^{k\sigma}x_1 : x_1 \in I_m^k\}$

$$\begin{aligned} \|w_k\|_{L_x^q L_t^2(2^{k\sigma}I_m^k \times \mathbb{R}^n)} &\lesssim 2^{k(1-\sigma)\delta(q)} \left(\|w_k\|_{L_{x_1}^\infty L_{x',t}^2(2^{k\sigma}I_m^k \times \mathbb{R}^n)} \right. \\ &\quad \left. + \|(-i\partial_1 + \tilde{p}(x, D_{x',t}))w_k\|_{L_{x_1}^1 L_{x',t}^2(2^{k\sigma}I_m^k \times \mathbb{R}^n)} \right). \end{aligned}$$

We can then rescale the estimate to obtain

$$(3.11) \quad \begin{aligned} \|\psi\Gamma_k u\|_{L_x^q L_t^2(I_m^k \times \mathbb{R}^n)} &\lesssim 2^{k\delta(q)} \left(\|\psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2(I_m^k \times \mathbb{R}^n)} \right. \\ &\quad \left. + \|(-i\partial_1 + p_k(\cdot, D_{x',t}))\psi\Gamma_k u\|_{L_{x_1}^1 L_{x',t}^2(I_m^k \times \mathbb{R}^n)} \right) \end{aligned}$$

In order to see that this implies the estimates in the hypothesis of Proposition 2.2, we will show following estimate

$$(3.12) \quad \begin{aligned} \|(-i\partial_1 + p_k(\cdot, D_{x',t}))\psi\Gamma_k u\|_{L_{x_1}^1 L_{x',t}^2(I_m^k \times \mathbb{R}^n)} &\lesssim \|\psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2} \\ &\quad + 2^{-k} \|\partial_1 \psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2} + 2^{-k} \|P(\psi\Gamma_k u)\|_{L_{x_1}^1 L_{x',t}^2} + \|\Gamma_k u\|_{L_{x,t}^2}. \end{aligned}$$

Combining this with (3.11), we can sum over m to obtain

$$\begin{aligned} \|\psi\Gamma_k u\|_{L_x^q L_t^2} &\lesssim \left(\sum_m \|\psi\Gamma_k u\|_{L_x^q L_t^2(I_m^k \times \mathbb{R}^n)}^q \right)^{\frac{1}{q}} \\ &\lesssim 2^{k(\delta(q) + \frac{\sigma}{q})} \left(\|\psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2} + 2^{-k} \|\partial_1 \psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2} \right. \\ &\quad \left. + 2^{-k} \|P(\psi\Gamma_k u)\|_{L_{x_1}^1 L_{x',t}^2} + \|\Gamma_k u\|_{L_{x,t}^2} \right) \end{aligned}$$

Proof of (3.12). Let $\widetilde{\Gamma}_1(\tau, \xi)$ be bump function that is identically one on the support of Γ_1 and is supported in a slightly larger cone containing $\text{supp}(\Gamma_1)$. With this, define $\widetilde{\Gamma}_k(\cdot) = \widetilde{\Gamma}_1(2^{-k}\cdot)$, and choose $\widetilde{\Gamma}_1^{-1}(1)$ large enough so that $\widetilde{\Gamma}_k \cdot \Gamma_k = \Gamma_k$ and

$$\widetilde{\Gamma}_k(D) \circ (-i\partial_1 - p_k(\cdot, D_{x',t}))\Gamma_k(D) = (-i\partial_1 - p_k(\cdot, D_{x',t}))\Gamma_k(D).$$

We can now write

$$\begin{aligned} & (-i\partial_1 + p_k(\cdot, D_{x',t}))\psi\Gamma_k u \\ &= \tilde{\Gamma}_k(-i\partial_1 + p_k(\cdot, D_{x',t}))\psi\Gamma_k u - [\tilde{\Gamma}_k, (-i\partial_1 + p_k(\cdot, D_{x',t}))\psi]\Gamma_k u \end{aligned}$$

By the symbolic calculus and (3.3), the commutator lies in $S_{1,\varsigma}^0(\mathbb{R}^{2n+2})$ and as such the L^2 norm of the second term can be bounded by the last term on the right hand side of (3.12). Rewrite the operator in the first term as

$$(3.13) \quad \begin{aligned} \tilde{\Gamma}_k(-i\partial_1 + p_k(\cdot, D_{x',t})) &= \tilde{\Gamma}_k(-i\partial_1 + q_k^-(\cdot, D_{x',t})) + \\ & \quad [\tilde{\Gamma}_k, (p_k(\cdot, D_{x',t}) - q_k^-(\cdot, D_{x',t}))] + (p_k(\cdot, D_{x',t}) - q_k^-(\cdot, D_{x',t}))\tilde{\Gamma}_k \end{aligned}$$

As above, the commutator term lies in $S_{1,\varsigma}^0(\mathbb{R}^{2n+2})$. In addition for $|(\xi', \tau)| \approx 2^k$ the argument used to give (3.1) also yields

$$|\partial_x^\gamma \partial_{\xi', \tau}^\beta (q_k^-(x, \xi', \tau) - p_k(x, \xi', \tau))| \leq C_{\alpha, \beta} \varrho(x_1) 2^{k\sigma + k\varsigma|\gamma| - k|\beta|}$$

as $1 - \varsigma(1 + \alpha) = \sigma$. By (3.4), $\|\varrho\|_{L^1(I_m^k)} \leq 2^{-k\sigma}$ which implies

$$(3.14) \quad \|(q_k^+(\cdot, D_{x',t}) - p_k^+(\cdot, D_{x',t}))\tilde{\Gamma}_k\psi\Gamma_k u\|_{L_{x_1}^1 L_{x',t}^2(I_m^k \times \mathbb{R}^n)} \lesssim \|\psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2(I_m^k \times \mathbb{R}^n)}$$

