

# $L^p$ norms and global harmonic analysis

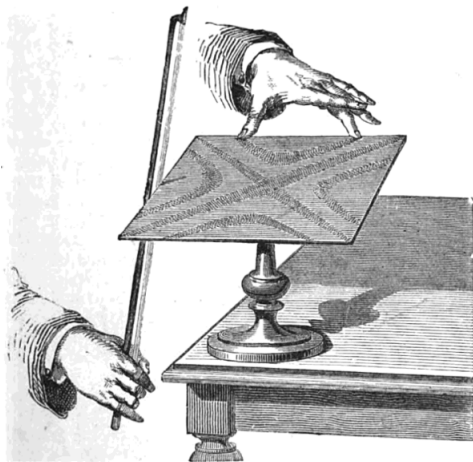
Matthew D. Blair

University of New Mexico

September 14, 2017

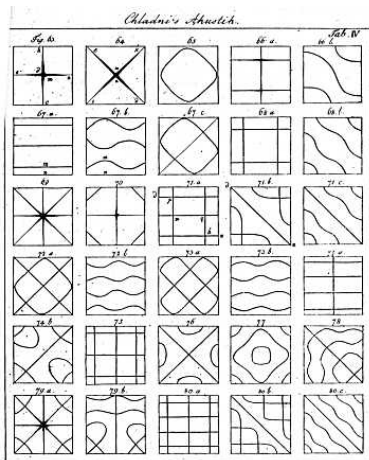
Research partially supported by NSF grants: DMS-1301717  
and DMS-1565436

# Chladni's experiment (circa 1787)



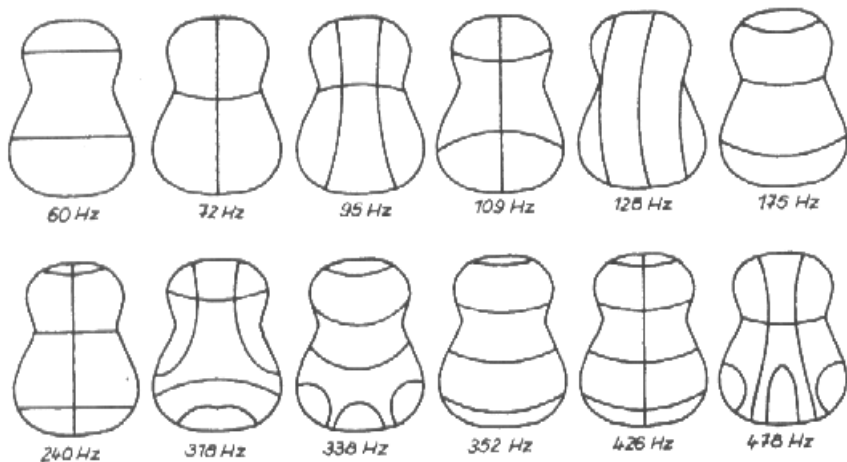
Bowing a dusted plate

# Chladni figures on a square plate



From the book *Die Akustik* by Ernst Chladni

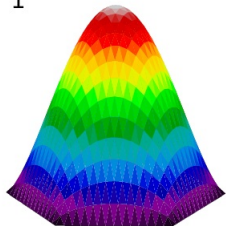
# Changing the shape of the plate



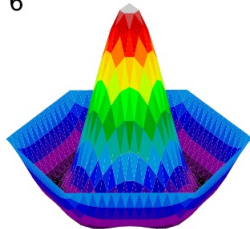
Uploaded to Wikipedia by Denis Diderot

# Amplitudes of vibrational modes

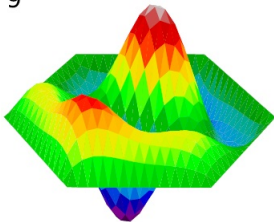
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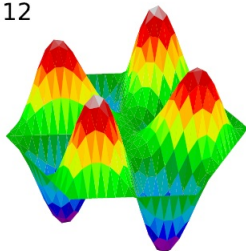
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9