To handle the first term in (3.13), let

$$B_k(x, \xi, \tau) = \frac{\tilde{\Gamma}_k(\xi, \tau)}{\xi_1 - q_k^+(x, \xi', \tau)} \quad \text{so that} \quad B_k, d_x B_k \in S_{1,\varsigma}^{-1}$$

which is well defined as $\tilde{\Gamma}_k$ is supported away from the zero set of the denominator. Integration by parts yields the following estimates on the integral kernel of B_k

$$(3.15) \quad 2^k |\partial_{x,t}^\alpha K_k(x, t; y, s)| \lesssim 2^{k(n+1)} (1 + 2^k |(x, t) - (y, s)|)^{-(n+2)}, \quad |\alpha| \leq 1.$$

Observe that by (3.3) and the symbolic calculus we have

$$\tilde{\Gamma}_k(D) = B_k(x, D) \circ (-i\partial_1 - q_k^+(x, D)) \quad \text{mod } S_{1,\varsigma}^{-1}$$

$$(-i\partial_1 - q_k^+(x, D_{x',t})) \circ (-i\partial_1 + q_k^-(x, D_{x',t})) = Q_k(x, D) \quad \text{mod } S_{1,\varsigma}^1.$$

This allows us to write

$$\tilde{\Gamma}_k(-i\partial_1 + q_k^-(\cdot, D_{x',t}))\psi\Gamma_k u = (B_k(x, D) \circ Q_k(x, D) + R_k(x, D))\psi\Gamma_k u$$

where $R_k(x, \xi, \tau) \in S_{1,\varsigma}^0$. We have now shown that

$$\begin{aligned} & \|(-i\partial_1 + p_k(\cdot, D_{x',t}))\psi\Gamma_k u\|_{L_{x_1}^1 L_{x',t}^2(I_m^k \times \mathbb{R}^n)} \\ & \lesssim \|\psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2(I_m^k \times \mathbb{R}^n)} + \|\Gamma_k u\|_{L_{x,t}^2} + \|B_k \circ Q_k(\psi\Gamma_k u)\|_{L_{x_1}^1 L_{x',t}^2(I_m^k \times \mathbb{R}^n)} \end{aligned}$$

To estimate the last term on the right hand side we write

$$B_k \circ Q_k = [B_k, Q_k - Q] + (Q_k - Q)B_k + B_k Q$$

Using (3.15), we can directly estimate the kernel of $[B_k, a]$ for Lipschitz functions a and use Schur's Lemma to get

$$\|[B_k, Q_k - Q](\psi\Gamma_k u)\|_{L_{x_1}^1 L_{x',t}^2(I_m^k \times \mathbb{R}^n)} \lesssim 2^{-2k} \|d_{x,t}^2(\psi\Gamma_k u)\|_{L_{x,t}^2} \lesssim \|\Gamma_k u\|_{L_{x,t}^2}$$

As in (3.14) we also have

$$\|(Q_k - Q)B_k(\psi\Gamma_k u)\|_{L_{x_1}^1 L_{x',t}^2(I_m^k \times \mathbb{R}^n)} \lesssim \|d_{x',t} B_k(\psi\Gamma_k u)\|_{L_{x_1}^\infty L_{x',t}^2} \lesssim \|\psi\Gamma_k u\|_{L_{x_1}^\infty L_{x',t}^2}$$

after an application of (3.15) and Schur's Lemma in the second inequality. We now use (3.15) again to complete the proof by arguing

$$\|B_k Q(\psi \Gamma_k u)\|_{L_{x_1}^1 L_{x',t}^2(I_m^k \times \mathbb{R}^n)} \lesssim 2^{-k} \|P(\psi \Gamma_k u)\|_{L_{x_1}^1 L_{x',t}^2} + 2^{-k} \|\partial_1 \psi \Gamma_k\|_{L_{x_1}^\infty L_{x',t}^2}$$

□

4. THE WAVE TRANSFORM

We have now reduced the problem of showing squarefunction inequalities for general solutions of the wave equation to estimating the $L_x^q L_t^2$ norm of solutions to a x_1 -hyperbolic pseudodifferential equation. To show Theorem 3.1, we proceed by constructing a parametrix for such operators that is suitable for proving these estimates. As mentioned previously, we employ the continuous wave packet transform of used by Smith in [8] for this purpose.

For ease of exposition of the arguments we will reverse the roles of t and x_1 and assume that our first order pseudodifferential equation is hyperbolic in t . Therefore, in this discussion and in section 5, we assume that $p(x, t, \xi)$ is a first order symbol close to $\sqrt{\xi_1^2 - |\xi'|^2}$, meaning that (3.6) and (3.8) take the form

$$(4.1) \quad \sup_{x \in \mathbb{R}^n, \xi \in \Xi} \sum_{|\alpha|+|\beta| \leq 2} |\partial_{x,t}^\alpha \partial_\xi^\beta (p(x, t, \xi) - \sqrt{\xi_1^2 - |\xi'|^2})| \leq \varrho(t)$$

$$(4.2) \quad \sup_{(x', \xi) \in \mathbb{R}^{n-1} \times \mathbb{S}^{n-1}} |\partial_{x,t}^\alpha \partial_\xi^\beta p(x, t, \xi)| \leq C_{\alpha,\beta} \varrho(t) \lambda^{\max(0, |\alpha| - 2)}$$

where $\Xi = \{\xi \in \mathbb{S}^{n-1} : -\xi_1 \geq 6|\xi'|\}$ and again $\|\varrho\|_{L^r(I)} \ll 1$. We thus seek to show that solutions $v(x, t)$ to

$$\begin{aligned} (\partial_t + ip(x, t, D_x))v(x, t) &= F(x, t) \in L_t^1(I; L_x^2(\mathbb{R}^n)) \\ v(x, 0) &= v_0(x) \in L^2(\mathbb{R}) \end{aligned}$$

such that

$$(4.3) \quad \widehat{v}(\xi, t) \subset \{\xi : 2^{-5}\lambda \leq |\xi| \leq 2^6\lambda\} \cap \{\xi : -\xi_1 \geq \widetilde{M}|\xi'|\}$$

satisfy over a time interval $t \in [-2, 2]$

$$\|v\|_{L_{x',t}^q(\mathbb{R}^{n-1} \times [-2,2]; L_{x_1}^2(\mathbb{R}))} \lesssim \lambda^{\delta(q)} (\|v_0\|_{L^2} + \|F\|_{L_t^1 L_x^2}).$$

We now introduce the continuous wave packet transform. Let g be a radial Schwartz function over \mathbb{R}^n such that $\text{supp}(\widehat{g}) \subset B_1(0)$ and $\|g\|_{L^2} = (2\pi)^{-\frac{n}{2}}$. For $\lambda \geq 1$, we define the operator $T_\lambda : \mathcal{S}'(\mathbb{R}^n) \rightarrow C^\infty(\mathbb{R}^{2n})$ by

$$(T_\lambda f)(x, \xi) = \lambda^{\frac{n}{4}} \int e^{-i\langle \xi, z-x \rangle} g(\lambda^{\frac{1}{2}}(z-x)) f(z) dz.$$

T_λ enjoys the property that its adjoint as a map from $L_{x,\xi}^2(\mathbb{R}^{2n}) \rightarrow L_z^2(\mathbb{R}^n)$ also serves as a left inverse for T_λ , that is, $T_\lambda^* T_\lambda = I$. This implies that T_λ is an isometry:

$$\|T_\lambda f\|_{L_{x,\xi}^2(\mathbb{R}^{2n})} = \|f\|_{L_z^2(\mathbb{R}^n)}.$$

In [8], H. Smith proved that for $v(x, t)$ with partial Fourier transform satisfying (4.3) there exists a function $G(x, t, \xi)$ with the property that

$$\|G(\cdot, t, \cdot)\|_{L_{x,\xi}^2(\mathbb{R}^{2n})} \lesssim \varrho(t) \|v(\cdot, t)\|_{L_z^2(\mathbb{R}^n)}$$

satisfying the equation

$$(T_\lambda v(\cdot, t))(x, \xi) = (T_\lambda f)(\chi_{0,t}(x, \xi)) + \int_0^t G(r, \chi_{r,t}(x, \xi)) + (T_\lambda F)(r, \chi_{r,t}(x, \xi)) dr.$$

Here $\chi_{r,t}(x, \xi) = (x_{r,t}, \xi_{r,t})$ denotes the time r solution to the initial value problem

$$\frac{dx}{ds} = d_\xi p(x, t, \xi), \quad \frac{d\xi}{ds} = -d_x p(x, t, \xi), \quad (x(t), \xi(t)) = (x, \xi).$$

See Lemmas 3.1, 3.2, and 3.3 of [8] for a proof of the result. Since T_λ^* is a left inverse for T_λ , this in turn implies that

$$v(\cdot, t) = T_\lambda^*(T_\lambda f \circ \chi_{0,t}) + \int_0^t T_\lambda^*(G_r \circ \chi_{r,t})(x, \xi) dr + \int_0^t T_\lambda^*((T_\lambda F_r) \circ \chi_{r,t})(x, \xi) dr$$

with $G_r(x, \xi) = G(x, r, \xi)$ and $F_r(z) = F(z, r)$. As in [8], we now define the operator W_r mapping functions on $\mathbb{R}_{x,\xi}^{2n}$ to $\mathbb{R}_{y,s}^{n+1}$ by

$$(W_r f)(y, s) = T_\lambda^*(f \circ \chi_{r,s})(y).$$

Let $\tilde{\Psi}_\lambda(\xi)$ denote a smooth cutoff identically 1 on the set in (4.3) and supported in a cone of slightly larger aperture in a set with $|\xi| \approx \lambda$. Since T_λ is an isometry, in order to control the $L_{y',s}^q L_{y_1}^2$ norm of u , it will suffice to show that

$$\tilde{\Psi}_\lambda(D)W_r : L_{x,\xi}^2(\mathbb{R}^{2n}) \rightarrow L_{y',s}^q(\mathbb{R}^{n-1} \times [-2, 2]; L_{y_1}^2(\mathbb{R}))$$

with operator norm bounded by $\lambda^{\delta(q)}$ uniformly for $r \in [-2, 2]$. By duality, this is equivalent to showing that

$$(4.4) \quad \tilde{\Psi}_\lambda(D)W_r W_r^* \tilde{\Psi}_\lambda(D) : L_{x',t}^{q'}(\mathbb{R}^{n-1} \times [-2, 2]; L_{x_1}^2(\mathbb{R})) \rightarrow L_{y',s}^q(\mathbb{R}^{n-1} \times [-2, 2]; L_{y_1}^2(\mathbb{R})).$$

also with uniform estimates on the operator norm.

Next observe that we can characterize the action of the operator $W_r W_r^* \tilde{\Psi}_\lambda(D)$ on a function $\tilde{F}(z, t)$ as integration of the function $(\tilde{\Psi}_\lambda \tilde{F})(z, t)$ against a kernel $K(y, s; z, t)$ where

$$K(y, s; z, t) = \lambda^{\frac{n}{2}} \int e^{i\langle \xi, y-x \rangle - i\langle \xi_{t,s}, z-x_{t,s} \rangle} g(\lambda^{\frac{1}{2}}(y-x)) g(\lambda^{\frac{1}{2}}(z-x_{t,s})) \gamma(\lambda^{-1}\xi) dx d\xi$$

and γ is a smooth cutoff function supported in a small cone about the negative ξ_1 -axis. This is accomplished by a straightforward computation which uses the fact that $\chi_{r,s}$ is a measure-preserving diffeomorphism and that $\chi_{t,r} \circ \chi_{r,s} = \chi_{t,s}$. Observe also that by choosing M sufficiently large in (2.6) we can assume that the support of γ is contained in a cone with aperture that can be made arbitrarily small.

We now seek to show the following pair of estimates

$$(4.5) \quad \left\| \int K(y, s; z, t) f(z) dz \right\|_{L_y^2(\mathbb{R}^n)} \lesssim \|f\|_{L_z^2(\mathbb{R}^n)}$$

$$(4.6) \quad \left\| \int K(y, s; z, t) f(z) dz \right\|_{L_{y'}^\infty L_{y_1}^2(\mathbb{R}^n)} \lesssim \lambda^{n-1} (1 + \lambda|t-s|)^{-\frac{n-1}{2}} \|f\|_{L_z^1 L_{z_1}^2(\mathbb{R}^n)}.$$

By interpolation and the fact that $\delta(q) \geq \frac{2(n+1)}{n-1}$, these estimates imply the inequality

$$\left\| \int K(y, s; z, t) f(z) dz \right\|_{L_{y'}^q L_{z_1}^2(\mathbb{R}^n)} \lesssim \lambda^{2\delta(q)} |t-s|^{-1+\frac{1}{q'}-\frac{1}{q}} \|f\|_{L_{z'}^{q'} L_{z_1}^2(\mathbb{R}^n)}.$$

The estimates (4.4) will then follow by the Hardy-Littlewood-Sobolev inequality of fractional integration.