12



# Helmholtz equation

- Led to study the Helmholtz equation on a domain

$$(\lambda^2 + \Delta)\varphi_\lambda(x) = 0$$

- Every  $\varphi_\lambda$  gives rise to solutions to the wave equation

$$(-\partial_t^2 + \Delta)(e^{\pm it\lambda}\varphi_\lambda(x)) = 0$$

- Fermat's principle suggests the mass of  $\varphi_\lambda$  should in some sense be invariant under the billiard dynamics (paths of least action) on the plate as  $\lambda \rightarrow \infty$

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- For a compact domain  $M$  and  $f : M \rightarrow \mathbb{C}$  continuous, define

$$\|f\|_{L^p(M)} := \left( \int_M |f(x)|^p dx \right)^{\frac{1}{p}}, \quad 1 \leq p < \infty$$

$$\|f\|_{L^\infty(M)} := \max_{x \in M} |f(x)|,$$

- Helmholtz sol'ns often normalized to have norm 1 in  $L^2(M)$
- For  $2 < p < \infty$ ,  $L^p$  norms will be sensitive to both the extrema of the function and the regions where the extrema is nearly obtained



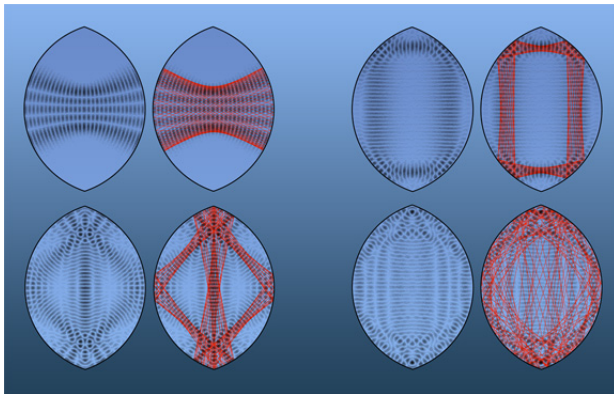
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# Scarring



By Eric Heller

# Laplacian on a compact manifold

- Let  $(M, g)$  be a  $C^\infty$ , *boundaryless*, compact Riemannian manifold,  $\dim(M) = n \geq 2$
- Let  $\Delta_g$  be the Laplace-Beltrami operator, in coordinates

$$\Delta_g u = \sum_{i,j=1}^n \frac{1}{\sqrt{\det g_{kl}}} \partial_i \left( g^{ij} \sqrt{\det g_{kl}} \partial_j u \right)$$

- Self-adjoint w.r.t. Riemannian measure  $dV_g$

$$\begin{aligned} \int_M \Delta_g u \bar{v} \, dV_g &= \int_M u \overline{\Delta_g v} \, dV_g, \\ \int_M \Delta_g u \bar{u} \, dV_g &= - \int_M |\nabla_g u|^2 \, dV_g \end{aligned}$$

# Eigenfunctions on a compact manifold

- Compactness of  $M \implies$  spectrum of  $\Delta_g$  is discrete: there exists a sequence  $\{\varphi_{\lambda_j}\}_{j=1}^{\infty}$  of eigenfunctions forming an O.N. basis for  $L^2(M)$

$$\Delta_g \varphi_{\lambda_j} = -\lambda_j^2 \varphi_{\lambda_j}, \quad \sqrt{-\Delta_g} \varphi_{\lambda_j} = \lambda_j \varphi_{\lambda_j}$$

- Eigenspaces are finite dimensional, but may have high multiplicity (perhaps even growing with  $\lambda$ )
- Notation:  $\varphi_{\lambda}$  is any admissible solution

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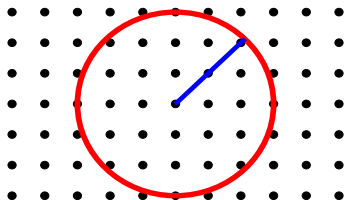
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# Example: Flat torus

- Let  $\mathbb{T}^n = \mathbb{R}^n / (2\pi\mathbb{Z})^n$ , endowed with the flat metric (a.k.a.  $[-\pi, \pi]^n$  with opposite sides “glued” together)
- Laplacian is the familiar  $\Delta u = \sum_1^n \frac{\partial^2 u}{\partial x_j^2}$
- Eigenfunctions are linear combos of Fourier modes

$$\varphi_\lambda(x) = \frac{1}{(2\pi)^n} \sum_{|m|=\lambda} a_m e^{ix \cdot m}, \quad \sum_{|m|=\lambda} |a_m|^2 = 1 \quad (m \in \mathbb{Z}^n)$$



Radius =  $\lambda$

# Example: $\mathbb{S}^n$ , canonical sphere in $\mathbb{R}^{n+1}$

- On  $\mathbb{S}^2$ , with spherical coordinates  $(\theta, \phi)$

$$\Delta_g u = \frac{1}{\sin \theta} \partial_\theta (\sin \theta \partial_\theta u) + \frac{1}{\sin^2 \theta} \partial_\phi^2 u$$

- E'functions are linear combos of  $Y_l^m(\theta, \phi) = e^{im\phi} P_l^m(\cos \theta)$  ( $P_l^m$  an associated Legendre polynomial)
- For  $n \geq 2$ , spectrum of  $\Delta_g$  on  $\mathbb{S}^n$  is:

$$\text{spec}(\sqrt{-\Delta_g}) = \left\{ \sqrt{k(k+n-1)} = k + \mathcal{O}(k^{-1}) : k = 0, 1, 2, \dots \right\}$$

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# $L^p$ bounds on eigenfunctions

- Main problem: given  $(M, g)$  and the unit sphere in an eigenspace

$$V_\lambda = \{\varphi_\lambda : \Delta_g \varphi_\lambda = -\lambda^2 \varphi_\lambda, \|\varphi_\lambda\|_{L^2} = 1\}$$

find (best possible) upper bounds on

$$\sup_{\varphi_\lambda \in V_\lambda} \|\varphi_\lambda\|_{L^p(M)}, \quad 2 < p \leq \infty$$

- Might expect the following type, for some  $\delta(p) > 0$

$$\|\varphi_\lambda\|_{L^p(M)} \leq C\lambda^{\delta(p)} \quad \text{for } \|\varphi_\lambda\|_{L^2(M)} = 1$$

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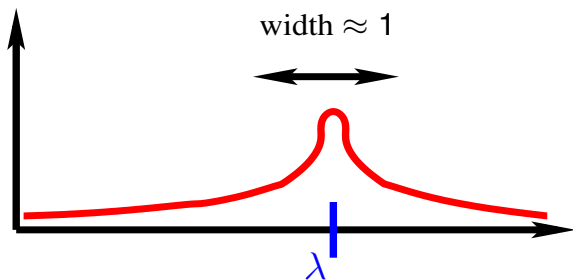
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# Wave approach to eigenfunctions (see Sogge, *Fourier Integrals in Classical Analysis*)

- Let  $\chi \in \mathcal{S}(\mathbb{R})$ , satisfy  $\text{supp}(\widehat{\chi}) \subset (-\epsilon, \epsilon)$ ,  $\chi(0) = 1$ , define

$$\chi(\lambda - \sqrt{-\Delta_g})f = \sum_{j=1}^{\infty} \chi(\lambda - \lambda_j) \langle f, \varphi_{\lambda_j} \rangle \varphi_{\lambda_j}$$

so that  $\chi(\lambda - \sqrt{-\Delta_g})\varphi_\lambda = \varphi_\lambda$



# Wave approach to eigenfunctions

- Formally, using Fourier integrals

$$\chi(\lambda - \sqrt{-\Delta_g}) = \frac{1}{2\pi} \int_{-\epsilon}^{\epsilon} e^{it\lambda} e^{-it\sqrt{-\Delta_g}} \widehat{\chi}(t) dt$$

- $e^{-it\sqrt{-\Delta_g}}$  is the half-wave evolution, a solution map for the wave eqn, i.e.

$$e^{-it\sqrt{-\Delta_g}} f = \sum_{j=1}^{\infty} e^{-it\lambda_j} \langle f, \varphi_{\lambda_j} \rangle \varphi_{\lambda_j}$$