To establish (4.5), we observe that

$$\int K(y, s; z, t) (\Gamma_\lambda f)(z) dz = T_\lambda^* ((T_\lambda \Gamma_\lambda f) \circ \chi_{t,s})(y).$$

The inequality then follows since $\chi_{t,s}$ is a symplectomorphism and by the L^2 mapping properties of T_λ .

Most of the work is thus involved in establishing (4.6). The main idea here is to obtain pointwise estimates on the $K(y, s; z, t)$ that will imply the desired mapping properties. These inequalities will reflect the decay of K away from the light cone as well as the dispersive properties of K . Additionally, since γ is localized to a cone about the negative ξ_1 axis, $K(y, s; z, t)$ should decay rapidly when $|y' - z'| \gg |y_1 - z_1|$. Therefore we should only need to examine properties of the part of the light cone whose projection down to the spatial variables lies in the complement of this set.

To accomplish this, we begin by establishing the existence of functions $\Phi_{z,t}^\pm(y)$, defined when $\mp(y_1 - z_1) \gg |y' - z'|$, such that $(y, \Phi_{z,t}^+(y))$ (or $(y, \Phi_{z,t}^-(y))$) is the graph of the forward (backward) light cone centered at (z, t) . We then seek to establish the estimate:

$$(4.7) \quad |K(y, s; z, t)| \leq C_N \lambda^n (1 + \lambda|s-t|)^{-\frac{n-1}{2}} (1 + \lambda V(y, s; z, t))^{-N}$$

where V is a nonnegative measurable function. V will have the property that

$$V(y, s; z, t) = |s - \Phi_{z,t}^\pm(y)|$$

for $|s|, |t| \leq 2$, $|y-z| \leq \frac{5}{2}$, and $y-z \in \tilde{\Omega}_\pm$ with $\tilde{\Omega}_+$ (resp. $\tilde{\Omega}_-$) a small cone about the negative (resp. positive) x_1 -axis to be determined later. Here $\Phi_{z,t}^+$ is used in the estimate if $s > t$ and $\Phi_{z,t}^-$ is used otherwise. For $(y, s; z, t)$ not in this domain, we have $V(y, s; z, t) = |y-z|$. As we will see, $\partial_{y_1} \Phi_{z,t}^\pm(y)$ is uniformly bounded from below on the region which it is defined, allowing us to get suitable control over

$$\sup_z \int |K(y, s; z, t)| dy_1$$

by changing variables. Hence the estimate (4.6) will follow by Schur's Lemma and the symmetry of (y, s) and (z, t) in the definition of K .

The estimates (4.7) will follow by a scaling argument of Smith [8]. Assume without loss of generality that $s = 0$ and $t \in [0, 2]$. Let ε be a scaling factor such that $\lambda^{-1} \leq \varepsilon \leq 1$ and consider the scaled kernel

$$K_\varepsilon(y; z, t) = \varepsilon^n K(\varepsilon y, 0; \varepsilon z, \varepsilon t).$$

In addition, set $R = \varepsilon\lambda^{\frac{1}{2}}$, $\mu = \varepsilon\lambda$ so that a change of variables puts K_ε in the form

$$(4.8) \quad K_\varepsilon(y; z, t) = R^n \int e^{i\langle \xi, y-x \rangle - i\langle \xi_t, z-x_t \rangle} g(R(y-x))g(R(z-x_t))\gamma(\mu^{-1}\xi) dx d\xi$$

where (x_t, ξ_t) is now an integral curve of the flow determined by the Hamiltonian function

$$H(x, t, \xi) = p(\varepsilon x, \varepsilon t, \xi).$$

with initial conditions $(x_0, \xi_0) = (x, \xi)$. We will not need to refer to the first component of x_t specifically, so for the remainder of this paper, $x_1(x, \xi)$ will refer to a vector, namely the projection of the time 1 solution to Hamilton's equations down to the x variables with initial conditions (x, ξ) . The inequalities (4.7) are now a consequence of the following lemma, which is akin to Lemma 3.4 in [8].

Lemma 4.1. *Let Λ_z denote the $s = 0$ time slice of the backward light cone centered at $(1, z)$. Then for some C_0 sufficiently large and $\mu \geq C_0$, we have the following pointwise estimates on $K_\varepsilon(y; z, 1)$*

$$(4.9) \quad |K_\varepsilon(y; z, 1)| \leq \begin{cases} C_N \mu^{\frac{n+1}{2}} (1 + \mu d(y, \Lambda_z))^{-N} & y - z \in \tilde{\Omega}_- \\ C_{N, N'} \mu^{-N'} (1 + \mu |y - z|)^{-N} & y - z \notin \tilde{\Omega}_- \end{cases}$$

with constants $C_N, C_{N, N'}$ independent of $\mu \geq C_0$. Additionally, when $\mu = 1$ we then have that for all $t \in [0, C_0]$

$$(4.10) \quad |K_\varepsilon(y; z, t)| \lesssim (1 + |y - z|)^{-N}.$$

To see that this Lemma implies (4.7), observe that when $t \geq C_0\lambda^{-1}$, we can set $\varepsilon = t$ and rescale (4.9) to get that

$$(4.11) \quad |K(y, 0; z, t)| \lesssim \begin{cases} C_N \lambda^n (1 + \lambda t)^{-\frac{n-1}{2}} (1 + \lambda d(y, \Lambda_{z,t}))^{-N} & y - z \in \tilde{\Omega}_- \\ C_N (1 + \lambda t)^{-\frac{n-1}{2}} (1 + \lambda |y - z|)^{-N} & y - z \notin \tilde{\Omega}_- \end{cases}$$

where $\Lambda_{z,t}$ denotes the time 0 slice of the light cone centered at (z, t) . It is clear this establishes the desired estimate in (4.7) when $t \geq C_0\lambda^{-1}$ and $y - z \notin \tilde{\Omega}_-$. When $y - z \in \tilde{\Omega}_-$ and $|y - z| \geq \frac{5}{2}$ then y is uniformly at distance $\frac{1}{8}$ away from the light cone for $t \in [0, 2]$ and hence we can conclude that $d(y, \Lambda_{z,t}) \gtrsim 1 + |y - z|$. In the case where $y - z \in \tilde{\Omega}_-$ and $|y - z| \leq \frac{5}{2}$, y is in the domain of Φ^- and hence we will be able to use the approximation (5.10) below to obtain $|\Phi_{z,t}^-(y)| \approx d(y, \Lambda_{z,t})$ so that (4.7) holds.