- Lax parametrix + stationary phase yields

$$\chi(\lambda - \sqrt{-\Delta_g}) v(x) = \int_M e^{i\lambda d_g(x,y)} a_\lambda(x,y) v(y) dy + \text{small error}$$
$$\text{supp}(a_\lambda) \subset \{(x,y) \in M \times M : d_g(x,y) \leq \epsilon\}$$

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# Limitations to the wave approach

- The dream: replace  $\chi(\lambda - \sqrt{-\Delta_g})$  by  $\tilde{\chi}_\lambda(\sqrt{-\Delta_g})$  where

$$\text{supp}(\tilde{\chi}_\lambda) \cap \text{spec}(\sqrt{-\Delta_g}) = \{\lambda\}$$

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- Fourier transform of  $\tilde{\chi}_\lambda$  loses the compact support, so must understand the wave kernel for all  $t \in \mathbb{R}$
- Weyl's law and its consequences:

$$\#\{\lambda_j : \lambda_j \leq \lambda\} = c_n \lambda^n + \mathcal{O}(\lambda^{n-1})$$

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so *lots* of eigenfunctions in the range of  $\chi(\lambda - \sqrt{-\Delta_g})$ , typically a *crude* approximation to eigenspace projection

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# $L^p$ bounds on spectral clusters

Theorem: “Universal”  $L^p$  bounds (Sogge ‘88)

If  $\dim(M) = n \geq 2$  and  $g \in C^\infty$ , then

$$\|\chi(\lambda - \sqrt{-\Delta_g})f\|_{L^p(M)} \leq C\lambda^{\delta(p)}\|f\|_{L^2(M)}$$

$$\delta(p) = \begin{cases} \frac{n-1}{2}(\frac{1}{2} - \frac{1}{p}), & 2 \leq p \leq \frac{2(n+1)}{n-1} \\ \frac{n-1}{2} - \frac{n}{p}, & \frac{2(n+1)}{n-1} \leq p \leq \infty \end{cases}$$

Consequently,  $\|\varphi_\lambda\|_{L^p(M)} = \mathcal{O}(\lambda^{\delta(p)})$  as  $\lambda \rightarrow \infty$

- At best, growth rate  $\delta(p)$  is sharp for the operator  $\chi(\lambda - \sqrt{-\Delta_g})$ , but not necessarily eigenspace projections
- $\mathbb{S}^n$  is exceptional:  
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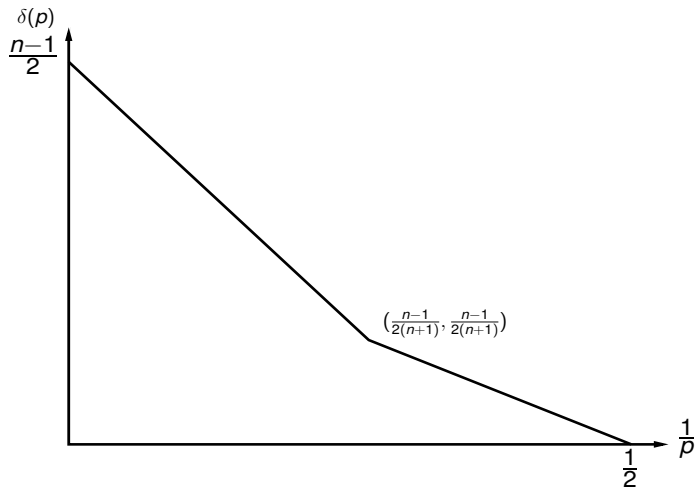
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# Zonal harmonics on $\mathbb{S}^n$

- Sogge: On  $\mathbb{S}^n$ , can find  $\varphi_\lambda$  satisfying for enough  $x$

$$|\varphi_\lambda(x)| \approx \lambda^{\frac{n-1}{2}} (1 + \lambda d_g(x, x_0))^{-\frac{n-1}{2}} \quad \text{near } x_0 \in M,$$

hence the  $\delta(p) = \frac{n-1}{2} - \frac{n}{p}$  cannot be improved for  $\frac{2(n+1)}{n-1} \leq p \leq \infty$

$$\|\varphi_\lambda\|_{L^p(M)} \approx \lambda^{\frac{n-1}{2} - \frac{n}{p}}$$

- On  $\mathbb{S}^n$ , these are zonal harmonics (exact eigenfunctions)

# Highest weight/sectoral harmonics on $\mathbb{S}^n$

- $\varphi_\lambda$  concentrated in a  $\lambda^{-\frac{1}{2}}$ -nbd of a closed geodesic  $\gamma$

$$\mathcal{T}_{\lambda^{-1/2}}(\gamma) = \{x : d_g(x, \gamma) \lesssim \lambda^{-\frac{1}{2}}\},$$

hence the  $\delta(p) = \frac{n-1}{2}(\frac{1}{2} - \frac{1}{p})$  growth rate cannot be improved for  $2 < p \leq \frac{2(n+1)}{n-1}$

$$\|\varphi_\lambda\|_{L^p(M)} \approx \left[ \text{Vol}(\mathcal{T}_{\lambda^{-1/2}}(\gamma)) \right]^{\frac{1}{p} - \frac{1}{2}} \approx \lambda^{\frac{n-1}{2}(\frac{1}{2} - \frac{1}{p})}$$

- On  $\mathbb{S}^n$ , these are the *highest weight harmonics*
- Underscores that  $p = \frac{2(n+1)}{n-1}$  is the *critical* exponent. *Both* mass concentration profiles saturate the growth rate.
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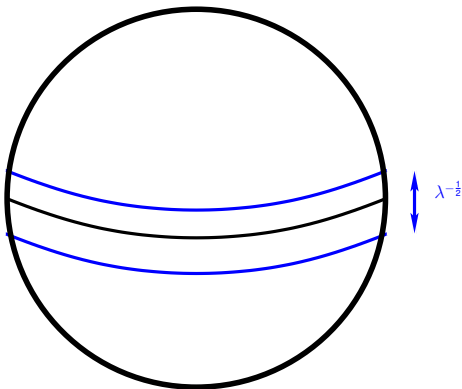
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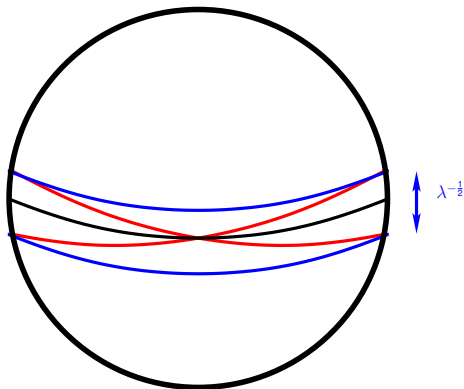
# Highest weight harmonics on $\mathbb{S}^n$

- Mass of  $\varphi_\lambda$  concentrated in a  $\lambda^{-\frac{1}{2}}$  neighborhood of the equator



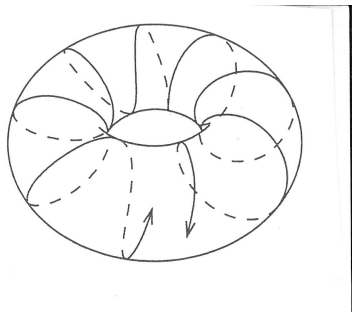
# Highest weight harmonics on $\mathbb{S}^n$

- Equator is a stable, degenerate orbit



# Flat and negatively curved manifolds

- No such stable orbits on the torus: e.g. image of lines with irrational slope under the quotient map are dense

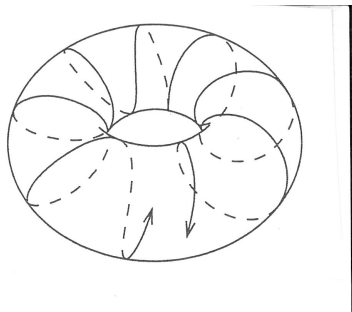


(Image from J. Lee, *Introduction to Smooth Manifolds*)

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Main question for general  $(M, g)$ : When can the  $\mathcal{O}(\lambda^{\delta(p)})$  bounds be improved?