In the case where $t \leq C_0\lambda^{-1}$, we set $\varepsilon = \lambda^{-1}$ so that rescaling (4.10) yields the desired inequalities in (4.7).

5. REGULARITY OF THE HAMILTONIAN FLOW

In this section we present the arguments needed to prove Lemma 4.1. The main idea here is to characterize integration against K_ε as being equivalent to the action of the adjoint of a Fourier integral operator. This means that we wish to find a phase function $\varphi_t(z, \eta)$ and symbol $a_t(z, \eta)$ (supported in a cone about η_1) such that

$$K_\varepsilon(y; z, t) = \int e^{i\langle y, \eta \rangle - i\varphi_t(z, \eta)} a_t(z, \eta) d\eta.$$

Once this is accomplished, we will be able to use a second dyadic decomposition of η -space to get the desired estimates on K_ε .

We begin with a simple lemma concerning the geometry of the integral curves (x_t, ξ_t) of the Hamiltonian flow induced by $H(x, t, \xi)$, that is, the solutions to the differential equation

$$(5.1) \quad \begin{aligned} \frac{dx}{dt} &= H_\xi(x, t, \xi) \\ \frac{d\xi}{dt} &= -H_x(x, t, \xi). \end{aligned}$$

Its proof will be omitted as it is an straightforward consequence of Gronwall's lemma and (4.1).

Lemma 5.1. *Let $(x_t(x, \xi), \xi_t(x, \xi))$ denote the solution to (5.1) with initial conditions $x(0) = x$, $\xi(0) = \xi$. Then for $-C_0 \leq t \leq C_0$ and $\xi \in \{\xi \in \mathbb{S}^{n-1} : -\xi_1 \geq 7|\xi'|\}$ we have*

$$(5.2) \quad \left| x_t - \left(x + t \frac{(\xi_1, -\xi')}{\sqrt{\xi_1^2 - |\xi'|^2}} \right) \right| + |\xi_t - \xi| \ll t$$

and

$$(5.3) \quad |d_x x_t - I| + |d_x \xi_t| + |d_\xi x_t - \int_0^t d_\xi^2 H(s, x_s, \xi_s) ds| + |d_\xi \xi_t - I| \ll 1.$$

By (5.3) we have that for any $\xi \in \mathbb{S}^{n-1}$ such that $-\xi_1 \geq 7|\xi'|$ and $t \in [-C_0, C_0]$, the map $x \mapsto x_t(x, \xi)$ is a local diffeomorphism on \mathbb{R}^n with uniform bounds on its differential. It is also a proper map since $|x_t(x, \xi) - x| \leq 5$. This implies that it is an open and closed map as well as that it is a covering map. By connectivity, its image is all of \mathbb{R}^n . Additionally, since \mathbb{R}^n is simply connected, it is also an injective map. Thus the map is a diffeomorphism and we let $z \mapsto \bar{x}_t(z, \xi)$ denote the inverse map.

We now define the generating function $\varphi_t(z, \xi)$ by

$$\varphi_t(z, \xi) = \langle \xi, \bar{x}_t(z, \xi) \rangle$$

which possesses the following properties for ξ such that $|\xi| = 1$ and $-\xi_1 \geq 7|\xi'|$

$$d_z \varphi_t(z, \xi) = \xi_t(\bar{x}_t(z, \xi), \xi) \quad d_\xi \varphi_t(z, \xi) = \bar{x}_t(z, \xi).$$

Also, by Lemma 4.4 of [8], we have that for $|\xi| = 1$ and $|t| \leq C_0$,

$$(5.4) \quad |\partial_z^\beta \partial_\xi^\alpha d_z^2 \varphi_t(z, \xi)| \leq C_{\alpha, \beta} R^{|\beta|} (1 + R)^{|\alpha|}$$

with $C_{\alpha, \beta}$ independent of choice of ξ, t . Such estimates follow by first differentiating Hamilton's equations to get estimates on derivatives of $(x_t(x, \xi), \xi_t(x, \xi))$ and subsequently applying the inverse function theorem.

In order to characterize K_ε as the integral kernel of the adjoint of a Fourier integral operator, we take the Fourier transform of the first factor of g in (4.8) and

write

$$(5.5) \quad \begin{aligned} K_\varepsilon(y; z, t) &= \int e^{i\langle \eta, y-x \rangle - i\langle \xi_t, z-x_t \rangle} h(R^{-1}(\eta - \xi)) g(R(z - x_t)) \gamma(\mu^{-1}\xi) dx d\xi d\eta \\ &= \int e^{i\langle y, \eta \rangle - i\varphi_t(z, \eta)} a_t(z, \eta) d\eta, \end{aligned}$$

where $h = \widehat{g}$ and

$$(5.6) \quad a_t(z, \eta) = \int e^{i\varphi_t(z, \eta) - i\langle x, \eta \rangle - i\langle \xi_t(x, \xi), z-x_t(x, \xi) \rangle} \cdot h(R^{-1}(\eta - \xi)) g(R(z - x_t(x, \xi))) \gamma(\mu^{-1}\xi) dx d\xi.$$

We note that since γ is compactly supported in a cone about the ξ_1 -axis, we have that for μ sufficiently large, $a_t(z, \cdot)$ must be supported in a cone of slightly larger aperture about the η_1 -axis. Observe also that by choosing M large in (2.6) and μ sufficiently large, we can assume that the support of $a_t(z, \cdot)$ lies in a cone of arbitrarily small (but fixed) aperture.

Citing Theorem 4.5 in [8] we have that for $t \in [-C_0, C_0]$ and $\mu \geq 1$, $a_t(z, \eta)$ satisfies the estimates

$$(5.7) \quad |\partial_\eta^\alpha \langle \eta, d_\eta \rangle^k a_t(z, \eta)| \leq C_{\alpha, k} (1 + R^2)^{\frac{|\alpha|}{2}} \mu^{-|\alpha|}$$

We pause to remark that the proofs of (5.4) and (5.7) in [8] depend on estimates on the derivatives of $x_t(x, \xi)$, $\xi_t(x, \xi)$, and $\bar{x}_t(z, \xi)$, but not on the specific geometry of the flow induced by $H(x, t, \xi)$.

We now introduce some technical notation. For $i = 0, 1, 2, 3, 4, 5$ we let Ω_i^+ denote small open cones about the negative ξ_1 -axis such that for any $z \in \mathbb{R}^n$,

$$\text{supp}(a_t(z, \cdot)) \subset \Omega_0^+ \subset \Omega_1^+ \subset \Omega_2^+ \subset \Omega_3^+ \subset \Omega_4^+ \subset \Omega_5^+$$

and the aperture of Ω_i^+ is strictly greater than that of Ω_{i-1}^+ . As with other conic sets of this type we assume that their aperture can be made arbitrarily small by choosing M in (2.6) and μ large. We also let

$$\Omega_i^- = \{\xi : (-\xi_1, \xi') \in \Omega_i^+\}.$$

Next we construct and analyze the $\Phi_{z,t}^\pm$ function introduced in the previous section. Let

$$\Lambda_i^\pm = \cup_{\pm t \in [0, 2]} \cup_{\xi \in \Omega_i^\pm} (x_t(0, \xi), t) \quad \text{and} \quad \Lambda_i^{\pm, t} = \cup_{\xi \in \Omega_i^\pm} (x_t(0, \xi), t)$$

be sections of the light cone or time slices of the light cone defined by H centered at the origin. We define Λ^\pm ($\Lambda^{\pm, t}$) similarly as the full (time slice of the) light cone at 0 by removing the condition $\xi \in \Omega_i^+$ above and replacing it with $\xi \in \mathbb{R}^n \setminus \{0\}$. We now have the following theorem, which is akin to Lemma 4.2 in [8].

Theorem 5.2. Λ_5^\pm is the graph $s = \Phi^\pm(y)$ of a function $\Phi^\pm(y)$ defined on a conic neighborhood containing $\Omega_4^\pm \cap \{z : |z| \leq 2\}$. The function $\Phi^\pm(y)$ has the following properties:

$$(5.8) \quad |\Phi^\pm(y) \mp \sqrt{y_1^2 - |y'|^2}| \ll |y|, \quad |d_y \Phi^\pm(y) \mp \frac{(y_1, -y')}{\sqrt{y_1^2 - |y'|^2}}| \ll 1, \quad |d_y^2 \Phi^\pm(y)| \lesssim \frac{1}{|y|}$$

Additionally, for $\pm s \in (0, 2)$, $\Lambda_t^{\pm, s}$ is a convex hypersurface in \mathbb{R}^n and for $x, y \in \Lambda_5^{\pm, s}$ we have

$$(5.9) \quad \left| \frac{d_x \Phi^\pm(x)}{|d_x \Phi^\pm(x)|} - \frac{d_y \Phi^\pm(y)}{|d_y \Phi^\pm(y)|} \right| \approx \pm s^{-1} |x - y|.$$

Lastly, if $y \in \Omega_4^\pm \cap \{z : |z| \leq \frac{5}{2}\}$, and $\pm s \in [0, 2)$, then

$$(5.10) \quad d(y, \Lambda_5^{\pm, s}) \approx |s - \Phi^\pm(y)|.$$

Proof. In what follows, we drop the ‘ \pm ’ and prove the claim for the case corresponding to the forward light cone, as the case of the backward light cone is similar. As in the proof of Theorem 4.2 of [8], we will prove the estimates in (5.8) for $\tilde{\Lambda}_5$, the part of Λ_5 corresponding to $\frac{1}{2} < t < \frac{5}{2}$, as a scaling argument over smaller dyadic intervals will handle the remaining part. Given positive numbers $r_1, r_2 > 0$, let $A_{(r_1, r_2)} = \{z : r_1 < |z| < r_2\}$. For z in the polar section $\Omega_5 \cap A_{(\frac{1}{2}, \frac{5}{2})}$, let $z = r\omega$ where $(r, \omega) \in (\frac{1}{2}, \frac{5}{2}) \times \mathbb{S}^{n-1} \cap \Omega_5$ and set

$$y(z) = x_r(0, \omega) \quad \xi(z) = \xi_r(0, \omega)$$

Let $\Psi(\xi) = \frac{(\xi_1, -\xi')}{\sqrt{\xi_1^2 - |\xi'|^2}}$, $T = \text{diag}(1, -1, \dots, -1)$ be an $n \times n$ matrix, $\tilde{\Pi}(\xi) = D\Psi(\xi) = \frac{T}{(\xi_1^2 - |\xi'|^2)^{\frac{1}{2}}} - \frac{(T\xi) \cdot (T\xi)^T}{(\xi_1^2 - |\xi'|^2)^{\frac{3}{2}}}$, and Π_ω^\perp denote orthogonal projection onto the subspace normal to ω . We then have that as a consequence of (5.3)

$$(5.11) \quad |\partial_r y - \Psi(\omega)| \ll 1, \quad |\partial_\omega y - r\tilde{\Pi}(\omega)| \ll 1, \quad |\partial_r \xi| \ll 1, \quad |\partial_\omega \xi - \Pi_\omega^\perp| \ll 1$$

which implies that over $\Omega_5 \cap A_{(\frac{1}{2}, \frac{5}{2})}$

$$\|y(z) - r\Psi(\omega)\|_{C^1} \ll 1, \quad \|\xi(z) - \omega\|_{C^1} \ll 1.$$

Hence the map $z \rightarrow y(z)$ is a C^1 diffeomorphism from $\Omega_5 \cap A_{(\frac{1}{2}, \frac{5}{2})}$ onto its image $y(\Omega_5 \cap A_{(\frac{1}{2}, \frac{5}{2})})$.