- Zygmund:  $\|\varphi_\lambda\|_{L^4(\mathbb{T}^2)} \leq C$
- Classical number theory:  $\|\varphi_\lambda\|_{L^\infty(\mathbb{T}^2)} \leq C_\epsilon \lambda^\epsilon$ , any  $\epsilon > 0$
- Bourgain, Bourgain-Demeter: improvements on  $\delta(p)$  exponent for flat tori
- Iwaniec-Sarnak: power improvements for certain (noncompact) arithmetic quotients of the hyperbolic plane
- Sogge-Zelditch: maximal  $L^\infty$  growth occurs only when "full measure" family of geodesics issuing from a point loop back at a common time

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Main question for general  $(M, g)$ : When can the  $\mathcal{O}(\lambda^{\delta(p)})$  bounds be improved?

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# Nonpositive curvature

- Negative curvature: geodesic flow behaves chaotically, stable and unstable manifolds are invariant under geodesic flow (Anosov)
- $L^p$  bounds expected to be *much* better than predicted by Sogge's bounds (perhaps even  $\mathcal{O}(\lambda^\epsilon)$  for any  $\epsilon > 0$ )
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# Wave approach for nonpositive curvature

- As before  $\chi \in \mathcal{S}(\mathbb{R})$  with  $\widehat{\chi} \in C^\infty$  and compactly supported in  $(-\epsilon, \epsilon)$ ,  $\chi(0) = 1$
- This time consider with  $T = T(\lambda) = \log \lambda$

$$\chi(T(\lambda - \sqrt{-\Delta_g})) = \frac{1}{2\pi T} \int_{-\epsilon T}^{\epsilon T} e^{it\lambda} e^{-it\sqrt{-\Delta_g}} \widehat{\chi}\left(\frac{t}{T}\right) dt$$

- Yields a gain of  $\frac{1}{T}$ , but have to analyze  $e^{-it\sqrt{-\Delta_g}}$  over large time scales
- Can lift the problem to the universal cover, treating it as a “periodic” problem and use the Hadamard parametrix
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- Coherent states/wave packets at frequency  $\lambda$ , consider Gaussian concentrated at scales  $\Delta x$  in space,  $\Delta \xi$  in frequency
- Take localization scales  $\Delta x \approx \lambda^{-1/2}$ ,  $\Delta \xi \approx \lambda^{1/2}$ , finest possible while respecting uncertainty principle and (approximate) invariance under bicharacteristic flow
- For unstable flows, wave packets lose their coherent structure when  $|t| \gg \log \lambda$ , e.g.

$$\Delta x_t + \lambda^{-1} \Delta \xi_t \leq C e^{\mu t} (\Delta x + \lambda^{-1} \Delta \xi) \approx e^{\mu t} \lambda^{-\frac{1}{2}}$$

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# Improvements beneath the critical exponent

Thm (B.-Sogge '15, to app. *Comm. Math. Phys., J. Diff. Geom.*)

Suppose  $(M, g)$  has nonpositive sectional curvatures. Then for any  $2 < p < \frac{2(n+1)}{n-1}$ , there exists an exponent  $\sigma = \sigma(p, n) > 0$  such that

$$\|\varphi_\lambda\|_{L^p(M)} = \mathcal{O}\left(\frac{\lambda^{\delta(p)}}{(\log \lambda)^\sigma}\right)$$

- Showed that mass does *not* concentrate in  $\lambda^{-1/2}$  neighborhoods
- Earlier work: Sogge, Sogge-Zelditch, B.-Sogge
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# Improvements at the critical exponent

## Theorem (B.-Sogge '17)

Suppose  $(M, g)$  has nonpositive sectional curvatures. Then there exists an exponent  $\sigma = \sigma(n) > 0$  such that

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# Restrictions to geodesic segments

Theorem (B.'16, to appear *Israel J. Math.*)

Suppose  $(M, g)$  has nonpositive sectional curvatures and  $n = 2$ . Then for any unit length geodesic segment  $\gamma$

$$\|\varphi_\lambda|_\gamma\|_{L^4(\gamma)} = \mathcal{O}\left(\frac{\lambda^{1/4}}{(\log \lambda)^{1/4}}\right)$$

- Improves universal bounds for geodesic restrictions of Burq-Gérard-Tzvetkov, Hu, Chen-Sogge
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