Hence $\Phi(y) = r$ is a well defined function on $y(\Omega_5 \cap A_{(\frac{1}{2}, \frac{5}{2})})$ such that $\{(y, \Phi(y)) : y \in y(\Omega_5 \cap A_{(\frac{1}{2}, \frac{5}{2})})\} = \tilde{\Lambda}_5$. By choosing c_0 sufficiently small in (1.13), $y(z)$ is a small perturbation of the map $r\Psi(\omega)$, meaning we can also assume that for $r \in (\frac{1}{2}, \frac{5}{2})$

$$(5.12) \quad y(\Omega_1 \cap \{z : |z| = r\}) \subset \Omega_2 \cap \Lambda^r \\ \subset y(\Omega_3 \cap \{z : |z| = r\}) \subset \Omega_4 \cap \Lambda^r \subset y(\Omega_5 \cap \{z : |z| = r\})$$

and that there is a positive distance between the boundary of any two of these sets.

The first estimate in (5.8) now follows by simple pointwise approximations. As in [8], we also have that the gradient of Φ is given by

$$(5.13) \quad (d_y \Phi)(y(z)) = \frac{\xi(z)}{H(y(z), r, \xi(z))}$$

since $\xi \cdot \partial_\omega y = 0$ and $\xi \cdot \partial_r y = H(y, r, \xi)$. The expression on the right is C^1 close to $\xi(z)/\sqrt{\xi_1(z)^2 - |\xi'(z)|^2}$ which is in turn C^1 close to $(Ty)/\sqrt{y_1^2 - |y'|^2}$. The remaining estimates in (5.8) now follow for $\frac{1}{2} < |z| < \frac{5}{2}$.

Suppose that $y \in \Lambda_5^1$, so that $y = y(\omega)$ for some $\omega \in \Omega_5$. By (5.13), we have that

$$\frac{d_y \Phi(y)}{|d_y \Phi(y)|} = \frac{\xi(\omega)}{|\xi(\omega)|}.$$

Since $|\frac{\partial \xi(\omega)}{\partial \omega} - \Pi_\omega^\perp| \ll 1$, we have that the map $\omega \mapsto \frac{\xi(\omega)}{|\xi(\omega)|}$ is a bilipschitz diffeomorphism from $\Omega_5 \cap \mathbb{S}^{n-1}$ to onto its image. $\frac{\xi(\omega)}{|\xi(\omega)|}$ is the unit normal to Λ_5^1 , so this means that the Gauss map associated to Λ_5^1 is injective, implying that it is a convex hypersurface. Since $\omega \mapsto y(\omega)$ is also bilipschitz from $\mathbb{S}^{n-1} \cap \Omega_5 \rightarrow \Lambda_5^1$, the estimate (5.9) follows for $s = 1$. The general case is now a result of a scaling argument.

To show (5.10), we observe again that by scaling we can restrict our attention to the case $s = 1$. Additionally, the claim follows easily when y is uniformly distant from the section of the hyperbola

$$\{x : x_1^2 = 1 + |x'|^2, x_1 < 0\} \cap \Omega_4$$

In other words, we can reduce the inequality to the case where

$$y \in \Omega_4 \cap A_{(\frac{3}{4}, \frac{5}{4})} \subset y(\Omega_5 \cap A_{(\frac{1}{2}, \frac{5}{2})}) \text{ and } \left| \sqrt{y_1^2 - |y'|^2} - 1 \right| < \tilde{c}_0$$

for some small $\tilde{c}_0 > 0$. Let x be a point on Λ_5^1 such that $|x - y| = d(y, \Lambda_5^1)$. Observe that by choosing the aperture of the Ω_i to be sufficiently small, we have that the restriction above implies that $|x - y| \leq .1$. We now claim that

$$(5.14) \quad |d(y, \Lambda_5^1) - |\langle Tx, y - x \rangle|| \leq \frac{1}{2}|y - x|$$

Suppose that $\omega, \tilde{\omega}$ are such that $x = x_1(0, \omega)$ and $\tilde{\omega} = \xi_1(0, \omega)/|\xi_1(0, \omega)|$. By choosing c_0 in (1.13) and \tilde{c}_0 above sufficiently small, we can assume that $\omega \in \Omega_5$ and that $|\langle \tilde{\omega}, y - x \rangle| = |y - x| = d(y, \Lambda_5^1)$. By approximating $Tx \approx \frac{\omega}{\sqrt{\omega_1^2 - |\omega'|^2}}$ and $\omega \approx \tilde{\omega}$, this now implies

$$\begin{aligned} |\langle Tx, y - x \rangle| &\leq Cc_0|y - x| + \left| \left\langle \frac{\omega}{\sqrt{\omega_1^2 - |\omega'|^2}}, y - x \right\rangle \right| \\ &\leq |\langle \tilde{\omega}, y - x \rangle| + \left(1 - \frac{1}{\sqrt{\omega_1^2 - |\omega'|^2}}\right)|y - x| + \tilde{C}c_0|y - x| \end{aligned}$$

By choosing the Ω_i to have sufficiently small angle about the negative ξ_1 -axis, this means that

$$|\langle Tx, y - x \rangle| - d(y, \Lambda_5^1) \leq \frac{1}{2}|y - x|$$

The other requisite inequality needed in (5.14) follows by similar reasoning.

Since

$$|\Phi(y) - 1| = \left| (y - x) \cdot \int_0^1 (d_y \Phi)(x + (1 - s)(y - x)) ds \right|$$

straightforward estimates on the integral involving the previous inequality yield (5.10) for $s = 1$. \square

We are now prepared to prove the desired estimates (4.9) and (4.10).

Proof of Lemma 4.1. We first observe that the estimate (4.10) follows by the proof of the analogous estimate in Lemma 3.4 in [8]. Indeed, we can write

$$K_\varepsilon(y; z, t) = \int e^{i\langle y-z, \eta \rangle} \left(e^{i\langle z, \eta \rangle - i\varphi_t(z, \eta)} a_t(z, \eta) \right) d\eta.$$

Since $\mu = 1$ implies $R \leq 1$, a simple integration by parts argument establishes the desired estimates.

We now prove (4.9) when $z = 0$ as a translation of coordinates will handle the case of arbitrary $z \in \mathbb{R}^n$. As in the analogous lemma from [8], we use a second dyadic decomposition to write $a_1(0, \eta) = \sum_\nu a^\nu(0, \eta)$ where $a^\nu(0, \eta)$ is supported in a cone of the form

$$\{\eta : |\eta|/|\eta| - \omega_\nu \leq \mu^{-\frac{1}{2}}\}.$$

The collection $\{\omega_\nu\}_\nu$ can be taken so that it is contained in Ω_1 as well as equally spaced distance $\mu^{-\frac{1}{2}}$ apart. By the estimates (5.7) on $a_1(0, \eta)$ this yields

$$|\langle \omega_\nu, \partial_\eta \rangle^j \partial_\eta^\beta a^\nu(0, \eta)| \lesssim \mu^{-j - \frac{|\beta|}{2}}$$

By homogeneity, we write

$$\varphi_1(0, \eta) = \langle (d_\eta \varphi_1)(0, \omega_\nu), \eta \rangle + R^\nu(0, \eta)$$

where

$$R^\nu(0, \eta) = \sum_{|\alpha|=2} \int_0^1 (\eta - |\eta|\omega_\nu)^\alpha \partial_\eta^\alpha \varphi_1(0, s\eta + (1-s)|\eta|\omega_\nu) ds.$$

By (5.4), this yields the estimates

$$|\langle \omega_\nu, \partial_\eta \rangle^j \partial_\eta^\beta R^\nu(0, \eta)| \lesssim \mu^{-j - \frac{|\beta|}{2}}.$$

Set $x_\nu = \bar{x}_1(0, \omega_\nu) = d_\eta \varphi_1(0, \omega_\nu)$. Integration by parts in (5.5) now yields

$$(5.15) \quad |K_\varepsilon(y; 1, 0)| \leq C_N \sum_\nu \mu^{\frac{n+1}{2}} (1 + \mu|\langle \omega_\nu, y - x_\nu \rangle| + \mu|y - x_\nu|^2)^{-N}.$$

We now set $\tilde{\Omega}_- = \Omega_2^-$. By (5.12), the distance between the complement of Ω_2^- and Λ_1^- is uniformly bounded from below, (4.9) is straightforward when either $y \notin \Omega_2^-$ or $d(y, \Lambda^-) \geq \tilde{c}_0$, hence from now on we assume $y \in \Omega_2^-$ and $d(y, \Lambda^-) \leq \tilde{c}_0$.

Given ν , set $\tilde{\omega}_\nu = d_z \varphi_1(0, \omega_\nu) / |d_z \varphi_1(0, \omega_\nu)|$ so that $x_\nu = x_{0,1}(0, \tilde{\omega}_\nu)$ and $\omega_\nu = \xi_{0,1}(0, \tilde{\omega}_\nu) / |\xi_{0,1}(0, \tilde{\omega}_\nu)|$, $(x_{0,1}(x, \omega), \xi_{0,1}(x, \omega))$ denoting the solution to Hamilton's equations (5.1) at time 0 with initial conditions (x, ω) at time 1. Since

$$\omega \mapsto x_{0,1}(0, \omega) \quad \text{and} \quad \omega \mapsto \xi_{0,1}(0, \omega) / |\xi_{0,1}(0, \omega)|$$

are both bilipschitz maps from \mathbb{S}^{n-1} onto their image we have that

$$|x_\nu - x_{\nu'}| \approx |\omega_\nu - \omega_{\nu'}|.$$

Now let $x_0 \in \Lambda_z$ be such that $d(y, \Lambda_z) = |y - x_0|$. Let ω_0 be the unit normal to Λ_z at x_0 with $(\omega_0)_1 < 0$. As in Lemma 5.2, we can choose c_0, \tilde{c}_0 small so that there exists $\tilde{\omega}_0 \in \Omega_3$ with $\omega_0 = \xi_{0,1}(x, \tilde{\omega}_0) / |\xi_{0,1}(x, \tilde{\omega}_0)|$ and $|\langle \omega_0, y - x_0 \rangle| = d(y, \Lambda_z)$. We claim that

$$(5.16) \quad |\langle \omega_0, y - x_0 \rangle| + |\omega_0 - \omega_\nu|^2 \lesssim |\langle \omega_\nu, y - x_\nu \rangle| + |y - x_\nu|^2.$$

First observe that

$$|y - x_0|^2 + |x_0 - x_\nu|^2 = |y - x_\nu|^2 + 2\langle y - x_0, x_0 - y \rangle + 2\langle y - x_0, y - x_\nu \rangle \leq 3|y - x_\nu|^2$$

and hence

$$|y - x_\nu|^2 \approx |y - x_0|^2 + |x_0 - x_\nu|^2 \approx |y - x_0|^2 + |\omega_0 - \omega_\nu|^2.$$

Now observe that

$$x_\nu - x_0 = \frac{\partial x_{0,1}}{\partial \omega}(0, \tilde{\omega}_0) \cdot (\tilde{\omega}_\nu - \tilde{\omega}_0) + E$$

where E denotes the second order error in the Taylor expansion of $\omega \mapsto x_{0,1}(0, \omega)$ about ω_0 . Since $\omega_0 \cdot \frac{\partial x_{0,1}}{\partial \omega}(0, \tilde{\omega}_0) = 0$ we have that

$$|\langle \omega_0, x_\nu - x_0 \rangle| \lesssim |\tilde{\omega}_\nu - \tilde{\omega}_0|^2 \lesssim |\omega_\nu - \omega_0|^2.$$

The claim now follows from the following inequality:

$$\begin{aligned} \left| |\langle \omega_0, y - x_0 \rangle| - |\langle \omega_\nu, y - x_\nu \rangle| \right| &\lesssim |\langle \omega_0 - \omega_\nu, y - x_\nu \rangle| + |\langle \omega_0, x_\nu - x_0 \rangle| \\ &\lesssim |y - x_\nu|^2 + |\omega_0 - \omega_\nu|^2. \end{aligned}$$

Since $|\langle \omega_0, y - x_0 \rangle| = d(y, \Lambda_z)$, we can now use (5.16) to control the sum in (5.15) to prove the Lemma. \square

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